

Neutrino's Non-Standard Interactions; another eel under a willow?



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Second eel (loach) under a willow?

- = more new physics from neutrinos?



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What is the natural time scale for discovery of new ν interactions?;

caution

- Meitner-Hahn (1911) – discovery of Neutral current (1973)
- Discovery of neutrino (1953) – discovery of neutrino mass (1998)



~ 50 years



ν 's Non-standard interactions (NSI)

- If there exists **NEW PHYSICS** at TeV scale there might be NSI of ν 's expressed by higher dimensional operators

$$\mathcal{L}_{\text{eff}}^{\text{NSI}} = -2\sqrt{2} \varepsilon_{\alpha\beta}^{\text{fP}} G_F (\bar{\nu}_\alpha \gamma_\mu P_L \nu_\beta) (\bar{f} \gamma^\mu P f),$$

$$P_L \equiv \frac{1}{2}(1 - \gamma_5) \text{ or } P_R \equiv \frac{1}{2}(1 + \gamma_5).$$

Wolfenstein, Grossman, Berezhiani-Rossi,
Davidson et al. ... many people

- But coefficient ε may be small  on
dimensional ground, $\varepsilon \sim \text{order } (M_W/M_{\text{NP}})^2$
 $(M_W/M_{\text{NP}})^2 = 0.01 \text{ (0.0001) if } M_{\text{NP}} = 1 \text{ (10) TeV}$

I report the progress in the last one year !

NSI theory; Some recent progress

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Some relevant references

Z. Berezhiani and A. Rossi, Phys. Lett. B **535**, 207 (2002), hep-ph/0111137.

S. Davidson, C. Pena-Garay, N. Rius and A. Santamaria, JHEP **03**, 011 (2003), hep-ph/0302093.

S. Antusch, J. P. Baumann and E. Fernandez-Martinez, Nucl. Phys. B **810**, 369 (2009), 0807.1003.

M. B. Gavela, D. Hernandez, T. Ota and W. Winter, Phys. Rev. D **79**, 013007 (2009), 0809.3451.

C. Biggio, M. Blennow and E. Fernandez-Martinez, arXiv:0902.0607 [hep-ph].

Note; many original references omitted, sorry!

Gauge invariance and NSI

- At high scale where NSI are originated, there exists SU(2) x U(1) gauge invariance
- Therefore, if there is a d=6 operator

$$\frac{1}{\Lambda^2}(\bar{\nu}_\alpha \gamma^\rho P_L \nu_\beta)(\bar{\ell}_\gamma \gamma_\rho \ell_\delta).$$

$\epsilon_{e\mu}^{ee}$

→ it must exist as a part of the gauge invariant operator

$$\frac{1}{\Lambda^2}(\bar{L}_\alpha \gamma^\rho L_\beta)(\bar{L}_\gamma \gamma_\rho L_\delta),$$

- But it involves 4 charged lepton operators

→ severe constraints from experiments

March 12, 2009 $\epsilon_{e\mu}^{ee} < 10^{-6}$ Neutrino Telescope 09

$\text{Br}(\mu \rightarrow eee) < 10^{-12}$

Smart way to avoid constraints from 4 lepton processes

- Anti-symmetrization; dim 6

Unique; Gavela
et al. 08

$$\begin{aligned}\mathcal{O}_6^a &= (\bar{L}_\gamma i\tau_2 L_\alpha^c)(\bar{L}_\beta^c i\tau_2 L_\delta) \\ &= (\bar{\ell}_\alpha \gamma^\mu \ell_\beta)(\bar{\nu}_\gamma \gamma_\mu \nu_\delta) + (\bar{\ell}_\gamma \gamma^\mu \ell_\delta)(\bar{\nu}_\alpha \gamma_\mu \nu_\beta) - (\bar{\ell}_\alpha \gamma^\mu \ell_\delta)(\bar{\nu}_\gamma \gamma_\mu \nu_\beta) - (\bar{\ell}_\gamma \gamma^\mu \ell_\beta)(\bar{\nu}_\alpha \gamma_\mu \nu_\delta)\end{aligned}$$

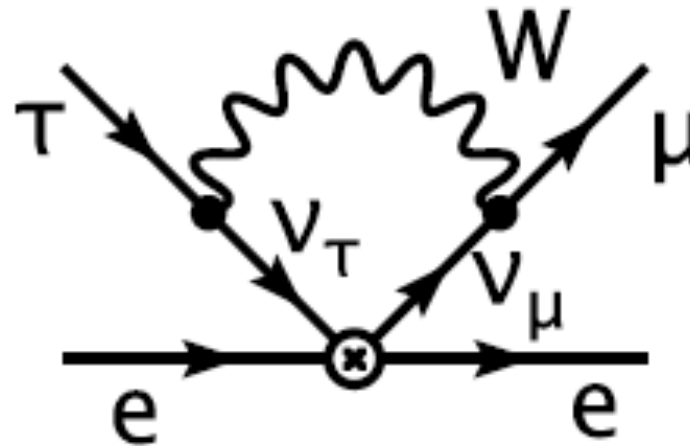
- Higgs projection; dim 8

$$: 2(\bar{L}_\beta \tilde{H})\gamma^\rho(\tilde{H}^\dagger L_\alpha)(\bar{L}_\delta \gamma_\rho L_\gamma) :$$

- Now, are we free from any constraints from
4 lepton processes?

Constraints from loops

- Dressing by SM particles gives you the bound!; loop level constraint



Davidson et al. 03

Figure 4: One-loop contributions to four-fermion interactions in the effective theory.



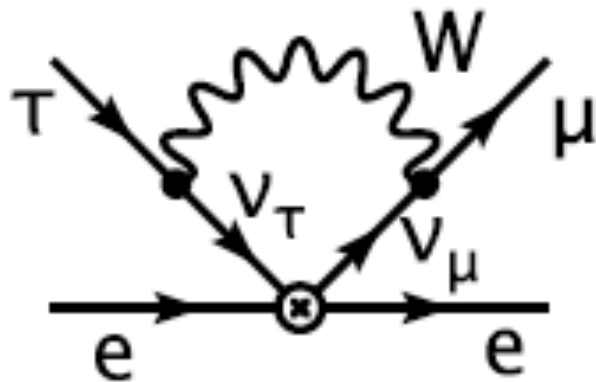
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$(\bar{e}\gamma^\rho P e)(\bar{\nu}_\tau \gamma_\rho L \nu_\mu)$	$ \epsilon_{\tau\mu}^{eP} < 1.2$ $(\tau \rightarrow \mu \bar{e} e)^*$ $ \epsilon_{\tau\mu}^{eP} < 0.1$ CHARM II	$ \epsilon_{\tau\mu}^{eL} < 0.04, \epsilon_{\tau\mu}^{eR} < 0.02$ leptonic s_W^2 at nufact
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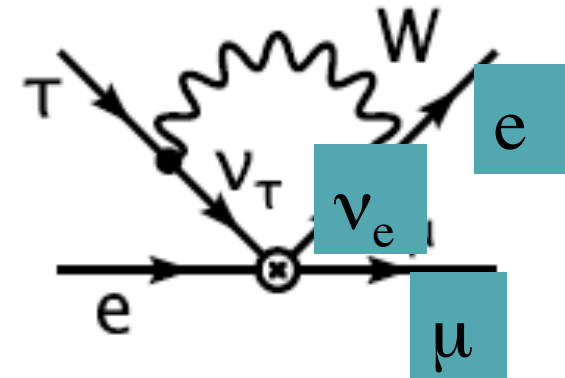
Constraints from loops

- Dressing by SM particles gives you the bound!; loop level constraint

Biggio, Blennow and Fernandez-Martinez (arXiv: 0902.0607)



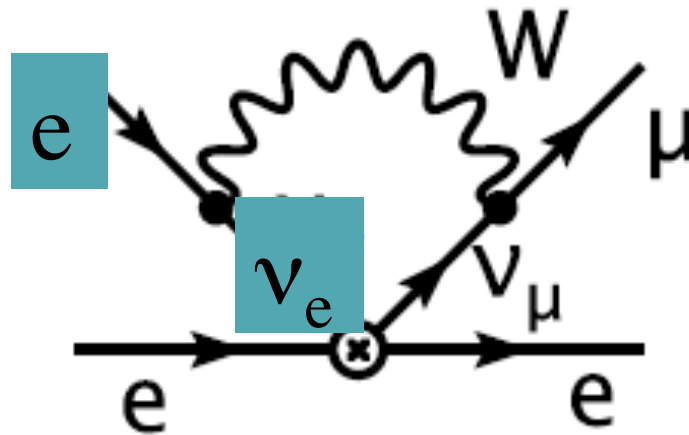
+ 3 other diagrams with anti-symmetry such as =>



$(\bar{e}\gamma^\rho P_L e)(\bar{\nu}_\tau \gamma_\rho L \nu_\mu)$	<p>Bound is loosened by factor $\sim 10^3$</p>
--	---

Bound on $\varepsilon^{\text{matter}}_{e\mu}$ is gone

- Used to be the most stringent loop level constraint
- But now, due to anti-symmetry, there is no such diagram as:



Biggio,
Blennow and
Fernandez-
Martinez
(0902.0607)

→ $(\bar{e}\gamma^\rho P e)(\bar{\nu}_\mu \gamma_\rho L \nu_e)$

GONE!

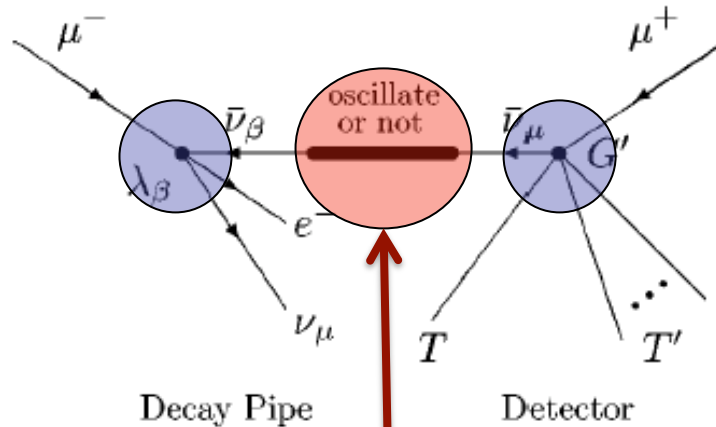
ν
oscillation
with NSI



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Non Standard Interaction (NSI)



- It comes in into 3 places, production, propagation, & detection

Grossmann, Ota-Sato-Yamashita, Huber et al. ...

I this talk, I concentrate on effects of NSI in ν propagation in matter governed by evolution eq.

$$i \frac{d}{dt} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \frac{1}{2E} \left[U \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger + a \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix} \right] \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

ν oscillation with NSI


- Given the structure

$$i \frac{d}{dt} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \frac{1}{2E} \left[U \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger + a \begin{pmatrix} 1 + \varepsilon_{ee} & \varepsilon_{e\mu} & \varepsilon_{e\tau} \\ \varepsilon_{e\mu}^* & \varepsilon_{\mu\mu} & \varepsilon_{\mu\tau} \\ \varepsilon_{e\tau}^* & \varepsilon_{\mu\tau}^* & \varepsilon_{\tau\tau} \end{pmatrix} \right] \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

➡ the system is highly nontrivial

- Note: when the new physics at TeV scale is identified, it will give us a few parameters
- all $\varepsilon_{\alpha\beta}$ are related, but likely to be present with similar magnitudes
- But, nobody seems to know how large is $\varepsilon_{\alpha\beta}$ effect in $P_{\gamma\delta}$ ➡ perturbation theory of ν

Formulating perturbation theory of ν oscillation

- We have no definite recipe because we still do not know how large is θ_{13}
- Only known small parameter is $\Delta m^2_{\text{solar}}/\Delta m^2_{\text{atm}} = r_{\Delta} \sim 0.03$  various possibilities (Note; $\sin(\pi/4 - \theta_{23}) < 0.14$)

$$s_{13} \sim \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \simeq 0.03 \quad \rightarrow \quad \sin^2 2\theta_{13} \simeq 3.6 \times 10^{-3}$$





$$s_{13} \sim \left(\frac{\Delta m_{21}^2}{\Delta m_{31}^2} \right)^{0.85} \simeq 0.05 \quad \rightarrow \quad \sin^2 2\theta_{13} \simeq 0.01$$

$$s_{13} \sim \sqrt{\frac{\Delta m_{21}^2}{\Delta m_{31}^2}} \simeq 0.17 \quad \rightarrow \quad \sin^2 2\theta_{13} \simeq 0.12$$

ε perturbation theory

- I take the assumption

$$s_{13} \sim \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \simeq 0.03 \quad \rightarrow \quad \sin^2 2\theta_{13} \simeq 3.6 \times 10^{-3}$$


- I assume $a/\Delta m_{\text{atm}}^2 \sim \Delta m_{\text{atm}}^2 L / E \sim O(1)$
- most natural perturbation theory of ν oscillation  $P_{e\mu}$ consists only of order ε^2 terms  widely used Cervera et al. formula (golden measurement paper)
- (If θ_{13} large, one may try $(s_{13})^2 \sim r_{\Delta} = \varepsilon$ 
 δ term $\sim \varepsilon^{3/2}$  $(s_{13})^3$ term required)

ε perturbation theory with NSI

- We include NSI by assuming

$$\frac{\Delta m_{21}^2}{\Delta m_{31}^2} \sim s_{13} \sim \varepsilon_{\alpha\beta} \sim \epsilon \quad (\alpha, \beta = e, \mu, \tau).$$

 and expand to ε^2 ; NSI 2nd order formula

- If dimension 8, 1st order formula would be enough
but, 2nd order formula is simpler !
- Are you happy with $\sin^2 2\theta_{13} \sim 0.004$?
- Well, in fact I argue sometimes that θ_{13} is large, e.g., reactor range
- Are you schizophrenia?  Yes!

SI formula -> NSI 2nd order formula

$$\begin{aligned}
 P(\nu_e \rightarrow \nu_\mu) = & 4c_{23}^2 \left[c_{12}s_{12} (\Delta m_{21}^2 / a) \right]^2 \sin^2 \frac{aL}{4E} \\
 & + 4s_{23}^2 \left[s_{13} e^{-i\delta} (\Delta m_{31}^2 / a) \right]^2 \left(\frac{a}{\Delta m_{31}^2 - a} \right)^2 \sin^2 \frac{\Delta m_{31}^2 - a}{4E} L \\
 & + 8c_{23}s_{23} \operatorname{Re} \left[c_{12}s_{12} (\Delta m_{21}^2 / a) \right] \left[s_{13} e^{-i\delta} (\Delta m_{31}^2 / a) \right] \\
 & \quad \times \frac{a}{\Delta m_{31}^2 - a} \sin \frac{aL}{4E} \cos \frac{\Delta m_{31}^2 L}{4E} \sin \frac{\Delta m_{31}^2 - a}{4E} L \\
 & + 8c_{23}s_{23} \operatorname{Im} \left[c_{12}s_{12} (\Delta m_{21}^2 / a) \right] \left[s_{13} e^{-i\delta} (\Delta m_{31}^2 / a) \right] \\
 & \quad \times \frac{a}{\Delta m_{31}^2 - a} \sin \frac{aL}{4E} \sin \frac{\Delta m_{31}^2 L}{4E} \sin \frac{\Delta m_{31}^2 - a}{4E} L
 \end{aligned}$$

Cervera et al , hep-ph/0002108

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SI formula -> NSI 2nd order formula

$$\begin{aligned}
 P(\nu_e \rightarrow \nu_\mu) = & 4c_{23}^2 \left[c_{12}s_{12}\frac{\Delta m_{21}^2}{a} + c_{23}\varepsilon_{e\mu} - s_{23}\varepsilon_{e\tau} \right]^2 \sin^2 \frac{aL}{4E} \\
 & + 4s_{23}^2 \left[s_{13}e^{-i\delta}\frac{\Delta m_{31}^2}{a} + s_{23}\varepsilon_{e\mu} + c_{23}\varepsilon_{e\tau} \right]^2 \left(\frac{a}{\Delta m_{31}^2 - a} \right)^2 \sin^2 \frac{\Delta m_{31}^2 - a}{4E} L \\
 & + 8c_{23}s_{23} \operatorname{Re} \left[\left(c_{12}s_{12}\frac{\Delta m_{21}^2}{a} + c_{23}\varepsilon_{e\mu} - s_{23}\varepsilon_{e\tau} \right) \left(s_{13}e^{i\delta}\frac{\Delta m_{31}^2}{a} + s_{23}\varepsilon_{e\mu}^* + c_{23}\varepsilon_{e\tau}^* \right) \right] \\
 & \quad \times \frac{a}{\Delta m_{31}^2 - a} \sin \frac{aL}{4E} \cos \frac{\Delta m_{31}^2 L}{4E} \sin \frac{\Delta m_{31}^2 - a}{4E} L \\
 & + 8c_{23}s_{23} \operatorname{Im} \left[\left(c_{12}s_{12}\frac{\Delta m_{21}^2}{a} + c_{23}\varepsilon_{e\mu} - s_{23}\varepsilon_{e\tau} \right) \left(s_{13}e^{i\delta}\frac{\Delta m_{31}^2}{a} + s_{23}\varepsilon_{e\mu}^* + c_{23}\varepsilon_{e\tau}^* \right) \right] \\
 & \quad \times \frac{a}{\Delta m_{31}^2 - a} \sin \frac{aL}{4E} \sin \frac{\Delta m_{31}^2 L}{4E} \sin \frac{\Delta m_{31}^2 - a}{4E} L.
 \end{aligned}$$

T. Kikuchi, H.M., S. Uchinami, [arXiv:0809.3312](https://arxiv.org/abs/0809.3312) => JHEP

Generalized atmospheric and solar variables

SI-NSI
confusion !

$$\Theta_{\pm} \equiv s_{13} \frac{\Delta m_{31}^2}{a} + (s_{23}\varepsilon_{e\mu} + c_{23}\varepsilon_{e\tau})e^{i\delta}$$

$$\Xi \equiv \left(c_{12}s_{12} \frac{\Delta m_{21}^2}{a} + c_{23}\varepsilon_{e\mu} - s_{23}\varepsilon_{e\tau} \right) e^{i\delta}$$

“ θ_{12} ”

“ θ_{13} ”

$$\tilde{H}^{\text{NSI}} = U_{23}^\dagger \begin{bmatrix} \varepsilon_{ee} & \varepsilon_{e\mu} & \varepsilon_{e\tau} \\ \varepsilon_{e\mu}^* & \varepsilon_{\mu\mu} & \varepsilon_{\mu\tau} \\ \varepsilon_{e\tau}^* & \varepsilon_{\mu\tau}^* & \varepsilon_{\tau\tau} \end{bmatrix} U_{23} \equiv \begin{bmatrix} \tilde{\varepsilon}_{ee} & \tilde{\varepsilon}_{e\mu} & \tilde{\varepsilon}_{e\tau} \\ \tilde{\varepsilon}_{e\mu}^* & \tilde{\varepsilon}_{\mu\mu} & \tilde{\varepsilon}_{\mu\tau} \\ \tilde{\varepsilon}_{e\tau}^* & \tilde{\varepsilon}_{\mu\tau}^* & \tilde{\varepsilon}_{\tau\tau} \end{bmatrix}$$

$$= \begin{bmatrix} \varepsilon_{ee} & c_{23}\varepsilon_{e\mu} - s_{23}\varepsilon_{e\tau} & s_{23}\varepsilon_{e\mu} + c_{23}\varepsilon_{e\tau} \\ c_{23}\varepsilon_{e\mu}^* - s_{23}\varepsilon_{e\tau}^* & c_{23}^2\varepsilon_{\mu\mu} + s_{23}^2\varepsilon_{\tau\tau} - c_{23}s_{23}(\varepsilon_{\mu\tau} + \varepsilon_{\mu\tau}^*) & c_{23}^2\varepsilon_{\mu\tau} - s_{23}^2\varepsilon_{\mu\tau}^* + c_{23}s_{23}(\varepsilon_{\mu\mu} - \varepsilon_{\tau\tau}) \\ s_{23}\varepsilon_{e\mu}^* + c_{23}\varepsilon_{e\tau}^* & c_{23}^2\varepsilon_{\mu\tau}^* - s_{23}^2\varepsilon_{\mu\tau} + c_{23}s_{23}(\varepsilon_{\mu\mu} - \varepsilon_{\tau\tau}) & s_{23}^2\varepsilon_{\mu\mu} + c_{23}^2\varepsilon_{\tau\tau} + c_{23}s_{23}(\varepsilon_{\mu\tau} + \varepsilon_{\mu\tau}^*) \end{bmatrix}$$

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How big is the contribution of $\epsilon_{\alpha\beta}$?

Bird-eye view

Channel	ϵ_{ee}	$\epsilon_{e\mu}$	$\epsilon_{e\tau}$	$\epsilon_{\mu\tau}$	$\epsilon_{\mu\mu}$	$\epsilon_{\tau\tau}$	a dep. (NSI)	a dep. (SI)
$P(\nu_e \rightarrow \nu_\alpha): \alpha = e, \mu, \tau$	ϵ^3	ϵ^2	ϵ^2	ϵ^3	ϵ^3	ϵ^3	ϵ^2	ϵ^2
$P(\nu_\alpha \rightarrow \nu_\beta): \alpha, \beta = \mu, \tau$	ϵ^3	ϵ^2	ϵ^2	ϵ^1	$\epsilon^1(\epsilon^2)$	$\epsilon^1(\epsilon^2)$	ϵ^1	ϵ^2

Matter hesitation


Direct transition by NSI If θ_{23} is maximal

- decoupling of $\epsilon_{\alpha\beta}$ in μ - τ sector to $P(\nu_e \rightarrow \nu_\mu)$ and $P(\nu_e \rightarrow \nu_\tau)$ to ϵ^2
- Simplify NSI measurement strategy

➡ Determine $\epsilon_{e\mu}$ and $\epsilon_{e\tau}$ from $P(\nu_e \rightarrow \nu_\mu)$

➡ then, one can go to μ - τ sector

However, in μ - τ sector,

- $\varepsilon_{\alpha\beta}$ (in μ - τ sector) dependent term is common to all $P(\nu_\mu \rightarrow \nu_\tau)$, $P(\nu_\mu \rightarrow \nu_\mu)$ and $P(\nu_\tau \rightarrow \nu_\tau)$  impossible to determine 3 parameters, $\varepsilon_{\mu\tau} = |\varepsilon_{\mu\tau}| e^{i\varphi}$ and $(\varepsilon_{\mu\mu} - \varepsilon_{\tau\tau})$, by rate only analysis

Symmetry in $P(\nu_\alpha \rightarrow \nu_\beta)$ with NSI !

- $P(\nu_e \rightarrow \nu_\tau)$ can be obtained from $P(\nu_e \rightarrow \nu_\mu)$ by transformation

$$C_{23} \rightarrow -S_{23}, \quad S_{23} \rightarrow C_{23}$$

- and undoing any transformation in the generalized solar and atmospheric variables



Parameter degeneracy

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Parameter degeneracy 1

- Intrinsic, sign- Δm^2 , θ_{23} octant degeneracy prevails in a generalized form which include NSI parameters
- In matter perturbative regime, for example, $P_{e\mu}$ is approx. invariant under

$$\begin{aligned}\Delta m_{31}^2 &\rightarrow -\Delta m_{31}^2, \\ \delta &\rightarrow \pi - \delta, \\ \phi_{e\alpha} &\rightarrow 2\pi - \phi_{e\alpha}.\end{aligned}$$

← sign- Δm^2


SI version: HM-Nunokawa 01

$$\begin{aligned}c_{23} &\rightarrow s_{23}, \\ s_{23} &\rightarrow c_{23}, \\ (\varepsilon_{\mu\mu} - \varepsilon_{\tau\tau}) &\rightarrow -(\varepsilon_{\mu\mu} - \varepsilon_{\tau\tau}).\end{aligned}$$

← θ_{23} octant

Fogli-Lisi 96


Parameter degeneracy 2

- A completely new type of degeneracy also exist  solar – atmospheric variable exchange degeneracy

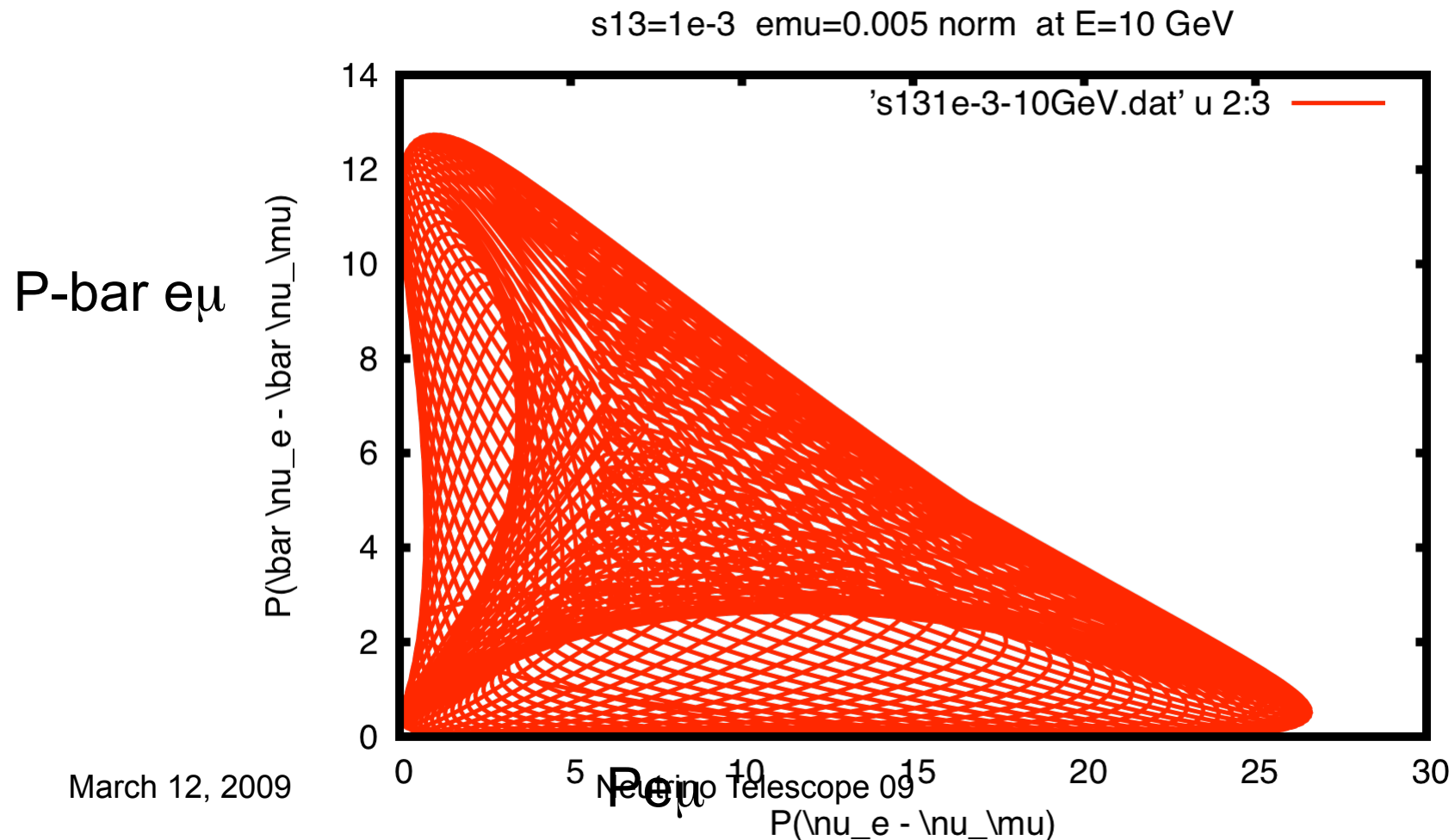
$$|\Theta_{\pm}^{(2)}| = \sqrt{\frac{Z}{X_{\pm}}} |\Xi^{(1)}| \quad \text{and} \quad |\Xi^{(2)}| = \sqrt{\frac{X_{\pm}}{Z}} |\Theta_{\pm}^{(1)}|$$

$$\begin{aligned} X_{\pm} &\equiv \left(\frac{a}{\delta m_{31}^2 \mp a} \right)^2 \sin^2 \frac{\delta m_{31}^2 \mp a}{4E} L, \\ Y_{\pm} &\equiv \left(\frac{a}{\delta m_{31}^2 \mp a} \right) \sin \frac{aL}{4E} \sin \frac{\delta m_{31}^2 \mp a}{4E} L, \\ Z &\equiv \sin^2 \frac{aL}{4E}. \end{aligned}$$

NSI perturbation theory; temporary summary

- Structure of ν oscillation with NSI is in fact very simple within ε perturbation theory
- But in real life, it is complicated even with single $\varepsilon_{\alpha\beta}$ because the system is enriched with 2 phases  lepton KM phase + NSI phase (see next page)

ν oscillation with NSI is already complicated phenomenon even with single $\varepsilon_{\alpha\beta}$ (example with $\varepsilon_{e\mu}$)





Solving the 2 phase confusion with 2- detector setting

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SI-NSI confusion

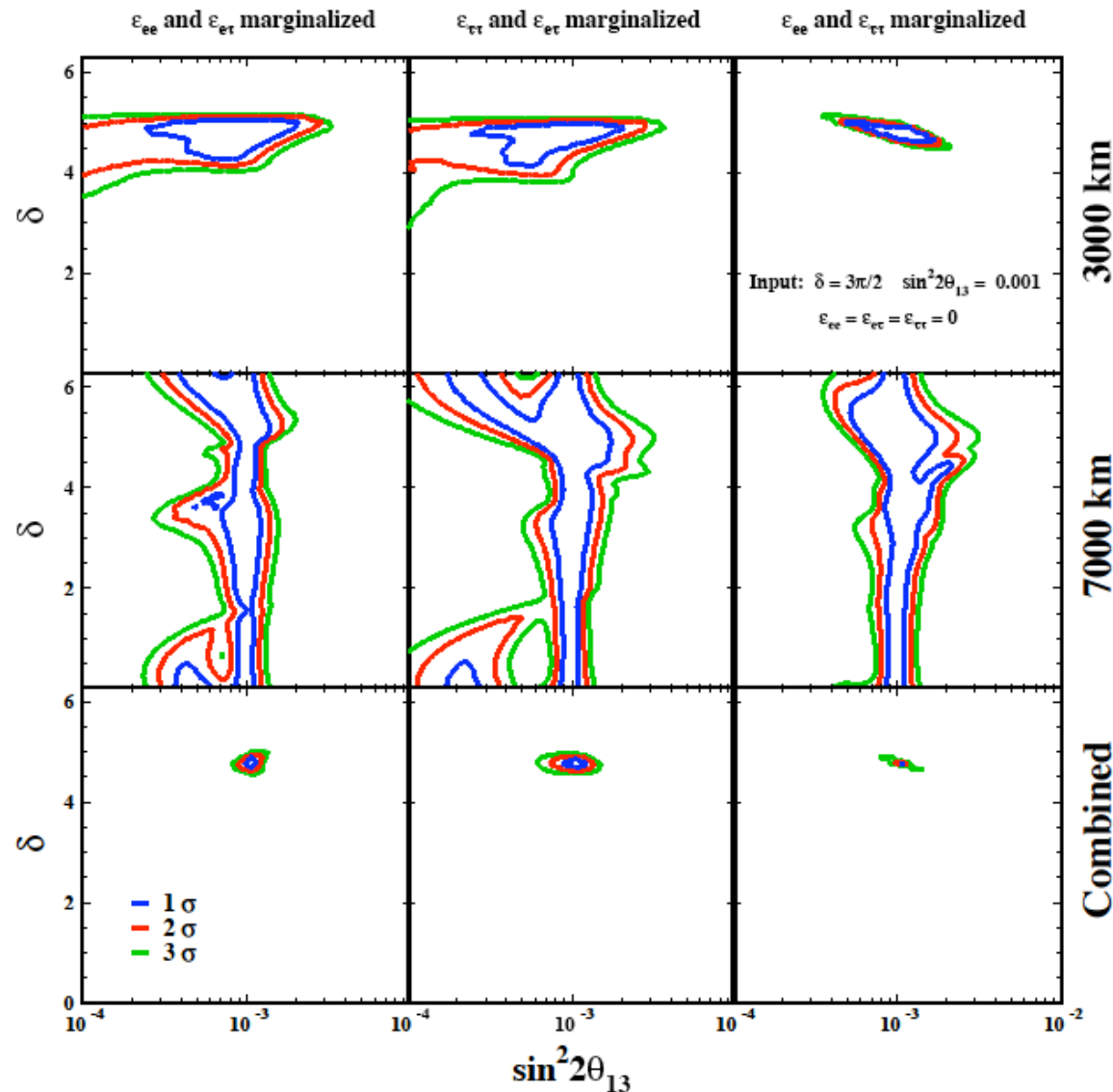
- We saw that SI and NSI parameters come together:  **SI-NSI confusion**

$$\Theta_{\pm} \equiv s_{13} \frac{\Delta m_{31}^2}{a} + (s_{23} \varepsilon_{e\mu} + c_{23} \varepsilon_{e\tau}) e^{i\delta}$$

$$\Xi \equiv \left(c_{12} s_{12} \frac{\Delta m_{21}^2}{a} + c_{23} \varepsilon_{e\mu} - s_{23} \varepsilon_{e\tau} \right) e^{i\delta}$$

- Last year I have reported that 2 detector setting (3000 + 7000 km) in neutrino factory can resolve θ_{13} - $|\varepsilon_{e\mu}|$ (or $|\varepsilon_{e\tau}|$) confusion
- Today, I discuss 2 phase confusion; **CPV** phase of $\varepsilon_{\alpha\beta}$ can be confused with lepton KM phase δ  single $\varepsilon_{\alpha\beta}$ system

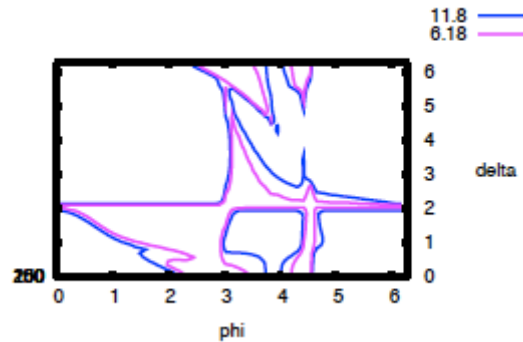
θ_{13} - NSI confusion can be solved: @NO-VE08



Set up: 10^{21} muons/year, 4+4 years of ν , $\bar{\nu}$, two 50 kt iron detectors (Cipriani Ribeiro-HM-Nunokawa-Uchinami-Zukanovich Funchal, Arxiv: 0709.1980=>JHEP 07)

2 phase confusion can be resolved

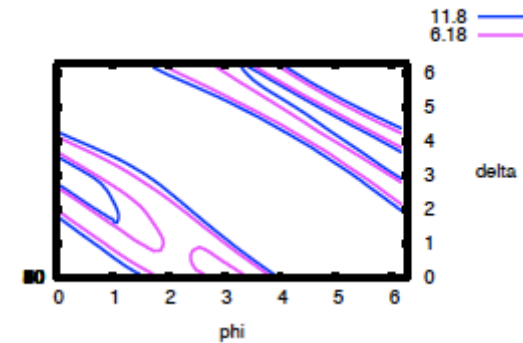
$\epsilon_{13}=1e-4$ $\delta\epsilon = 1\text{piby}2$ $\phi = 3\text{piby}4$ $\epsilon\tau = 5.e-3 - 3000$ km



3000 km

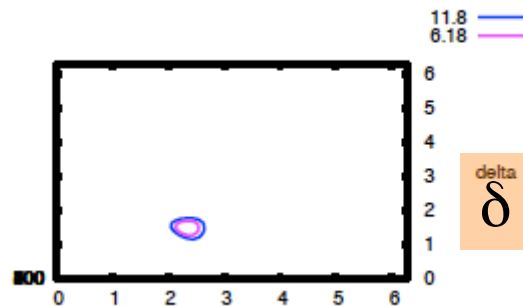
$$|\epsilon_{e\tau}| = 5 \times 10^{-3}$$

$\epsilon_{13}=1e-4$ $\delta\epsilon = 1\text{piby}2$ $\phi = 3\text{piby}4$ $\epsilon\tau = 5.e-3 - 7000$ km



7000 km

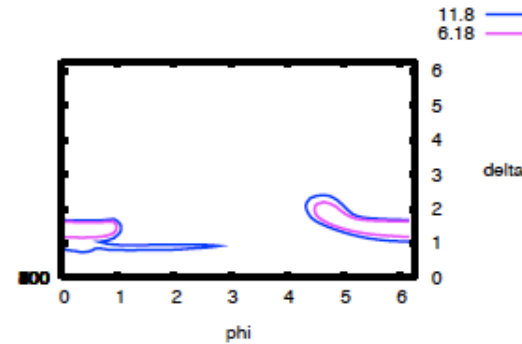
$\epsilon_{13}=1e-4$ $\delta\epsilon = 1\text{piby}2$ $\phi = 3\text{piby}4$ $\epsilon\tau = 5.e-3 - \text{combined}$



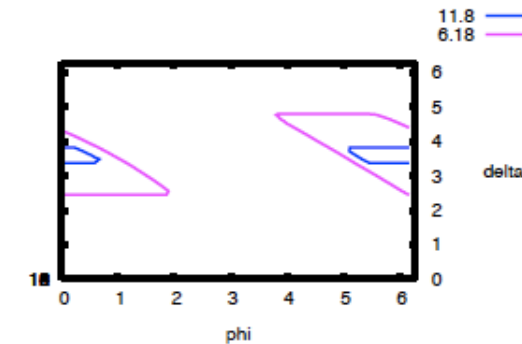
combined

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 ϕ

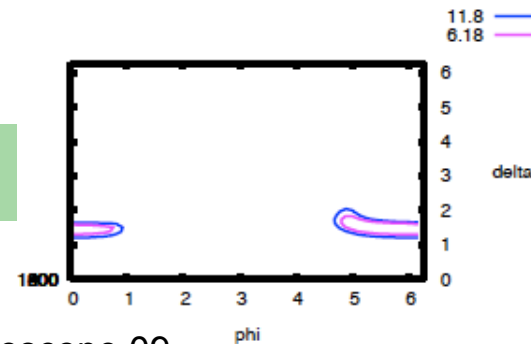
$\epsilon_{13}=1e-4$ $\delta\epsilon = 1\text{piby}2$ $\phi = 7\text{piby}4$ $\epsilon\mu = 5.e-4 - 3000$ km



$\epsilon_{13}=1e-4$ $\delta\epsilon = 1\text{piby}2$ $\phi = 7\text{piby}4$ $\epsilon\mu = 5.e-4 - 7000$ km



$\epsilon_{13}=1e-4$ $\delta\epsilon = 1\text{piby}2$ $\phi = 7\text{piby}4$ $\epsilon\mu = 5.e-4 - \text{combined}$



$$|\epsilon_{e\mu}| = 5 \times 10^{-4}$$

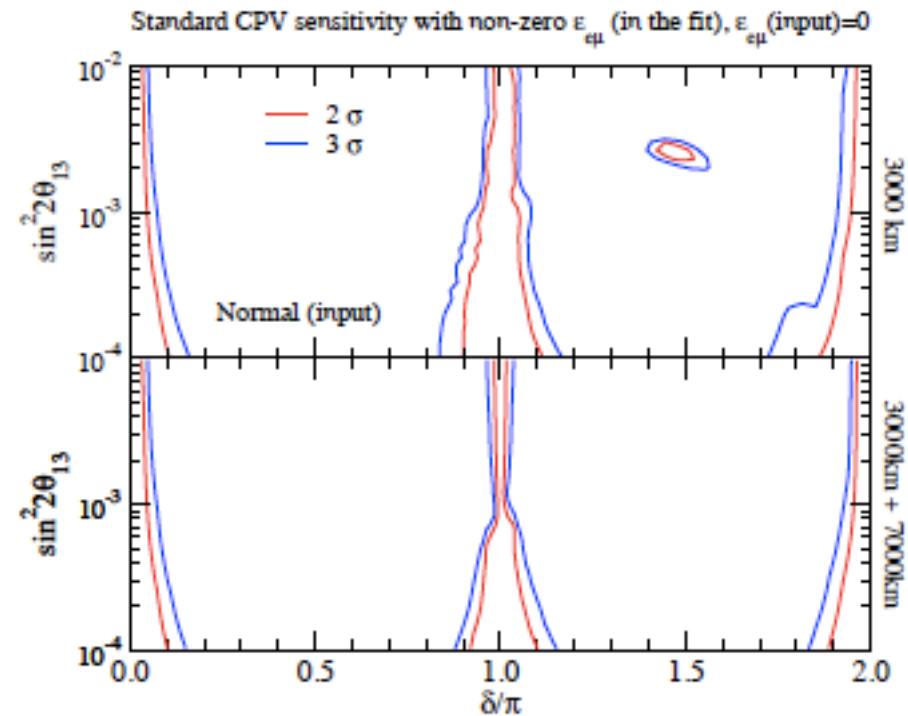
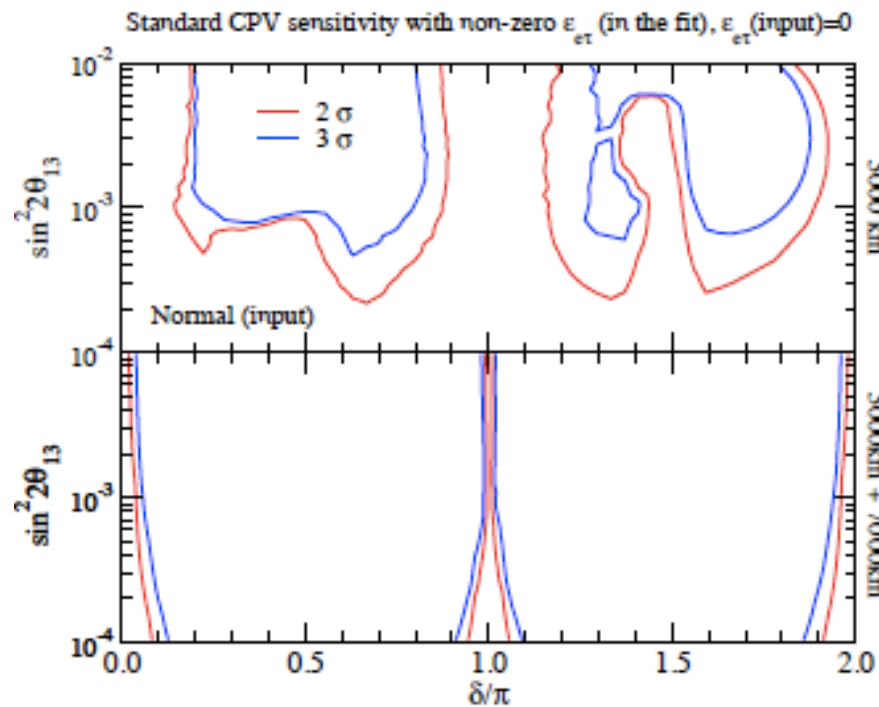
Cipriani
Ribeiro
et al.
coming
soon !

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Lepton KM CPV discovery potential

$$\varepsilon_{e\tau}$$

$$\varepsilon_{e\mu}$$

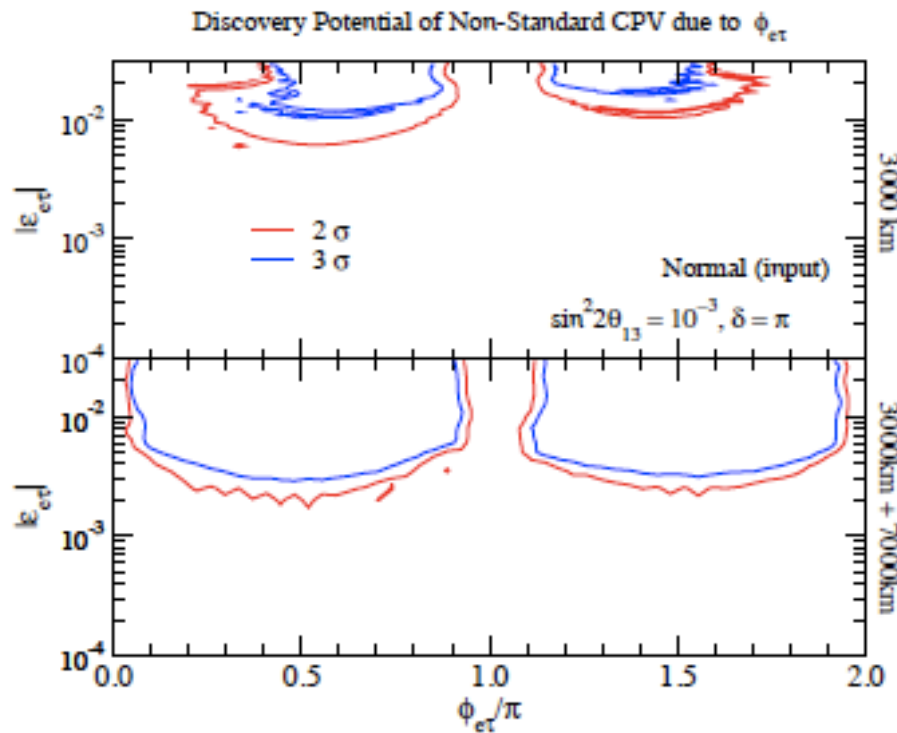


Upper 3000 km
lower combined

θ_{13} , $\varepsilon_{e\mu}$, etc.
marginalized

NSI phase CPV discovery potential

$$\varepsilon_{e\tau}$$

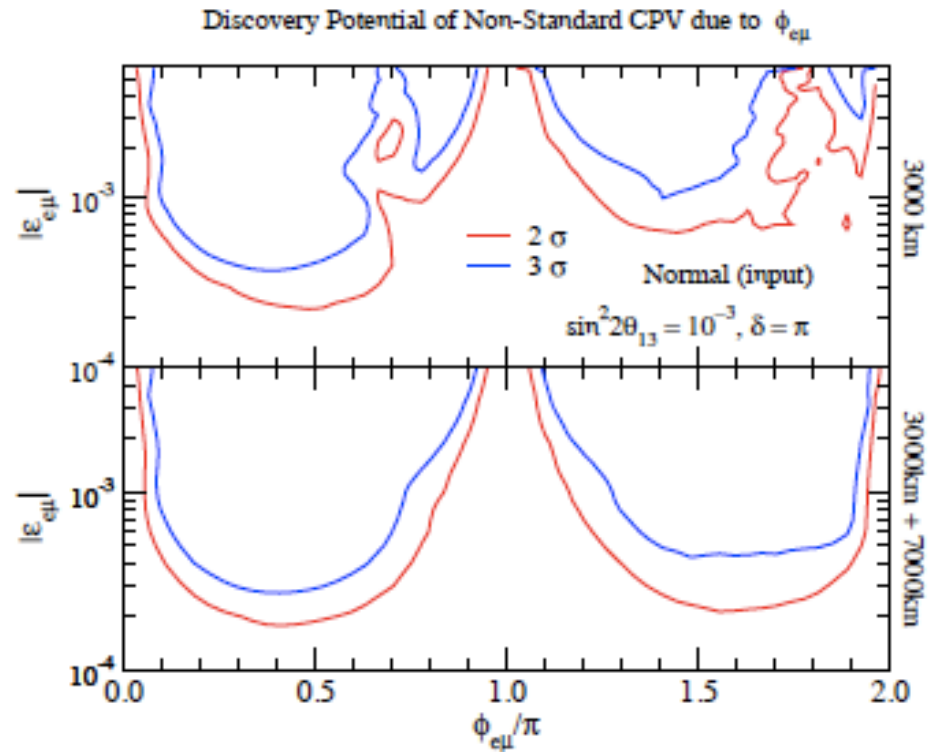


Upper 3000 km
lower combined

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$$\varepsilon_{e\mu}$$



θ_{13}, δ , etc.
marginalized

Conclusion

- New consideration lead to loosing the bound on NSI by a factor of $\sim 10^{-3}$
- Global overview for ν oscillation with NSI is obtained by perturbative formulation
- problem of 2 phase confusion addressed with 2 detector setting 3000 + 7000 km in neutrino factory

 synergy between 2 detectors seems strong enough to solve it with single $\varepsilon_{\alpha\beta}$

 Can establish 2 new CPV!

Second loach can be BIG!



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