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SM Extensions with Gauged B-L

NO-VE, 17 April 2008

Mostly a critical review
after discussions with:

- L.Basso, M.Passera
- G.Villadoro



Plan

- Motivations
- Constraints on extra Z' bosons
- A minimal non-susy model
- A minimal susy model
- Kinetic mixing
- Conclusions and outlook

Choose here $M_R \sim \text{TeV}$ and $y_N \sim y_e$

Hierarchy problem? Today's energy frontier!

Alone, no big impact at high-energy colliders:

$y_N \sim y_e \rightarrow \text{SM \& nu-R very weakly coupled}$

How to generate a more interesting coupling?

$M_R = \text{the scale of (B -) L breaking}$

Promote (B-L) to an extra U(1) gauge symmetry

A reason for right-handed neutrinos:

$Y' = a Y + b \text{ (B-L)}$ automatically anomaly-free
if fermions in SM families with nu-R

Weinberg, QFT-II, p.388: "a neutral vector boson somewhat heavier than the Z^0 and coupled to B-L seems like the most plausible addition to the SM"

Pragmatic motivation: “easy” LHC signal ?

$Z' \rightarrow e^+e^-$ with SM-like couplings (Z_{SSM})

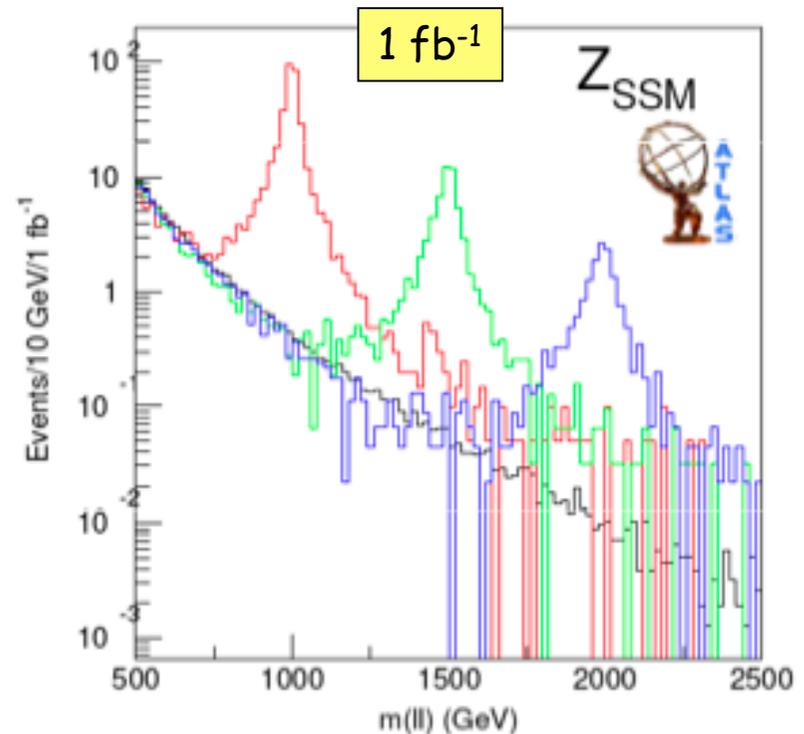
Mass	Expected events for 1 fb ⁻¹ (after all analysis cuts)	Integrated luminosity needed for discovery (corresponds to 10 observed evts)
1 TeV	~ 160	~ 70 pb ⁻¹
1.5 TeV	~ 30	~ 300 pb ⁻¹
2 TeV	~ 7	~ 1.5 fb ⁻¹

- with 100 pb⁻¹ large enough signal for discovery up to $m > 1$ TeV
- signal is (narrow) mass peak on top of small Drell Yan background
- ultimate calorimeter performance not needed

Ultimate ATLAS reach (300 fb⁻¹): ~ 5 TeV

[F. Gianotti, CERN-SPC, 17/9/07]

Similar reach for the CMS experiment



Further theoretical motivations

GUTs

Embeddable in **SO(10) grand unification**:

$$\begin{aligned} \text{SO}(10) &\rightarrow \text{SU}(3)_C \times \text{SU}(2)_L \times \text{SU}(2)_R \times \text{U}(1)_{B-L} \\ &\rightarrow \text{SU}(3)_C \times \text{SU}(2)_L \times \text{U}(1)_Y \times \text{U}(1)_{Y'} \end{aligned}$$

e.g., with a Higgs in the adjoint 45 representation

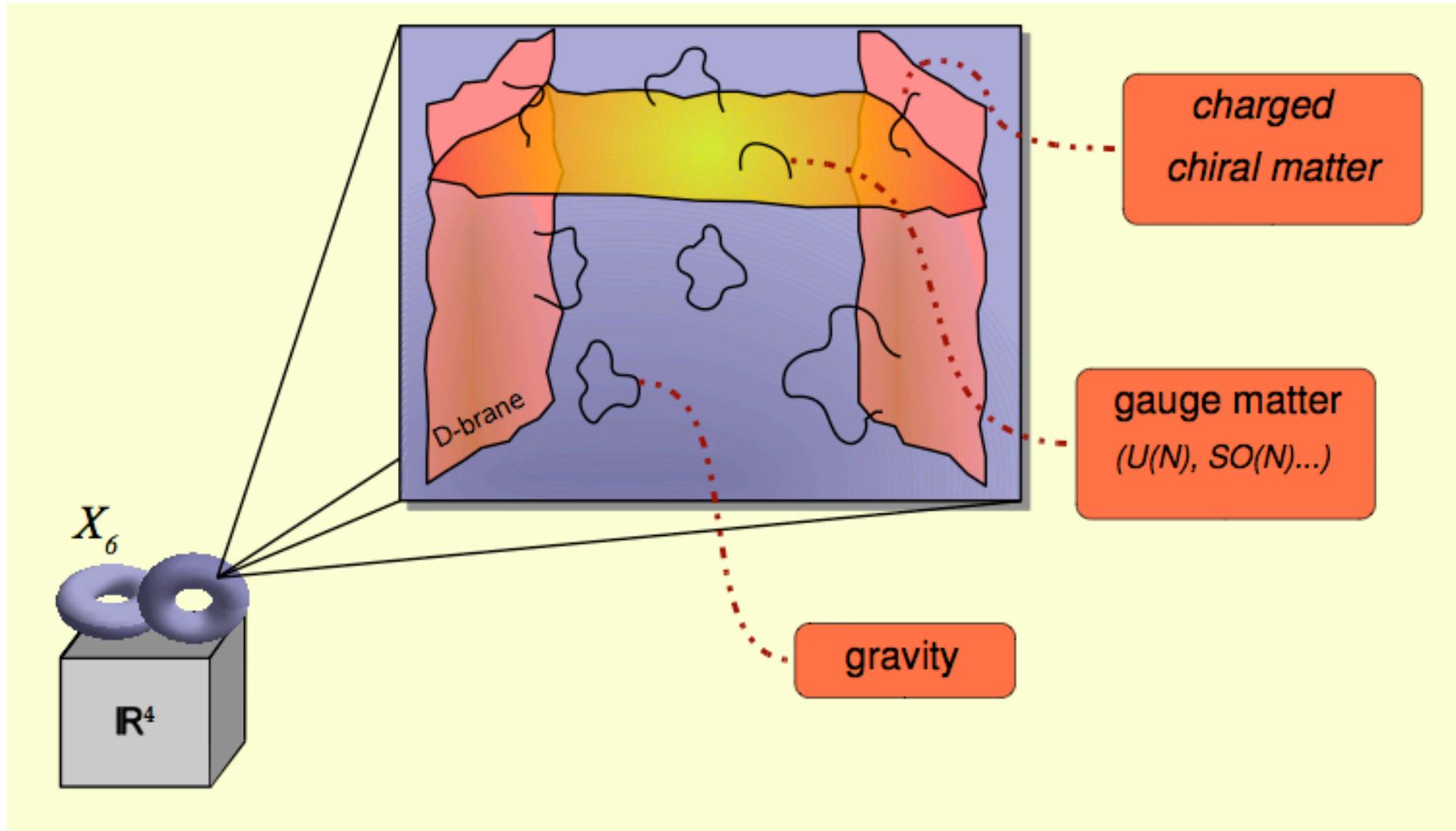
Type-II string models with D-branes

Gauge group for a stack of **N parallel D-branes**:

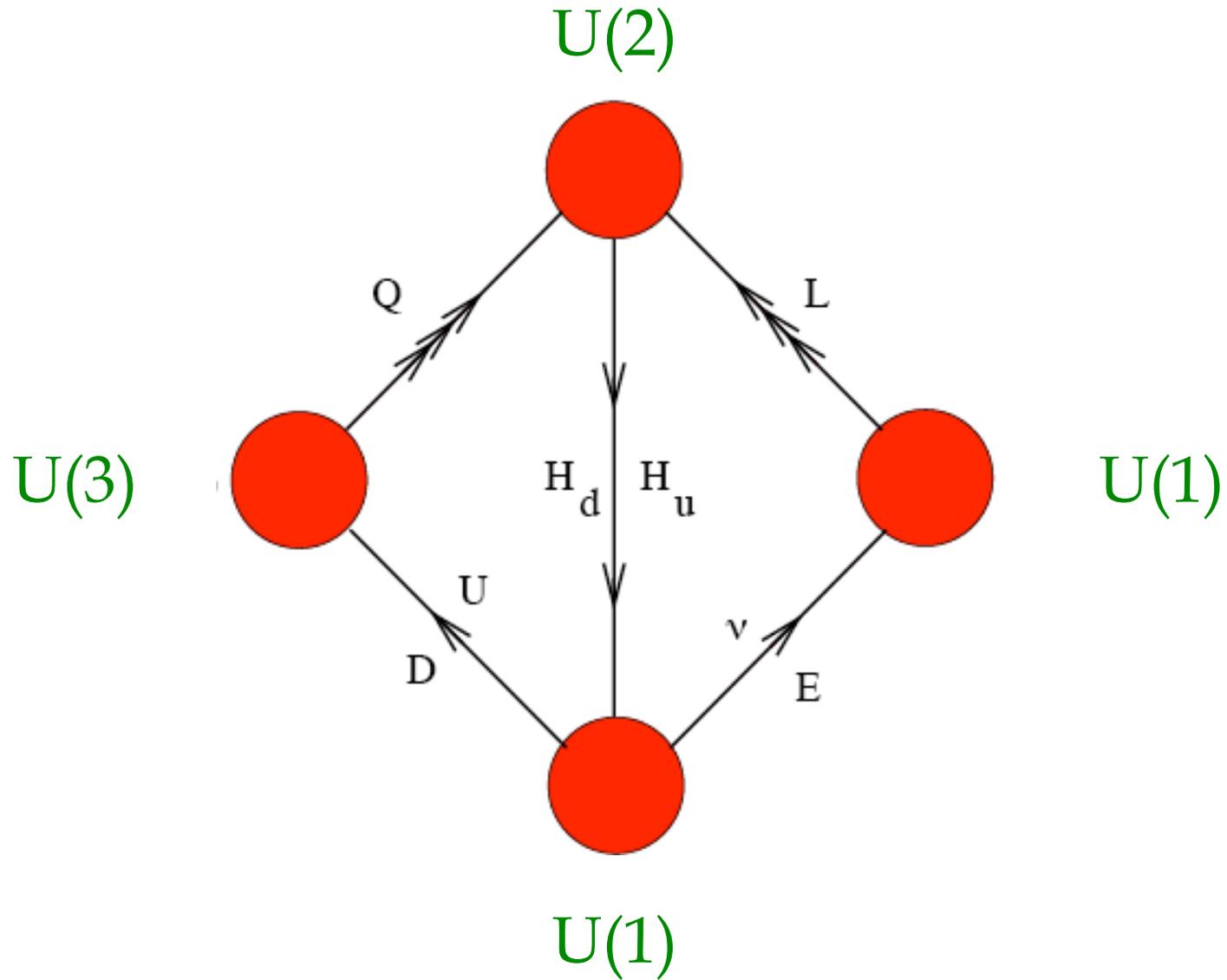
$$\text{U}(N) \rightarrow \text{SU}(N) \times \text{U}(1)$$

Multiple U(1) factors frequent in realistic models often including a residual non-anomalous $\text{U}(1)_{B-L}$

A picture of the brane-world (IIA)



An example of SM-like “quiver”



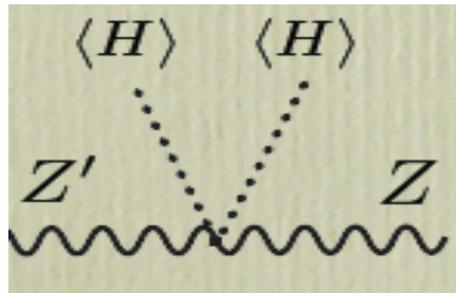
LEP bounds on Z'

$$\mathcal{L} = \mathcal{L}_{SM} - \frac{1}{4} Z'_{\mu\nu} Z'^{\mu\nu} - \frac{1}{2} M_{Z'}^2 Z'_\mu Z'^\mu \quad (\text{diagonal kin. and mass term})$$

$$+ g_{Z'} Z'_\mu \sum_f \bar{f} z_f \gamma^\mu f + g_{Z'} (H^\dagger z_H Z'_\mu i D^\mu H + h.c.) + \dots$$

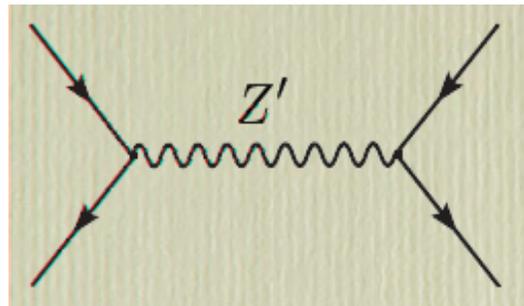
(couplings to SM fermions) (Z-Z' mixing after EWSB)

LEP-1 Z-pole data
mostly constrain
Z-Z' mixing
 $|\theta| < \mathcal{O}(10^{-3})$



$$\theta \sim \frac{g_{Z'}}{g_Z} \frac{M_Z^2}{M_{Z'}^2} z_H$$

LEP-2 (off-pole) data
constrain **4-fermion**
effective **operators**



$$\sim \frac{g_{Z'}^2}{M_{Z'}^2} z_e z_f$$

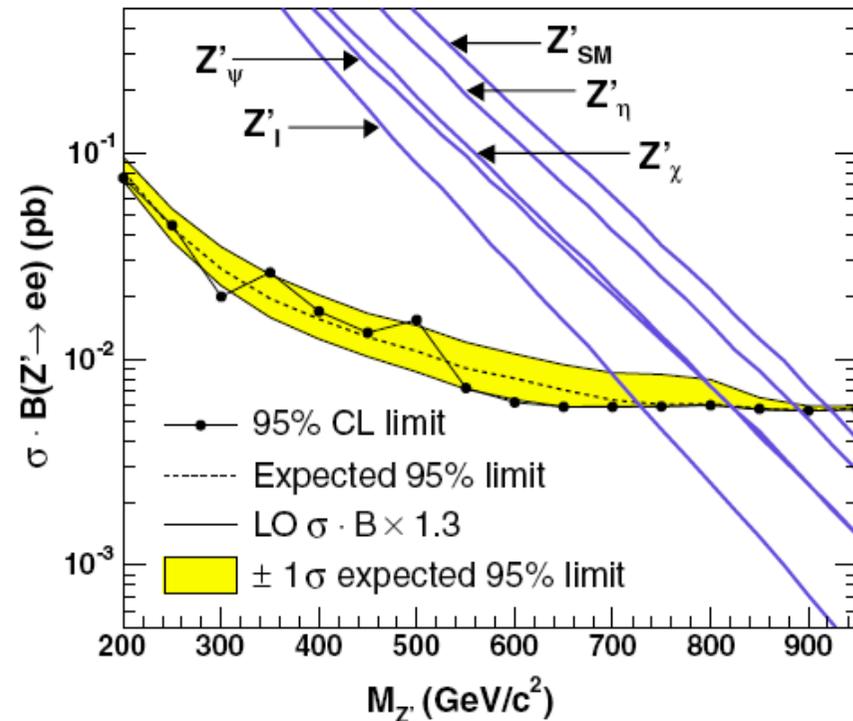
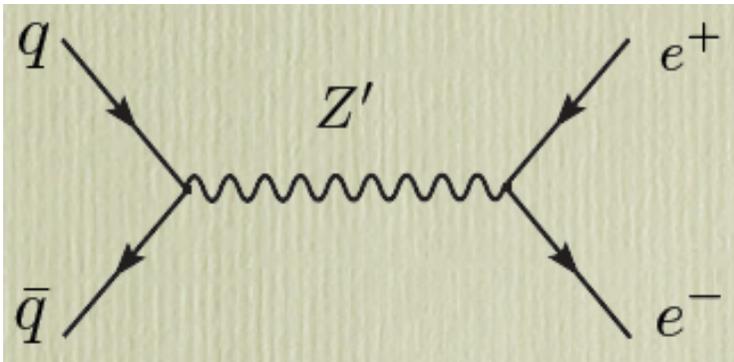
Tevatron bounds on Z'

More difficult to parametrize in a simple way!

Typical bounds are on $\sigma(Z') \cdot BR(Z' \rightarrow l^+l^-)$

But (already at leading order):

$$\sigma(Z') = g_{Z'}^2 f(z_q, z_u, z_d, s, M_{Z'}^2)$$



where f depends on the PDF

CDF Run II 1.3 fb^{-1}
PRL 99 (2007) 171802

Quantum numbers of the SM particles

	SU(3)	SU(2)	Y	$B - L$	T_{3R}	X
Q	3	2	1/6	1/3	0	-1
D^c	$\bar{3}$	1	1/3	-1/3	1/2	3
U^c	$\bar{3}$	1	-2/3	-1/3	-1/2	-1
L	1	2	-1/2	-1	0	3
E^c	1	1	1	1	1/2	-1
N^c	1	1	0	1	-1/2	-5
$H = H_2$	1	2	1/2	0	1/2	2
$\chi = \chi_2$	1	1	0	2	-1	-10
H_1	1	2	-1/2	0	-1/2	-2
χ_1	1	1	0	-2	1	10

$$Q = T_{3L} + Y \quad Y = T_{3R} + \frac{B - L}{2} \quad X = 4Y - 5(B - L) = 4T_{3R} - 3(B - L)$$

LEP1 bounds more easily evaded for $Y' = B - L$ [$z_H = 0$]

LEP2 bounds on $M_{Z'}(\text{TeV}) / g_{Z'}$: 6-7 for B-L, 15 for X
(most favourable case would be "leptophobic" Z')

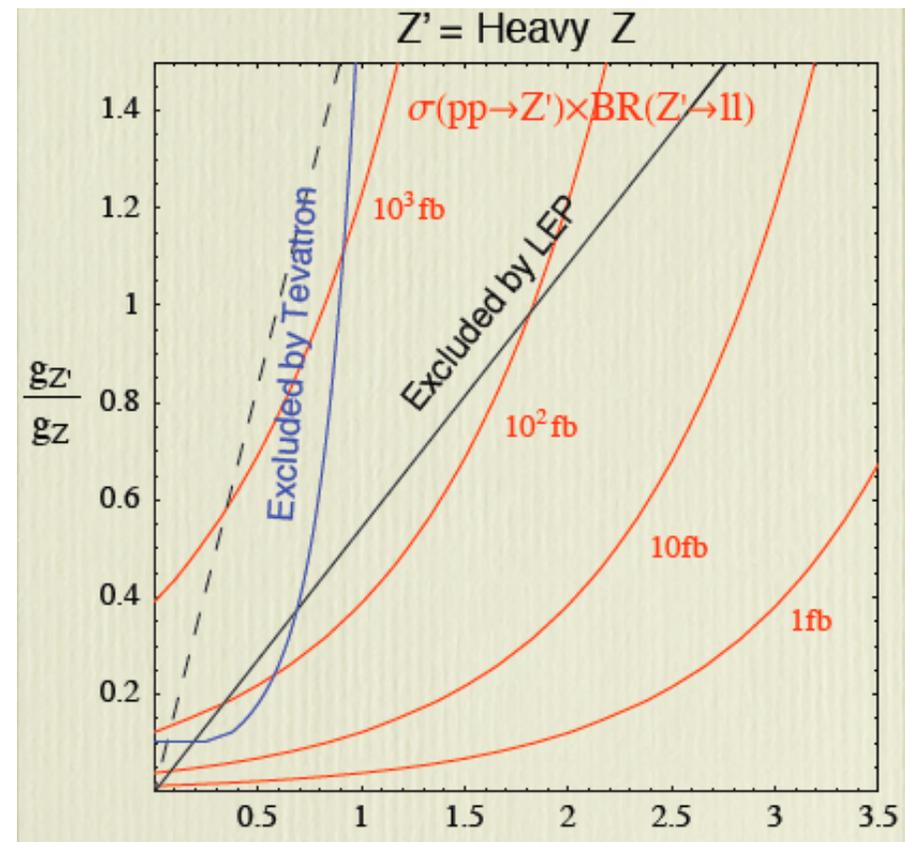
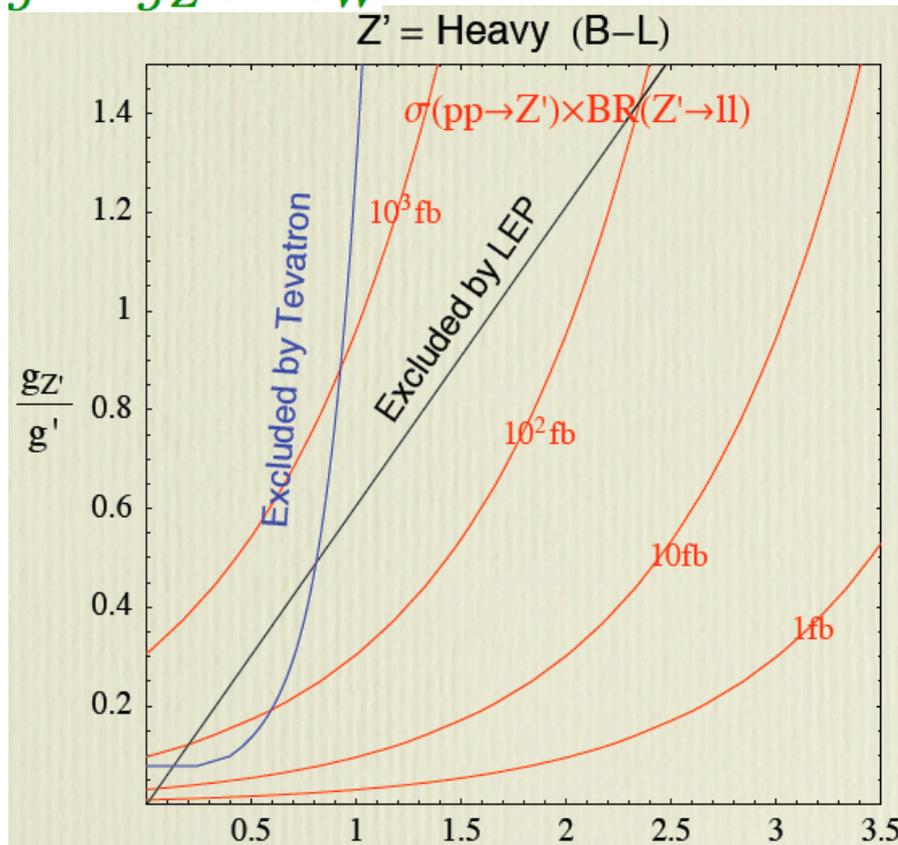
How much room left for the LHC ?

Normalization of $g_{Z'}$ and identity of Y' model-dependent

Direct SO(10) breaking would give $Y' \sim X$ and $g_{Z'}/g' \sim 0.2$

Different possibilities within brane-world constructions

$$g' = g_Z \sin \theta_W$$



[Contino, talk at the Perugia LHC workshop, 30/1-2/2/08] $M_{Z'} (TeV)$

A minimal non-SUSY model

[Buchmuller-Greub-Minkowski 1991]

SM (with 3 right-neutrinos & Higgs doublet) + extra $U(1)_{Y'}$
+ complex SM-singlet Higgs χ with $B-L=-2$ to generate M_R

$$\Delta\mathcal{L}_{Yuk} = y_R \nu_R \nu_R \chi + h.c.$$

$$V = m^2 |H|^2 + \mu^2 |\chi|^2 + \lambda_1 |H|^4 + \lambda_2 |\chi|^4 + \lambda_3 |H|^2 |\chi|^2$$

The gauged version of the singlet Majoron model

[Chikashige-Mohapatra-Peccei 1980]

Acceptable symmetry breaking for suitable parameter choices

Main phenomenological features:

[see, e.g., recent studies by Khalil et al, Basso et al]

- **Z' phenomenology** as discussed before [with possible decays into right-handed neutrinos, no Z-Z' tree-level mixing for canonical gauge kinetic terms & $Y' = B-L, \dots$]
- An **extended Higgs spectrum**: two neutral scalars h_1 & h_2 with complementary couplings to SM states controlled by their mixing angle. Typically, **weakened signals at the LHC**.
- The possibility of a **purely radiative symmetry-breaking** of the gauge symmetry via the **Coleman-Weinberg mechanism** (setting to zero the mass parameters in the scalar potential)

A minimal SUSY model

[Babu-Dutta-Mohapatra 2003, Khalil-Masiero 2007]

Enlarge the Higgs sector as required by supersymmetry:

$$H \rightarrow (H_1, H_2) \quad \chi \rightarrow (\chi_1, \chi_2)$$

Write general gauge-invariant renormalizable W :

$$W = W_{MSSM} + y_N L N^c H_2 + y_R N^c N^c \chi_1 + \mu' \chi_1 \chi_2$$

(automatically conserving baryon and lepton number)

After introducing **soft SUSY breaking** as usual:

- Can realize radiative breaking of gauge symmetry
- Link $SU(2) \times U(1)$, (B-L) and SUSY-breaking scales
- Richer spectrum of neutralinos & neutral Higgses
- An enlarged sneutrino sector within the TeV scale

Kinetic mixing

[Holdom 1986; DelAguila-Quiros-FZ 1987; ...]

In the presence of (at least) **two U(1) factors**, can write

$$\mathcal{L}_{kin} = -\frac{1}{4} (g^{-2})_{mn} F_{\mu\nu}^m F^{n\mu\nu}$$

with $(g^{-2})_{mn}$ defining a **matrix of coupling constants** besides the **two diagonal U(1) couplings**, the **third off-diagonal coupling g_x** can be reabsorbed into g_Z, Y' but this is **not stable** against quantum corrections:

$$\frac{d(g^{-2})_{mn}}{dt} = -\frac{b_{mn}}{8\pi^2} \quad b_{mn} = \frac{2}{3} \sum_f Q_m^f Q_n^f + \frac{1}{3} \sum_b Q_m^b Q_n^b$$

Only for **orthogonal U(1) generators** mixing postponed to 2-loop and threshold effects: almost true (excluding Higgs sector) for the (Y, X) or $(T_{3R}, B-L)$ pairs in the table

SUSY kinetic mixing

In the **supersymmetric** case, **gauge kinetic mixing** extends to **gaugino masses** and kinetic terms also to the **D-term** part of the scalar potential

$$\mathcal{L}_{kin} = -\frac{1}{4} \int d^2\theta h_{mn} W^m W^n + h.c.$$

[old SUGRA literature; Dienes-Kolda-Russell 1997]

The **MSSM RGE** can be fully generalized [Villadoro, FZ]

Consequence: minimal models discussed above cannot be extrapolated as such to **very large scales** (e.g. **GUT** scale) kinetic mixing effects must be properly included, e.g.:

$$\theta \sim \frac{g_x}{g_Z} \frac{m_Z^2}{M_{Z'}^2}$$

Conclusions and outlook

- SM extensions with **right-handed neutrinos** and an extra **U(1) gauging B-L** [or $Y' = aY + b(B-L)$] are quite plausible
- **LEP constraints** are quite strong (often more than Tevatron) but leave still **room for possible discoveries at the LHC**
- **Kinetic mixing** effects cannot be neglected in general especially when **extrapolating models to high scales**
- Interesting to explore more systematically **ranges of $g_{Z'}$ and Y' combinations allowed by brane-world models**
- Can one build a **natural** (not fine-tuned) SUSY model of this kind, compatible with grand unification, precision tests and cosmological constraints (baryogenesis,...) ?
If so, would be worth exploring in detail its predictions!