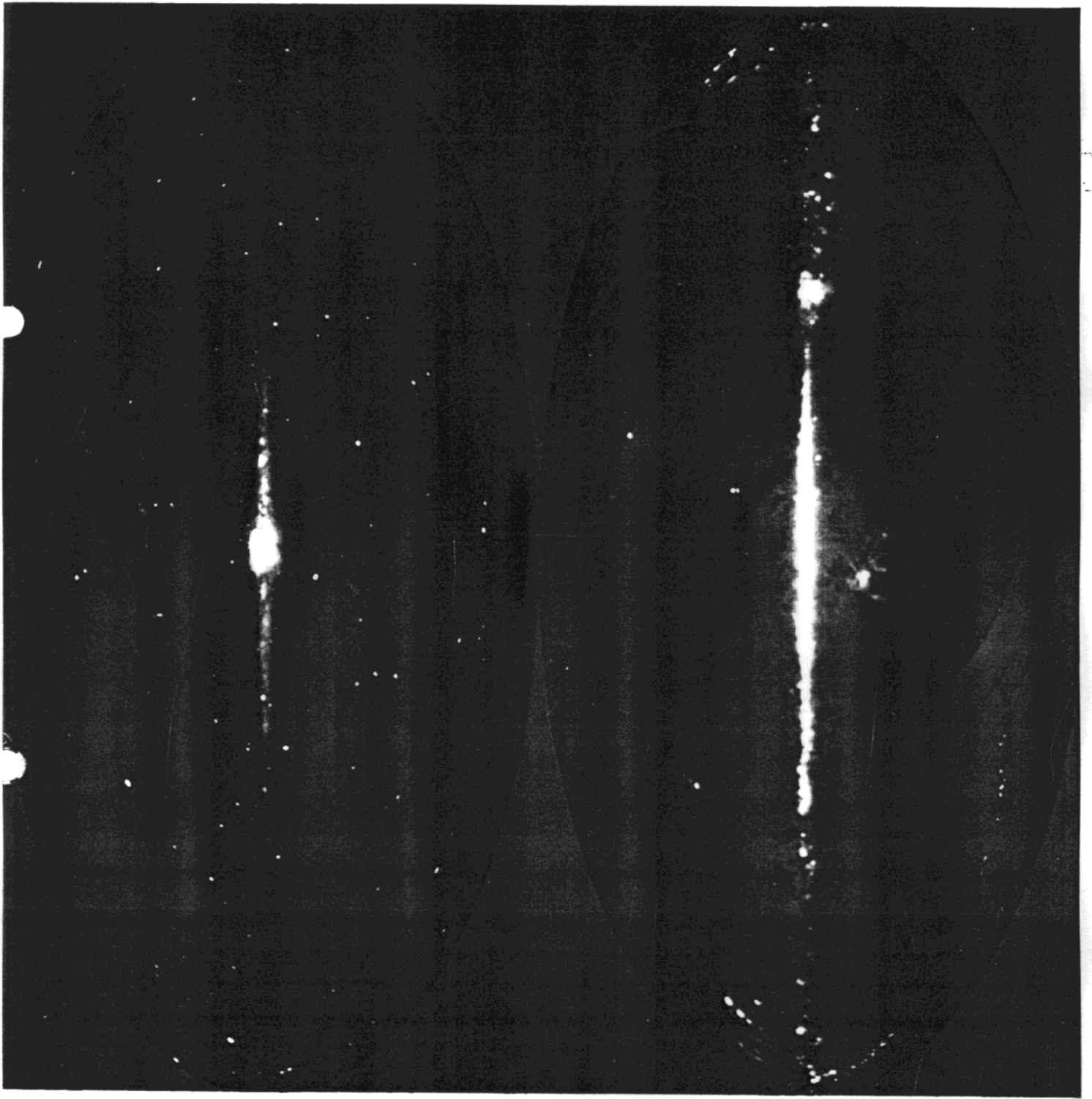


The Antiproton Signal of the Galactic Supersymmetric Dark Matter

Neutrino Telescopes - Venezia - February 25 1999

P. Salati, Laboratoire LAPTH and Université de Savoie

- 1) Antiproton production in our galaxy
- 2) BESS 95 & supersymmetric antiprotons
Bottino, Donato, Fornengo & Salati PR D58 (1998) 123503
- 3) Primary vs secondary antideutons
Donato, Fornengo & Salati
- 4) The instruments : balloon vs space
- 5) Discussion & future prospects

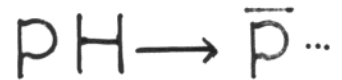


Antiproton production in our galaxy

Diffusion process :

$$\frac{\partial n}{\partial t} = 0 = K \Delta n + 2h\delta(z)q - 2h\delta(z)\Gamma n$$

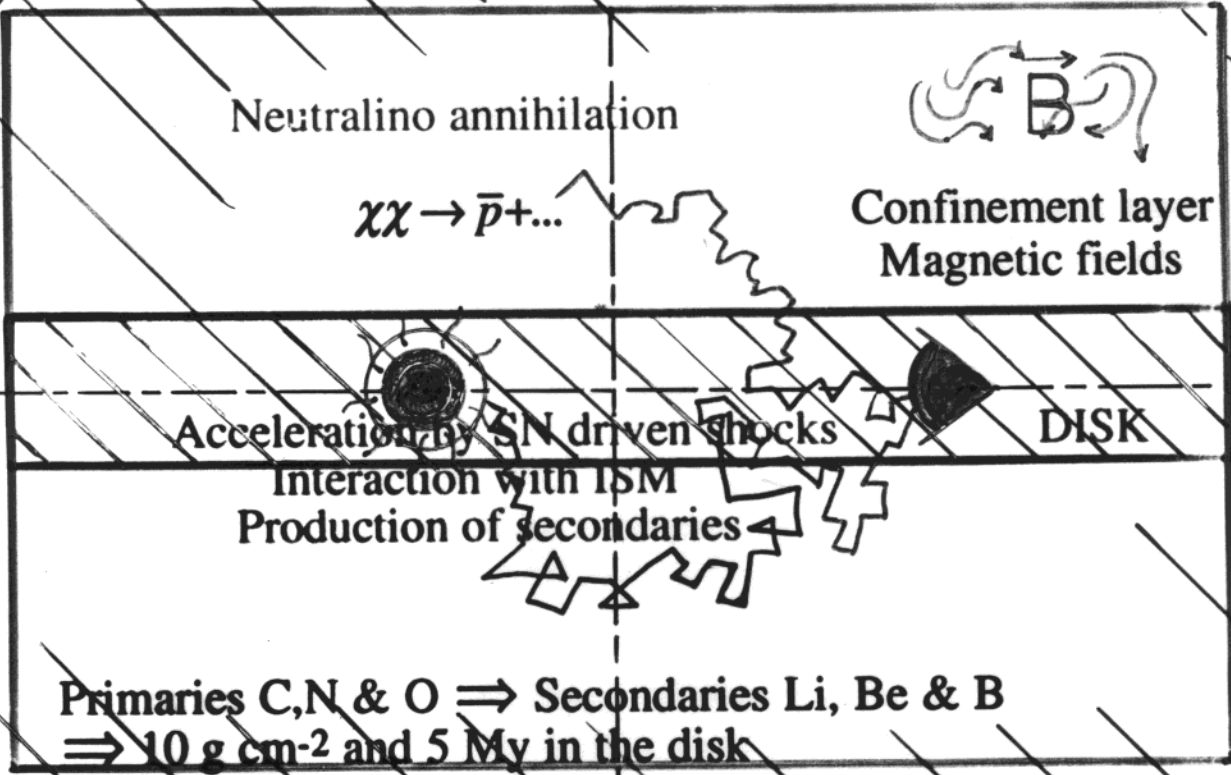
Steady state



$$\bar{p} \text{ source} = q = 4\pi n_H \int \frac{d\sigma}{dE_{\bar{p}}} \phi_p dE_p$$

Spallation $\Gamma = \sigma n_H v$

3 kpc
100 pc



^{10}Be (1.6 My) vs ^9Be \Rightarrow 50 My in the galaxy
 \Rightarrow Extended layers of confinement.

DM SUSY halo

$$K_0 = 6 \times 10^{27} \text{ cm}^2 \text{ s}^{-1}$$

$$K = K_0 \left\{ 1 + R \right\}^{0.6}$$

Ficenec et al

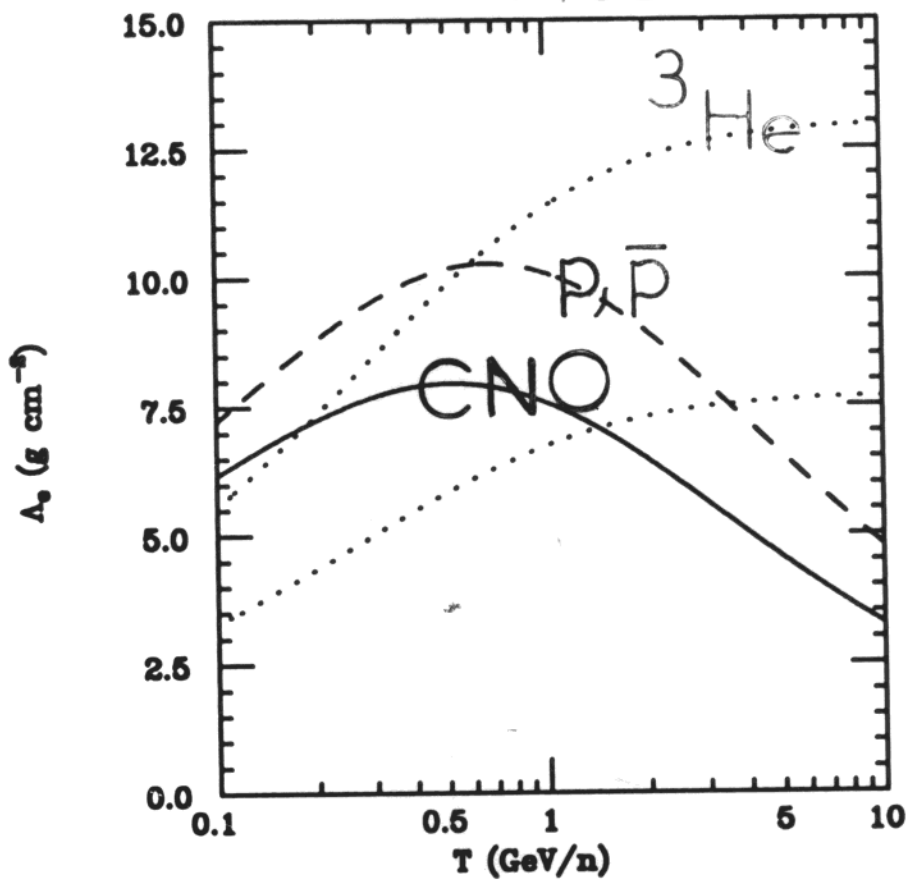


Figure 4

$$\chi^2_{red} = 0.83$$

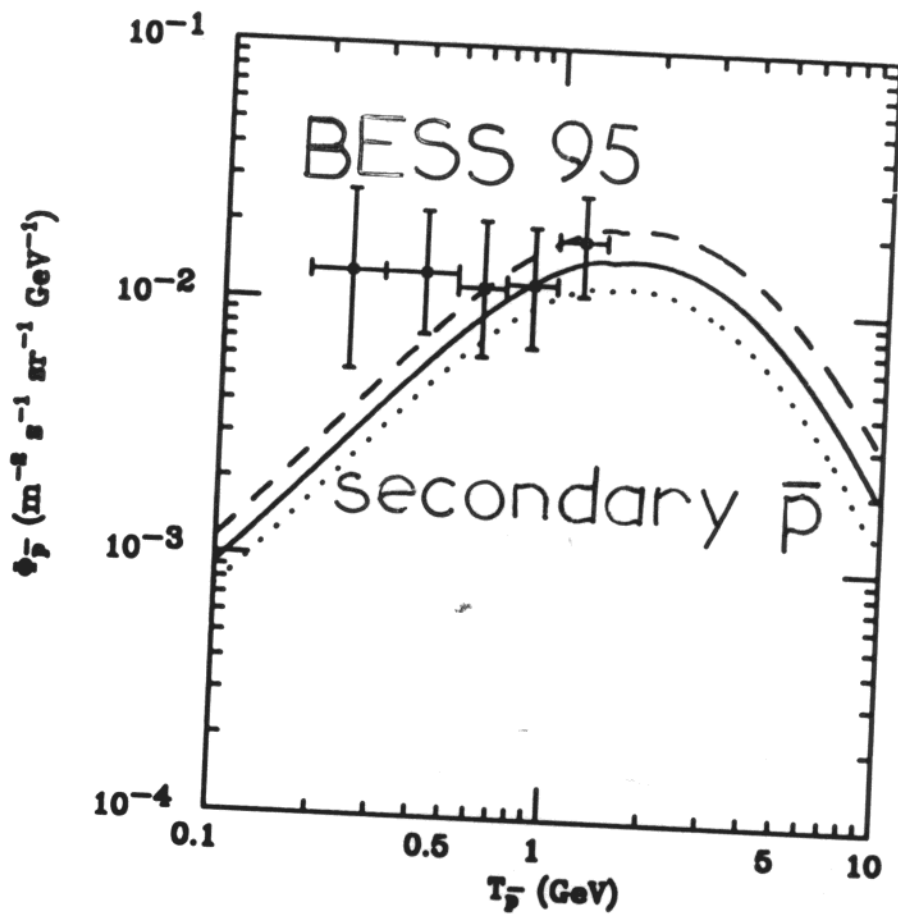


Figure 1

SUSY analysis

a. $\chi^2_{red} < 2.2$

b. $0.03 < \Omega h^2 < 0.7$
 χ

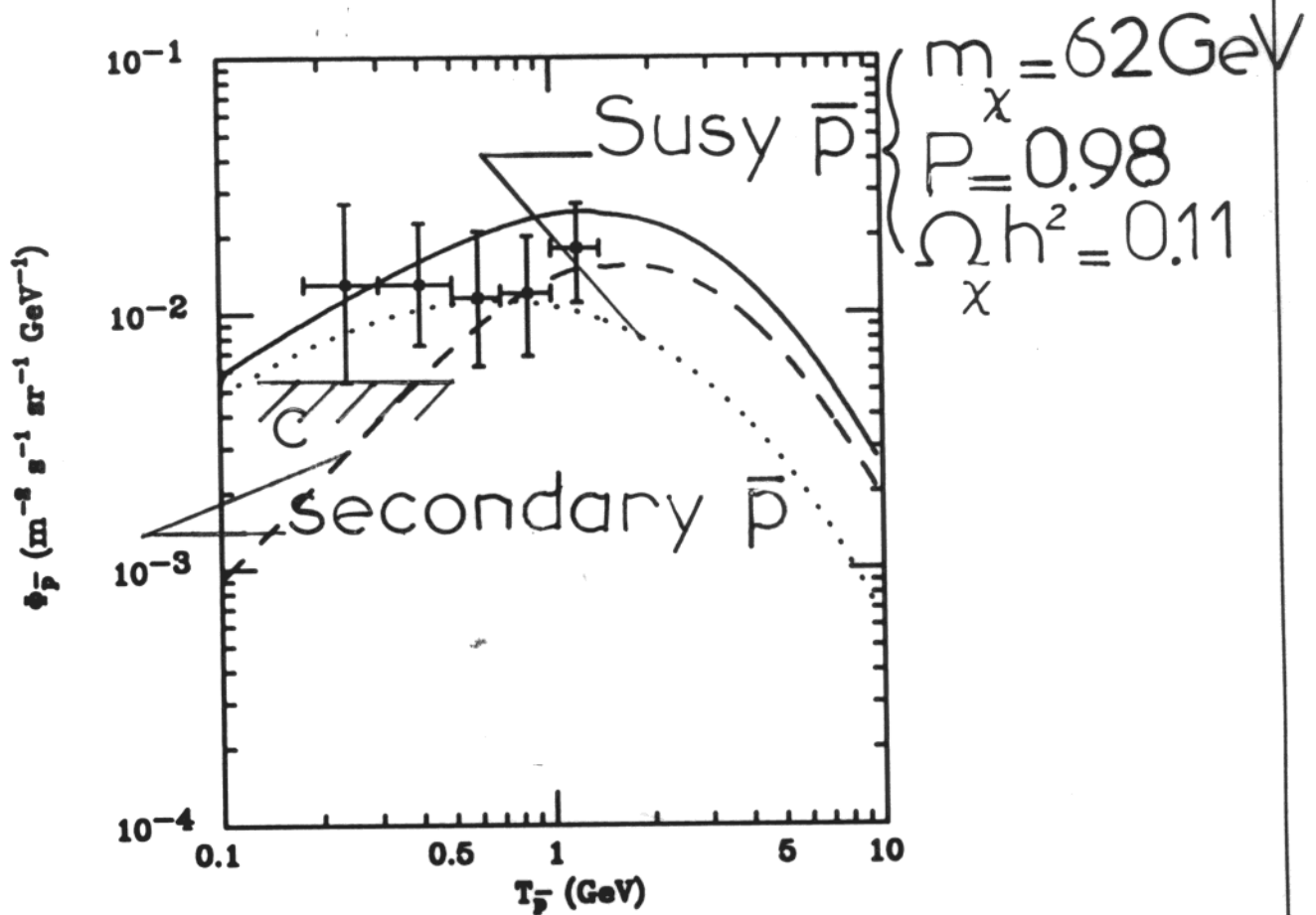


Figure 10

Compatible with BESS

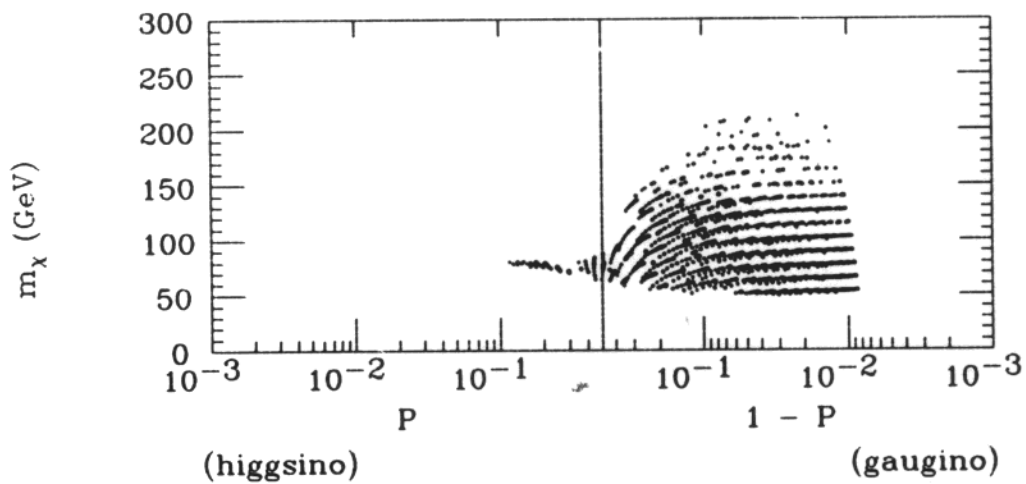


Figure 12 (b)

$$\chi_{red}^2 > 4$$

Excluded

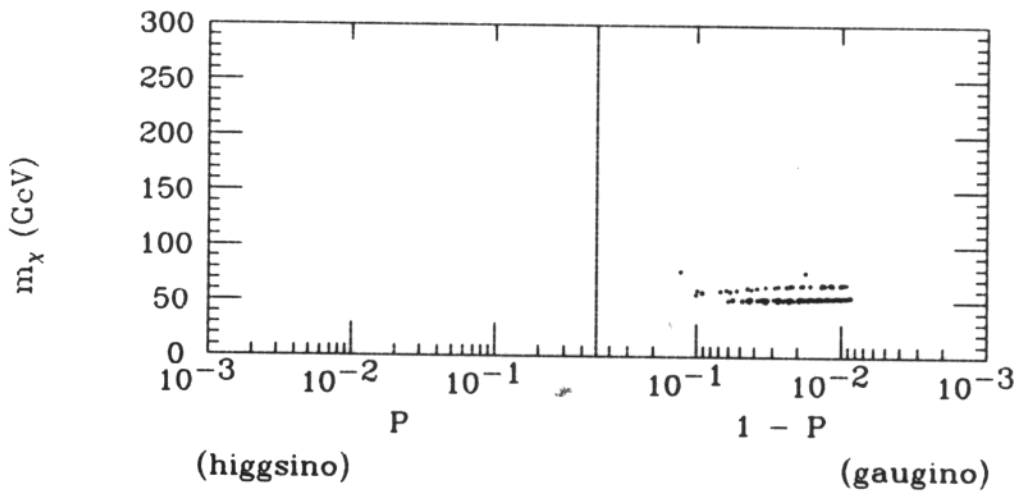


Figure 13 (b)

Excluded Configurations

$$\chi_{\text{red}}^2 \geq 4$$

Flattened Galactic Halo with $f = 0.5$

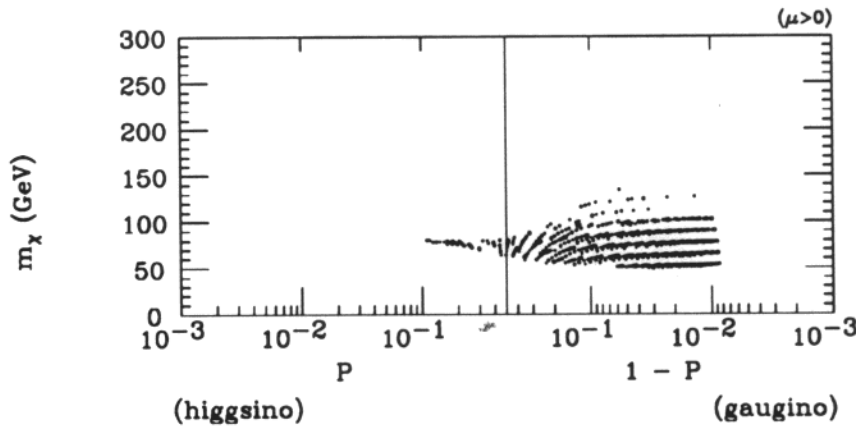


Figure 14 (b)

Figure 1: Scatter plot in the P - m_χ plane for the configurations excluded by the BESS 95 data when the galactic halo is flattened by a factor of $f = 0.5$.

Excluded Configurations

$$\chi_{\text{red}}^2 \geq 4$$

Flattened Galactic Halo with $f = 0.5$

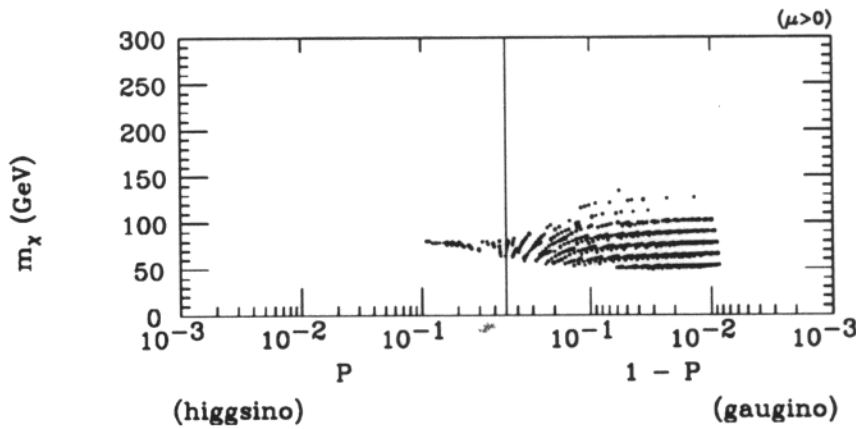


Figure 14 (b)

Figure 1: Scatter plot in the P - m_χ plane for the configurations excluded by the BESS 95 data when the galactic halo is flattened by a factor of $f = 0.5$.

Antiprotons & Direct Detection

$$\rho_l = 0.4 \text{ GeV cm}^{-3}$$

$$\chi_{\text{red}}^2 \leq 2.2 \quad 0.03 \leq \Omega h^2 \leq 0.7 \quad \text{c)}$$

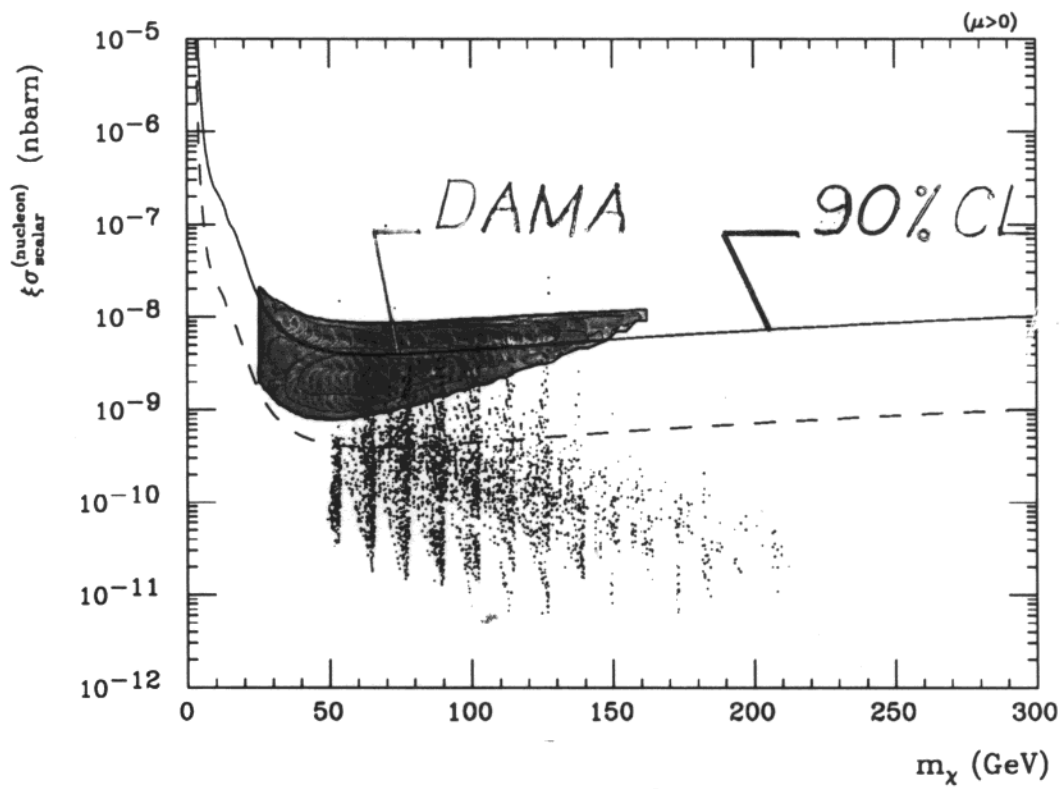
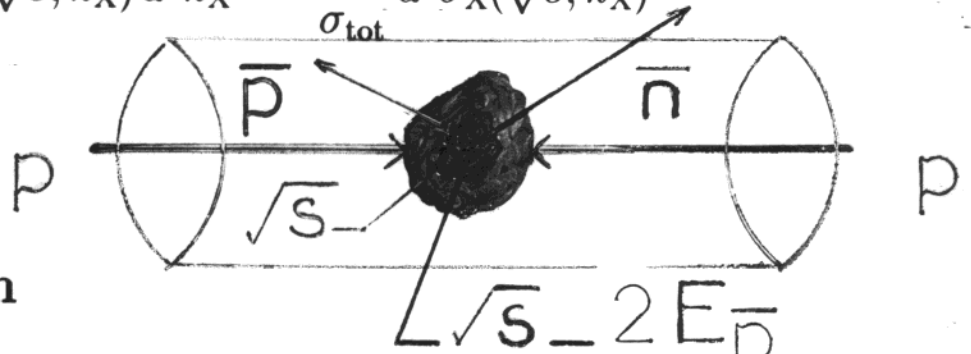


Figure 15 (b)

Figure 2: Scatter plot of the values of $\xi \sigma_{\text{scalar}}^{\text{nucleon}}$ versus the neutralino mass for the configurations compatible with BESS 95. The region delimited by a closed contour is the one singled out by the DAMA experiment as possibly indicative of an annual modulation effect.

Antideuteron Production in p-p Collisions & Neutralino Annihilations

$$dN_X = \mathcal{F}_X(\sqrt{s}, \vec{k}_X) d^3 k_X = \frac{1}{\sigma_{\text{tot}}} d^3 \sigma_X(\sqrt{s}, \vec{k}_X)$$



1) Factorization

$$\mathcal{F}_{\bar{p}, \bar{n}}(\sqrt{s}, \vec{k}_{\bar{p}}, \vec{k}_{\bar{n}}) = \frac{1}{2} \mathcal{F}_{\bar{p}}(\sqrt{s}, \vec{k}_{\bar{p}}) \mathcal{F}_{\bar{n}}(\sqrt{s} - 2E_{\bar{p}}, \vec{k}_{\bar{n}}) + (\vec{k}_{\bar{p}} \leftrightarrow \vec{k}_{\bar{n}})$$

2) Coalescence



$$\mathcal{F}_{\bar{D}}(\sqrt{s}, \vec{k}_{\bar{D}}) d^3 k_{\bar{D}} = \int d^3 k_{\bar{p}} d^3 k_{\bar{n}} C(\vec{k}_{\bar{p}}, \vec{k}_{\bar{n}}) \mathcal{F}_{\bar{p}, \bar{n}}(\sqrt{s}, \vec{k}_{\bar{p}}, \vec{k}_{\bar{n}})$$

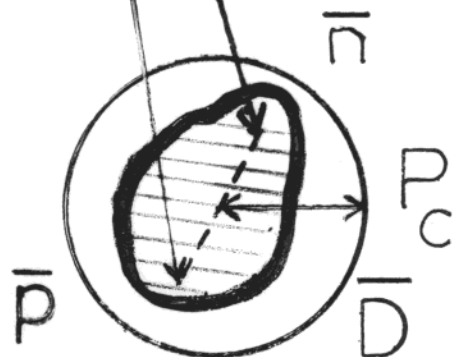
Because $\vec{k}_{\bar{p}} + \vec{k}_{\bar{n}} = \vec{k}_{\bar{D}}$ and $\vec{k}_{\bar{p}} - \vec{k}_{\bar{n}} = 2\vec{\Delta}$, then

$$\mathcal{F}_{\bar{D}}(\sqrt{s}, \vec{k}_{\bar{D}}) = \int d^3 \vec{\Delta} C(\vec{\Delta}) \mathcal{F}_{\bar{p}, \bar{n}}\left(\sqrt{s}, \vec{k}_{\bar{p}} = \frac{\vec{k}_{\bar{D}}}{2} + \vec{\Delta}, \vec{k}_{\bar{n}} = \frac{\vec{k}_{\bar{D}}}{2} - \vec{\Delta}\right)$$

$$\mathcal{F}_{\bar{D}}(\sqrt{s}, \vec{k}_{\bar{D}}) \simeq \left\{ \int d^3 \vec{\Delta} C(\vec{\Delta}) \right\} \mathcal{F}_{\bar{p}, \bar{n}}\left(\sqrt{s}, \vec{k}_{\bar{p}} = \frac{\vec{k}_{\bar{D}}}{2}, \vec{k}_{\bar{n}} = \frac{\vec{k}_{\bar{D}}}{2}\right)$$

with

$$\int \frac{E_{\bar{D}}}{E_{\bar{p}} E_{\bar{n}}} d^3 \vec{\Delta} C(\vec{\Delta}) \simeq \left(\frac{m_{\bar{D}}}{m_{\bar{p}} m_{\bar{n}}} \right) \left(\frac{4}{3} \pi \bullet \right)$$



- p-p collisions

$$p + p \rightarrow \bar{p} \text{ or } \bar{D}$$

$$E_{\bar{D}} \frac{d^3\sigma_{\bar{D}}}{d^3k_{\bar{D}}} = \left(\frac{m_{\bar{D}}}{m_{\bar{p}} m_{\bar{n}}} \right) \left(\frac{4}{3} \pi \right) \times \frac{1}{2\sigma_{\text{tot}}} \times$$

$$\times \left\{ E_{\bar{p}} \frac{d^3\sigma_{\bar{p}}}{d^3k_{\bar{p}}} (\sqrt{s}, \vec{k}_{\bar{p}}) E_{\bar{n}} \frac{d^3\sigma_{\bar{n}}}{d^3k_{\bar{n}}} (\sqrt{s} - 2E_{\bar{p}}, \vec{k}_{\bar{n}}) + (\vec{k}_{\bar{p}} \leftrightarrow \vec{k}_{\bar{n}}) \right\}$$

- Neutralino Annihilations

$$\chi + \chi \rightarrow \bar{p} \text{ or } \bar{D}$$

$$\sqrt{s} = 2m_{\chi}$$

and

$$\frac{dN_{\bar{D}}}{dE_{\bar{D}}} = \left(\frac{4}{3} \pi \right) \left(\frac{m_{\bar{D}}}{m_{\bar{p}} m_{\bar{n}}} \right) \times \sum_h B_{\chi h} \left\{ \frac{dN_{\bar{p}}^h}{dE_{\bar{p}}} (E_{\bar{p}} = E_{\bar{D}}/2) \right\}^2$$

whereas

$$\frac{dN_{\bar{p}}}{dE_{\bar{p}}} = \sum_h B_{\chi h} \frac{dN_{\bar{p}}^h}{dE_{\bar{p}}}$$

$$q_X^{\text{susy}} = \langle \sigma v \rangle \frac{dN_X}{dE_X} \left\{ \frac{\rho_X}{m_X} \right\}^2$$

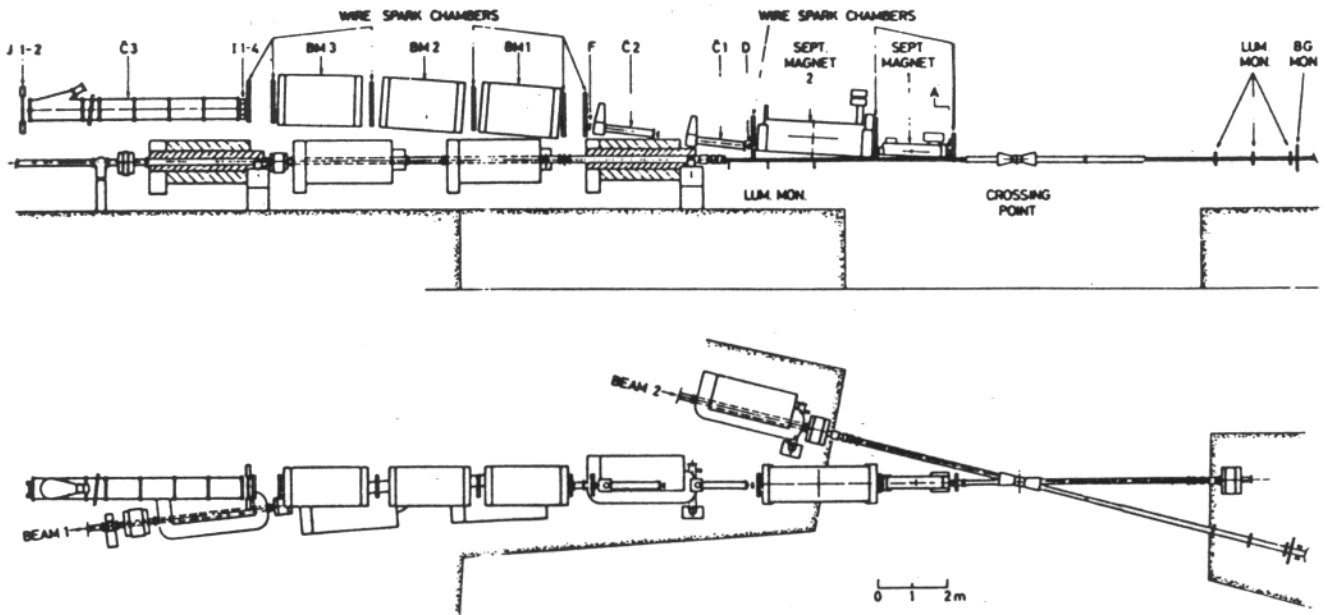


Fig. 1. Side and top views of the small angle spectrometer (SAS), located above the ISR beam 1. BM 1,2,3 are bending magnets, LUM MON is a set of counters for measuring the ISR luminosity. BG MON is a set of counters used to distinguish beam/beam from beam/gas events (see text).

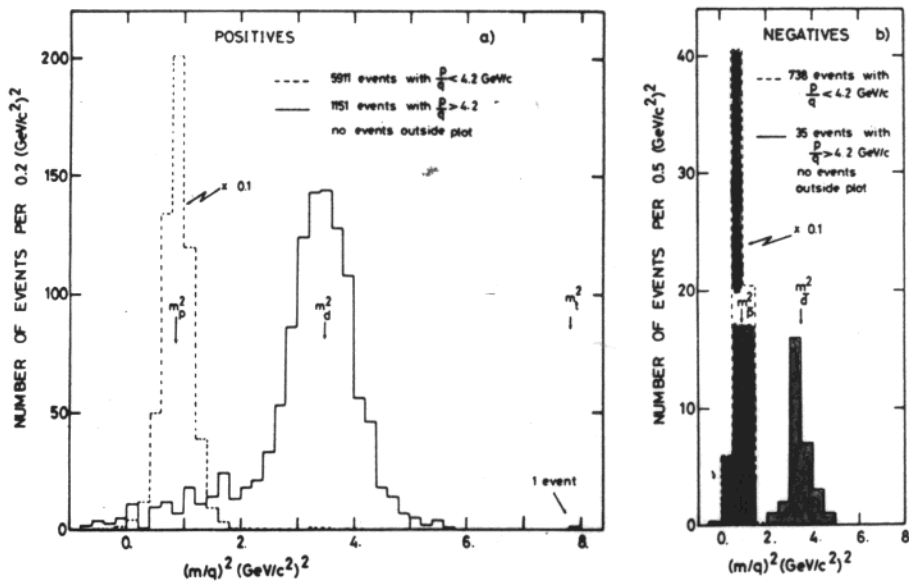


Fig. 3. (a) Distribution of $(m/q)^2$ for positive particles with $(p/q) < 4.2$ GeV/c (dashed lines) and $(p/q) > 4.2$ GeV/c (full lines). m is the particle mass and q its charge per electron charge. (b) Same as (a) for negative particles.

ISR 1975

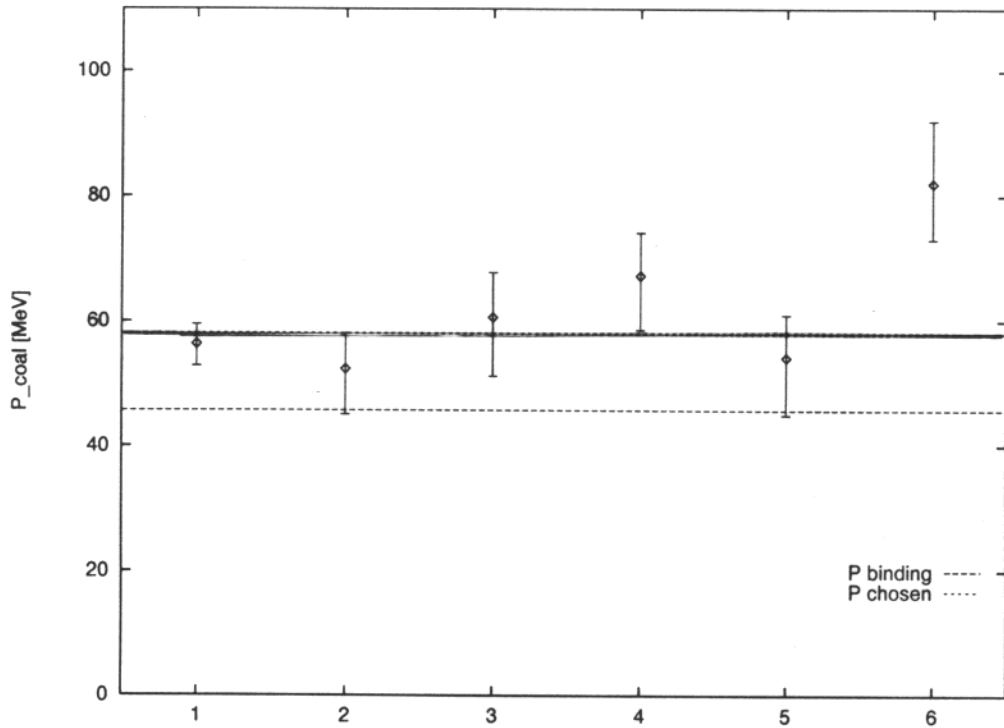


Figure 1: This figure displays various experimental constraints on the coalescence momentum P_{coal} , the only free parameter of the model discussed in section (2). Points 1 and 2 are from Serpukhov with $\sqrt{s} = 11.5$ GeV while all the other data have been collected at the ISR at $\sqrt{s} = 53$ GeV. A coalescence momentum P_{coal} of order 60 MeV provides a reasonable fit of all the points but the last one.

$$pp \longrightarrow \overline{D} X$$

$$\text{Valeurs de } E \frac{d^3\sigma}{dp^3} \Big|_{\overline{D}}$$

#	p_L	p_T	$\mu\text{b}/\text{GeV}^2$
1	0	1.14	$(9.1 \pm 1.6) \times 10^{-5}$
2	0	1.50	$(2.2 \pm 0.8) \times 10^{-5}$
3	4.8	0.16	$(6.8 \pm 2.7) \times 10^{-2}$
4	5.56	0.21	$(5.6 \pm 1.9) \times 10^{-2}$
5	6.62	0.30	$(1.4 \pm 0.6) \times 10^{-2}$

Susy & Secondary Antideuteron Flux

$$m_\chi = 60.4 \text{ GeV}$$

$$P = 0.84 : \text{ mostly gaugino}$$

$$\Omega_\chi h^2 = 0.02$$

Donato, Fornengo & Salati, preliminary

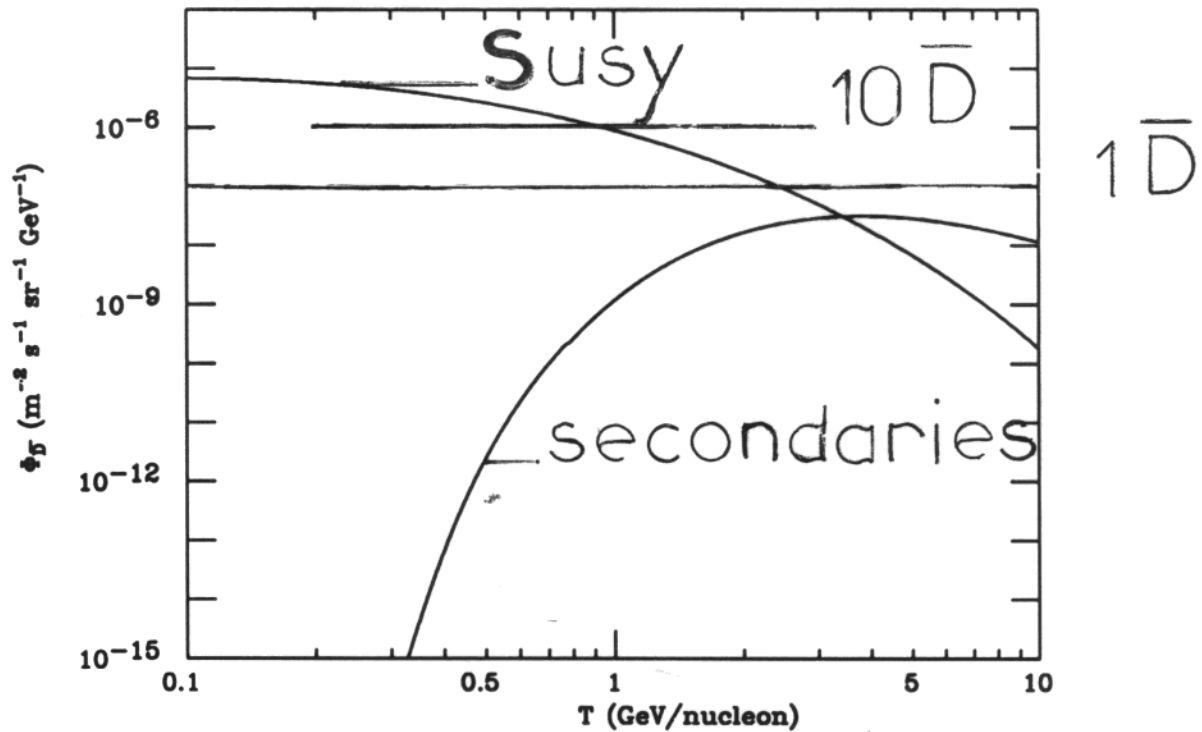


Figure 9: The IS antideuteron flux is presented as a function of the kinetic energy per nucleon. The primary SUSY flux emerges at low energy while the secondary antideuteron fluxes are most visible at energies in excess of a few GeV/n.

The instruments

- Balloon-borne experiments.

IMAX : 16 antiprotons accumulated with an acceptance of $\Omega_{ST} = 260 \text{ m}^2 \text{ sr s}$

BESS 95 : 43 low-energy antiprotons with an acceptance of $\Omega_{ST} = 2 \times 10^3 \text{ m}^2 \text{ sr s}$

Stratosphere $\Rightarrow 5 \text{ g cm}^{-2}$ of air

Geomagnetic suppression factor

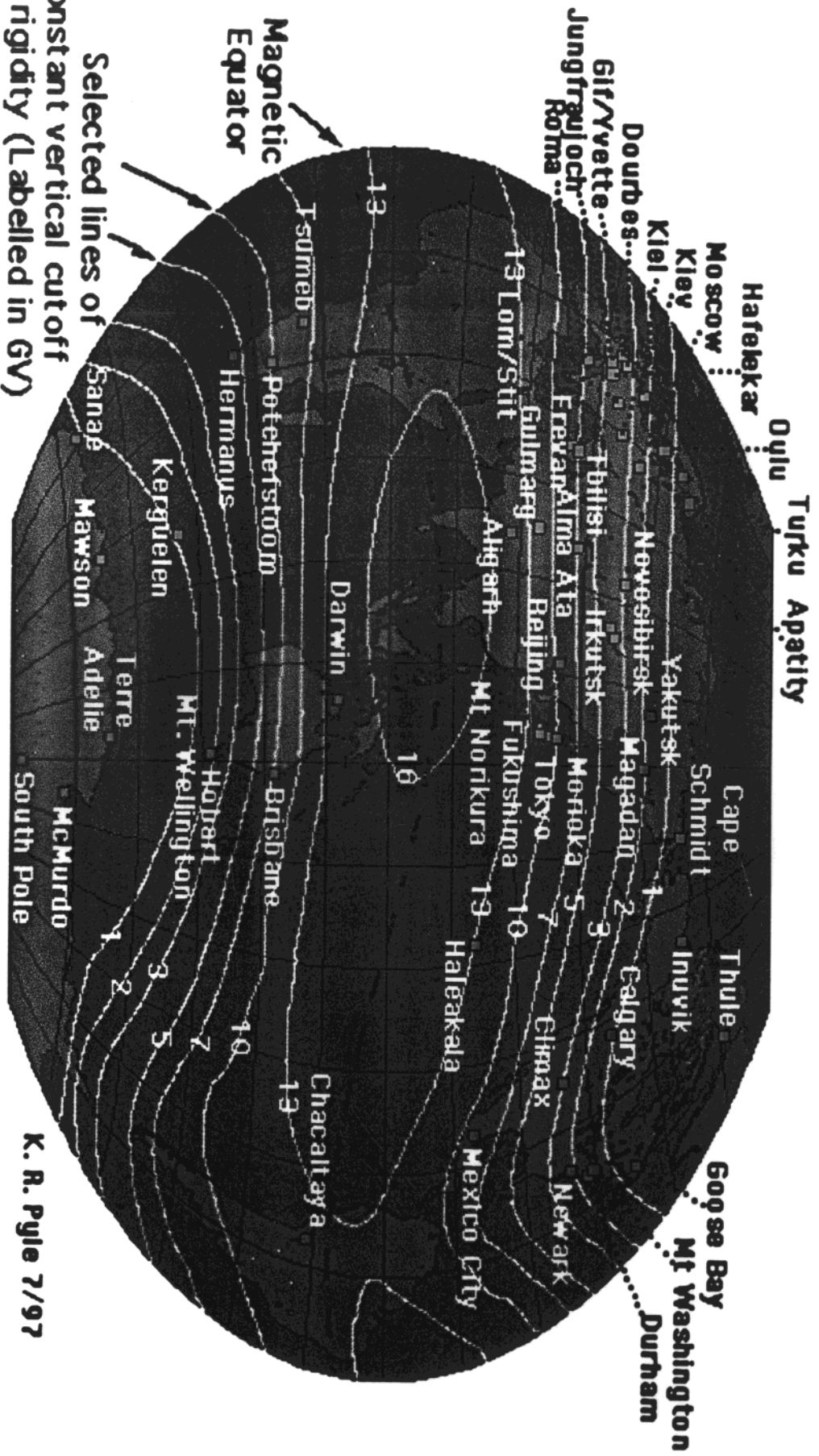
$$\left. \frac{pc}{Ze} \right|_{\min} = \frac{\mu}{R^2} \frac{\cos^4 L}{\left[1 + \sqrt{1 + \cos \theta \cos^3 L} \right]^2} \quad \text{with} \quad \frac{\mu}{R_0^2} = 60 \text{ GV}$$

- Space instruments.

PAMELA : small acceptance but polar orbit
 $\Rightarrow 10^{-7}$ level on anti- ^4He

AMS : large acceptance $\Omega_{ST} = 1.8 \times 10^4 \text{ m}^2 \text{ sr s}$ but mostly equatorial $\Rightarrow 200$ antiprotons during the shuttle flight of 1998 $\Rightarrow 10^{-9}$ level on anti-nuclei during the space station stage with $\Omega_{ST} = 2 \times 10^7 \text{ m}^2 \text{ sr s}$

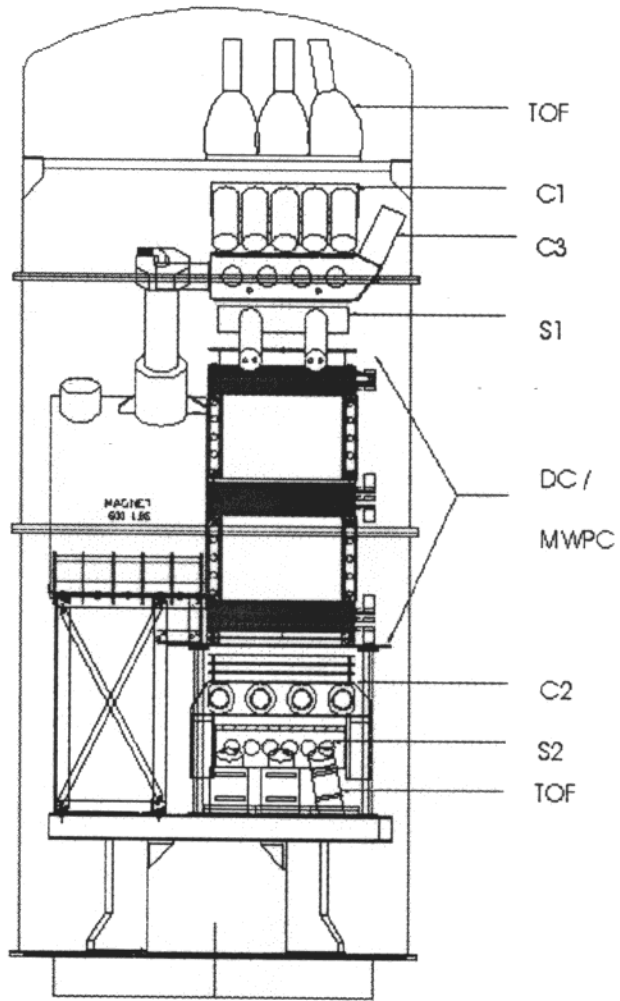
Cosmic Ray Neutron Monitors, 1997



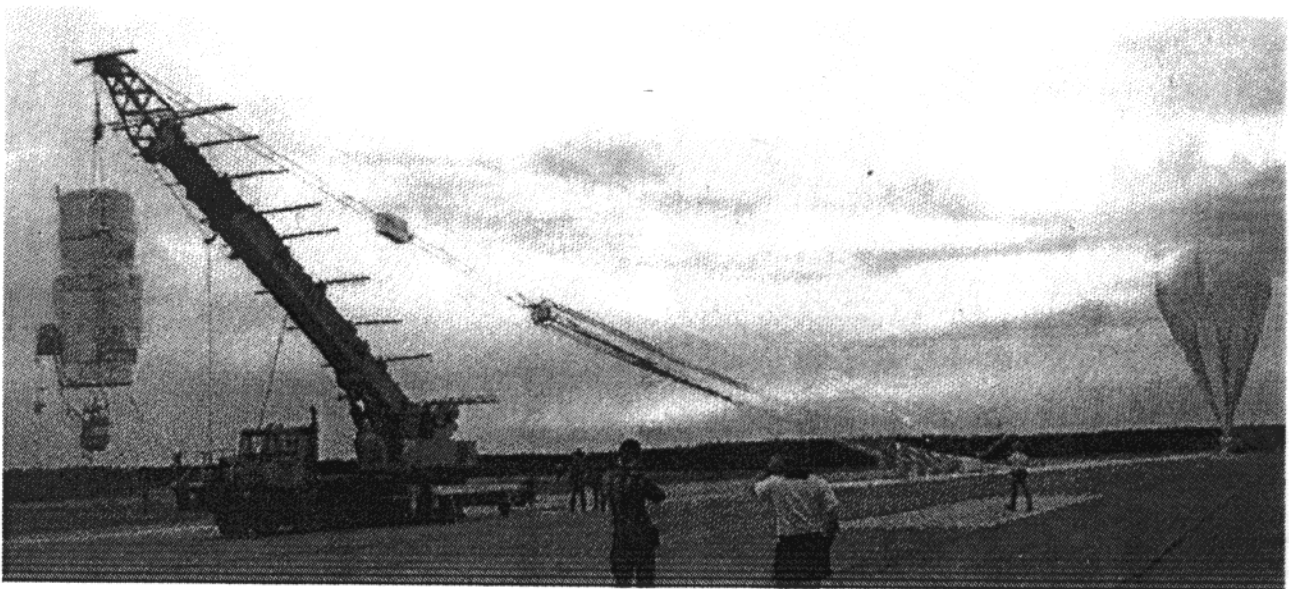
Selected lines of constant vertical cutoff rigidity (Labelled in GV)

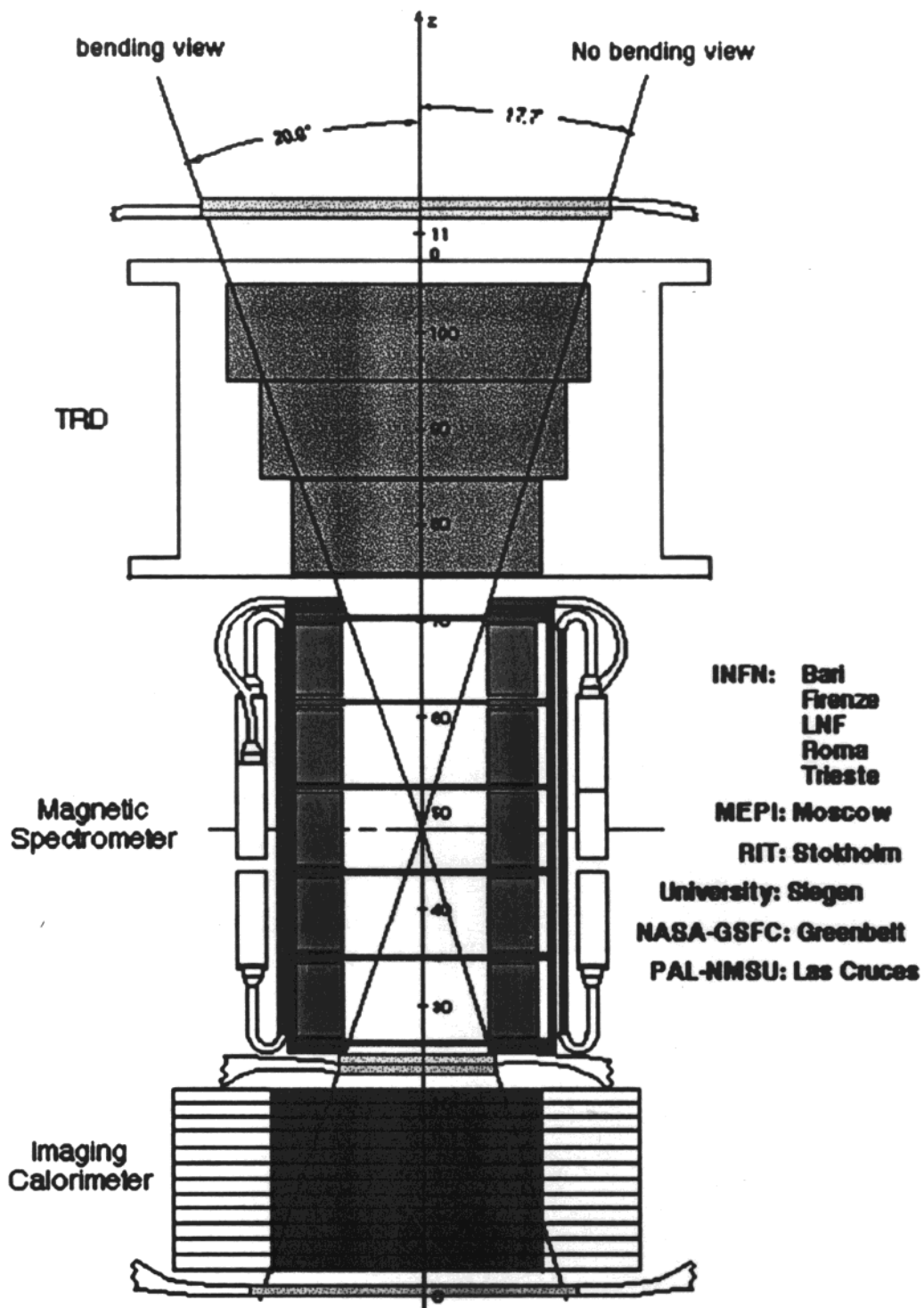
Magnetic Equator

K. R. Pyle 7/97

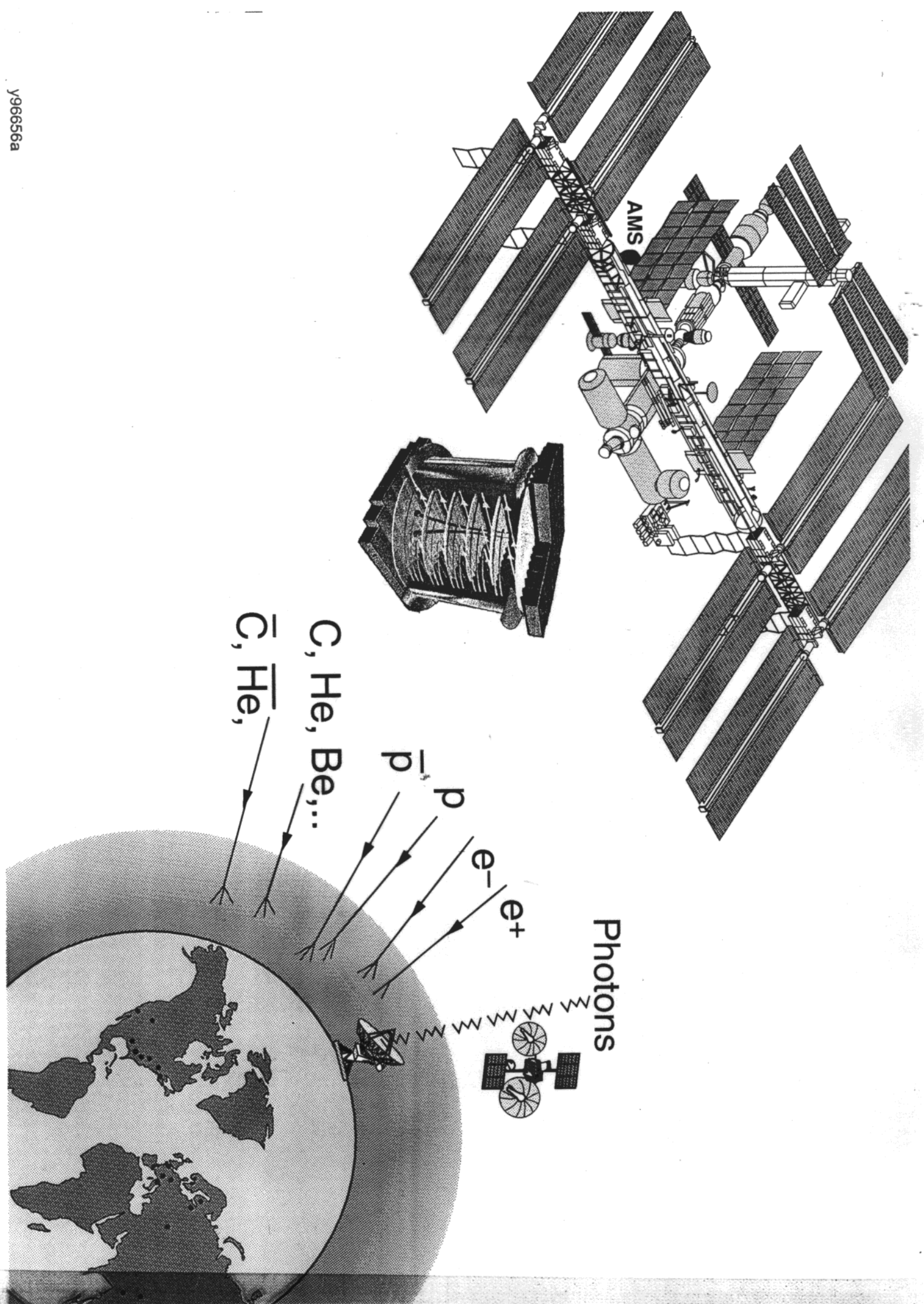


Balloon-borne IMAX





PAMELA



AMS

Photons

e^-
 e^+

p
 \bar{p}

C, He, Be, \dots

\bar{C}, \bar{He}, \dots

Geomagnetic cut-off & AMS acceptance included.

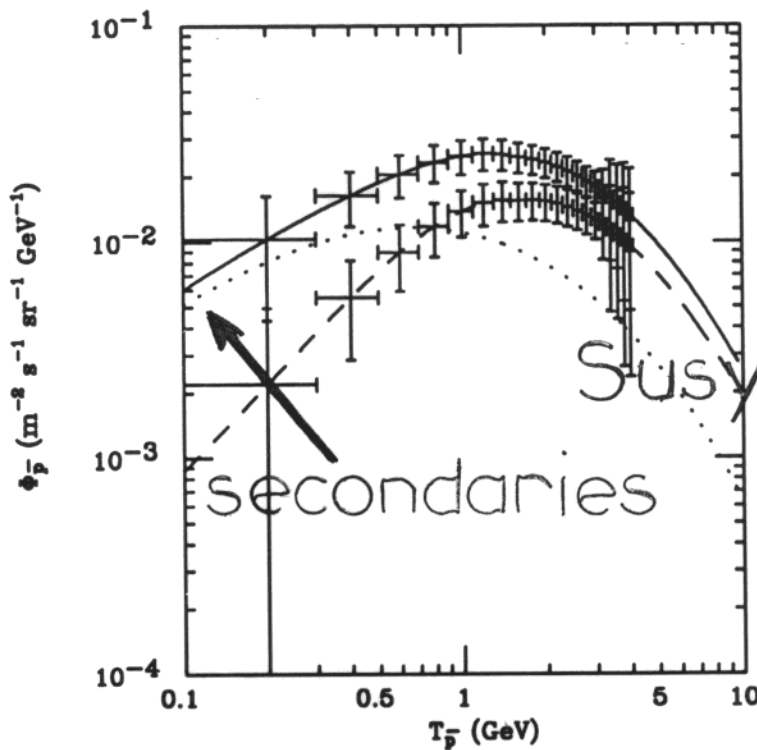


Figure 18

Figure 3: Expected distribution of measurements with the AMS Shuttle flight. The lower sequence of crosses corresponds to secondary antiprotons alone while the upper sequence includes a significant contribution of primary antiprotons from neutralino annihilations.

IS proton flux \rightarrow 25% uncertainty on $\Phi_{\bar{p}}$.

Uncertainty on the diffusion coefficient K

$\rightarrow \Delta\Phi_{\bar{p}}/\Phi_{\bar{p}} = 45\%$ at 100 MeV and 20% at 600 MeV.

Antideutrium detectability

On board AMS during the space station phase

- Between 100 MeV & 3 GeV : SUSY antideutons.

$$\Omega_{\text{STE}} = 1.6 \times 10^7 \text{ m}^2 \text{ sr s GeV} \quad \text{solar minimum}$$

$$\Omega_{\text{STE}} = 9.8 \times 10^6 \text{ m}^2 \text{ sr s GeV} \quad \text{solar maximum}$$

$$\Phi_{\bar{\text{D}}} = 10^{-6} \text{ m}^{-2} \text{ sr}^{-1} \text{ s}^{-1} \text{ GeV}^{-1} \quad \Rightarrow \quad 10 \bar{\text{D}}$$

- Between 3 GeV & 10 GeV : secondary antideutons.

$$\Omega_{\text{STE}} = 1.3 \times 10^8 \text{ m}^2 \text{ sr s GeV} \quad \text{solar minimum}$$

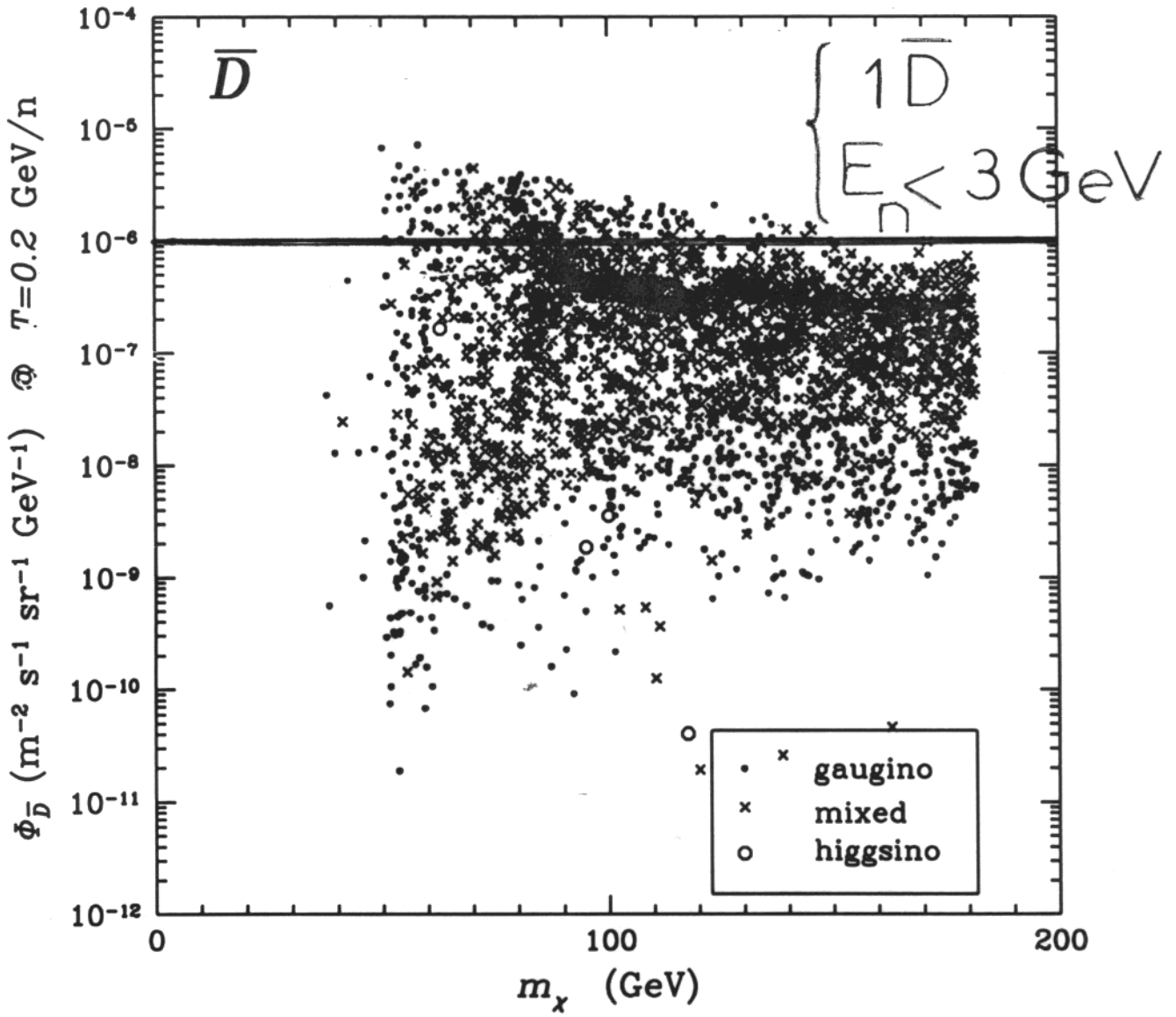
$$\Omega_{\text{STE}} = 1.1 \times 10^8 \text{ m}^2 \text{ sr s GeV} \quad \text{solar maximum}$$

$$\Phi_{\bar{\text{D}}} = 10^{-8} \text{ m}^{-2} \text{ sr}^{-1} \text{ s}^{-1} \text{ GeV}^{-1} \quad \Rightarrow \quad 1 \text{ to } 2 \bar{\text{D}}$$

On board BESS

$$\Omega_{\text{STE}} = 1\text{-}2 \times 10^4 \text{ m}^2 \text{ sr s GeV}$$

Should run for a long time !



A word of caution !

- Better determination of the proton flux needed between 5 and 20 GeV.
- Consistent determination of the diffusion coefficient K from B/C measurements.
- Solar modulation $\phi_{\bar{p}}^{\text{TOA}}$ vs $\phi_{\bar{p}}^{\text{IS}}$
- Antiproton energy losses
Ionisation
Elastic scattering on IS nuclei
- Diffusive reacceleration
Heinbach & Simon ApJ 441 (1995) 209
- Nucleus-nucleus interactions
Simon, Molnar & Roesler \neq Gaisser & Schaefer

However !

The larger the secondary antiproton flux at low energy, the more constrained is SUSY by the BESS and AMS data.

Proton flux poorly determined at high energy.

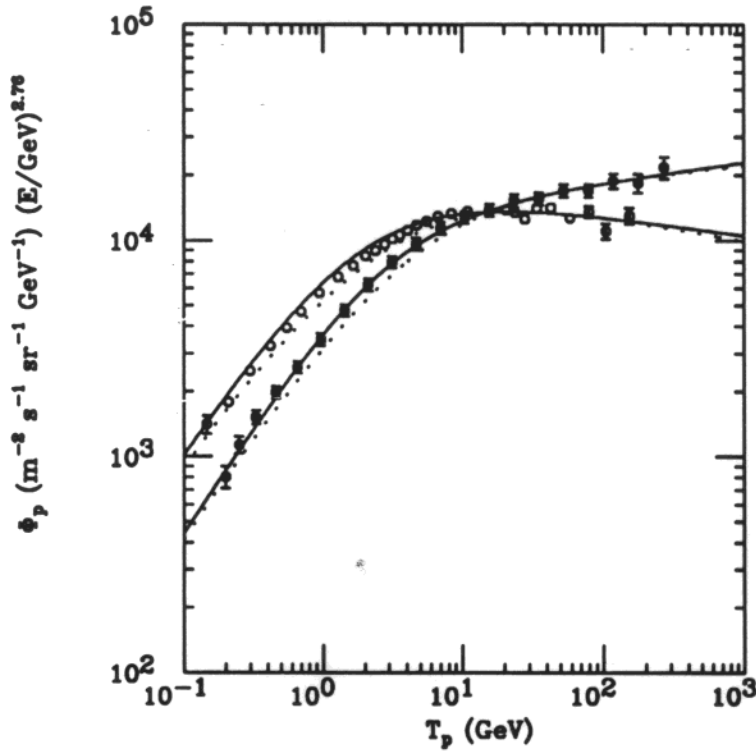


Figure 2

Figure 4: TOA spectra of IMAX (full circles) and of CAPRICE (open circles) with our best-fit curves with parametrization as a function of the energy (solid lines) and the rigidity (dotted lines).

$$\Phi_p^{\text{IS}} = A \beta E^{-\gamma}$$

$$A = (1.595 \pm 0.365) \text{ protons cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1} \text{ GeV}^{-1}$$

$$\gamma = 2.76 + 0.13 - 0.15$$

The B/C Ratio and the Diffusion Coefficient K

- $\mathcal{R} \geq \mathcal{R}_0 : K = K_0 \mathcal{R}^{0.6}$
- $\mathcal{R} \leq \mathcal{R}_0 : K = K_0 \mathcal{R}_0^{0.6-\delta} \mathcal{R}^\delta$

D. Maurin (LAPTH) : preliminary

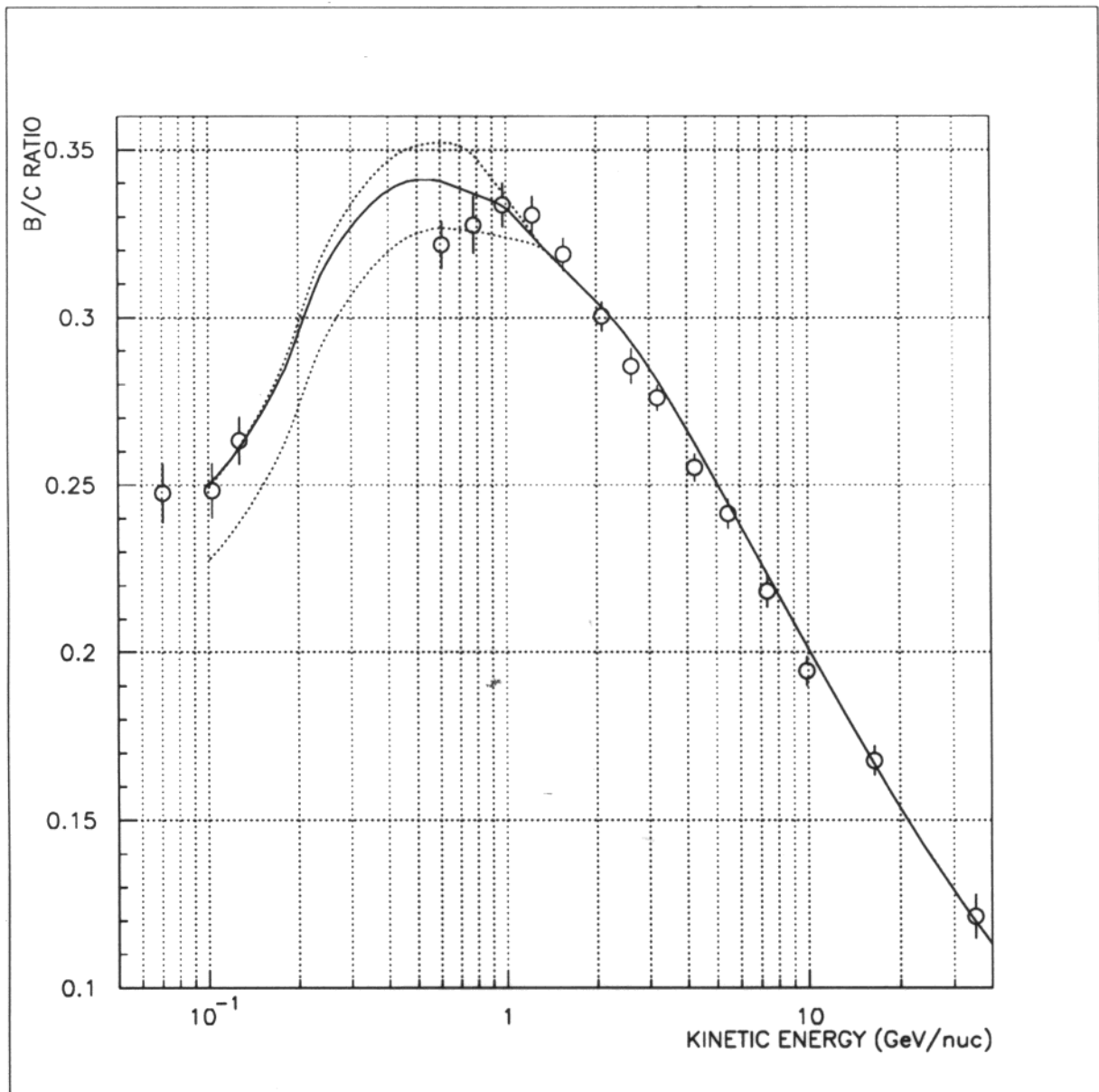


Figure 8: The B/C ratio as a function of the kinetic energy per nucleon. Ionization losses and reacceleration have not been included. The isotopes ^{10}B and ^{11}B are produced from the spallation of the CNO elements. Cross sections are borrowed from Read & Viola (1984) while the CNO abundances are from Duvernois *et al.* (1996). The solid curve corresponds to $\delta = 0.08$ and $\mathcal{R}_0 = 3.8$ Gv whereas, for the dashed curves, $\delta = 0$ and $\mathcal{R}_0 = 3.2$ Gv (max) and 4.2 Gv (min).

Solar Modulation : Force Field Approximation.

secondaries

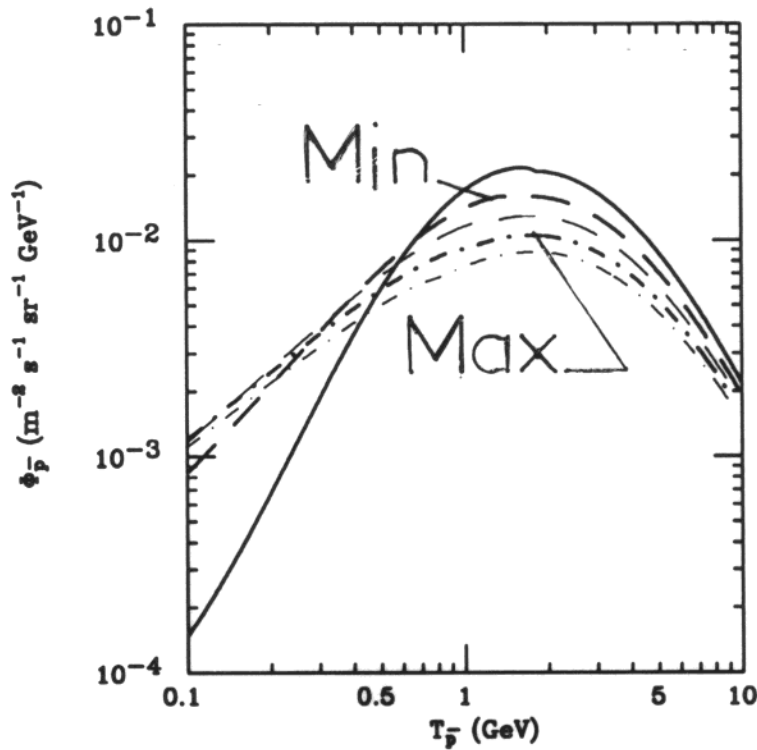


Figure 8

Figure 5: Solar modulation of the IS median secondary antiproton flux (solid line). Thick dashed (thick dot-dashed) line is the solar-modulated spectrum at minima (maxima) when the modulation parameter Δ is obtained from a parametrization of the primary proton spectrum as a function of energy. The light curves refer to a parametrization as a function of rigidity.

$$E_n^{\text{TOA}} = E_n^{\text{IS}} - \frac{Z}{A} \cdot ? \quad \text{and} \quad \Phi^{\text{TOA}} / ? = (p_n^{\text{TOA}} / p_n^{\text{IS}})^2$$

for a rigidity ≤ 200 MV and with $0.32 \text{ GV} \leq \Delta \leq 0.8 \text{ GV}$

Solar Modulation Effect on Primary Susy Antiprotons

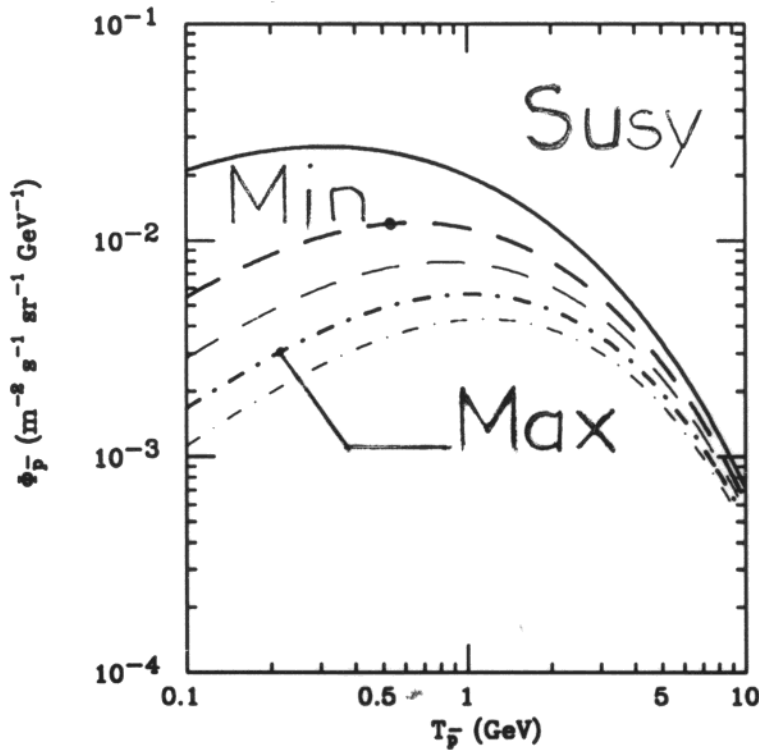


Figure 9

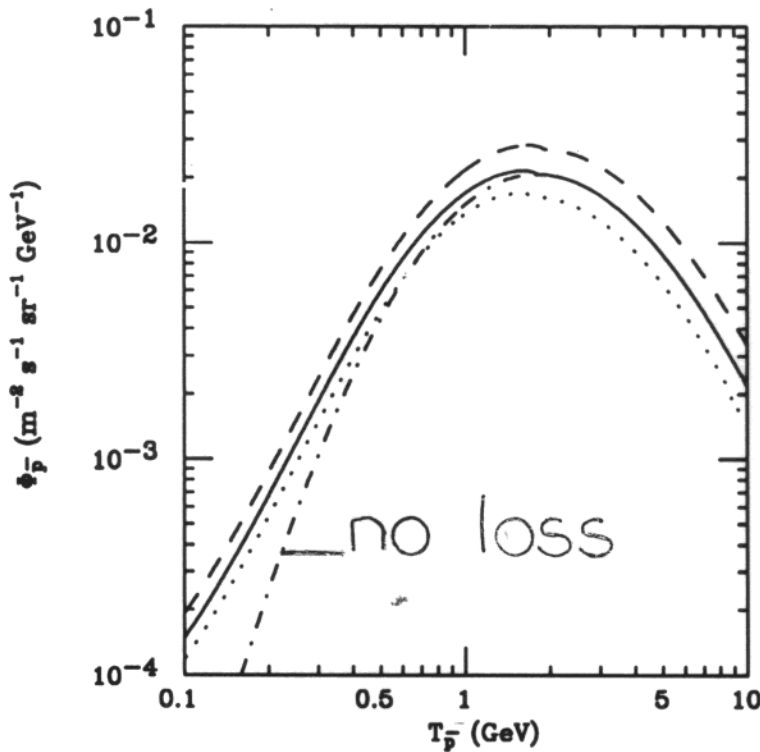
Figure 6: Solar modulation of the IS antiproton flux, due to neutralino annihilation for the representative neutralino configuration with $m_{\chi} = 62$ GeV, $P = 0.98$ and $\Omega_{\chi} h^2 = 0.11$. Solid line is the IS spectrum. Thick dashed (thick dot-dashed) line is the solar-modulated spectrum at minima (maxima) when the modulation parameter Δ is obtained from a parametrization of the primary proton spectrum as a function of energy. The light curves correspond to a parametrization as a function of rigidity.

Energy losses

- Ionisation losses
- Elastic scattering on Interstellar H

inelastic_no annihilation.

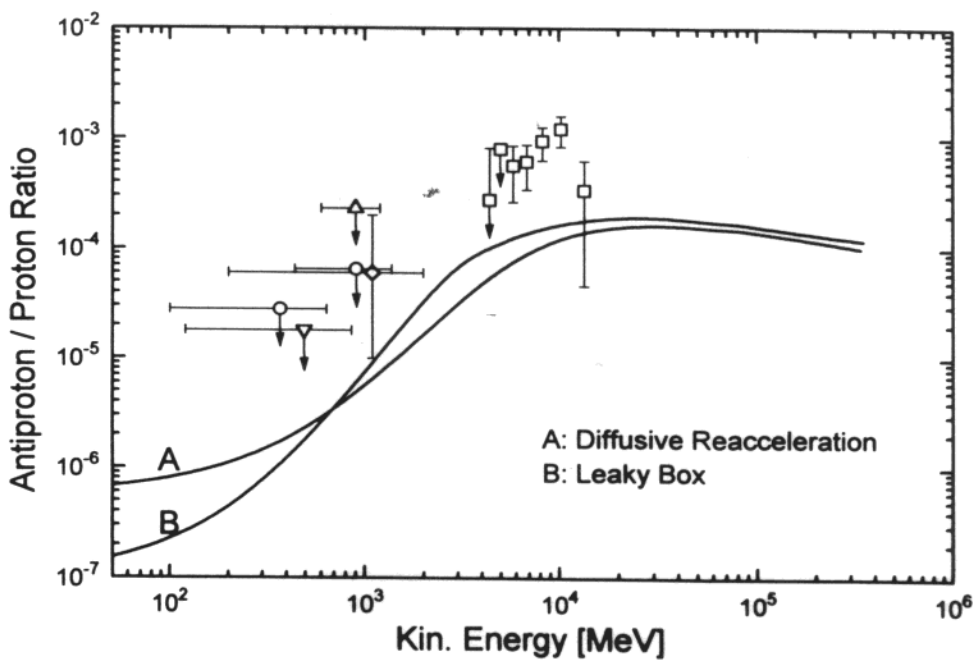
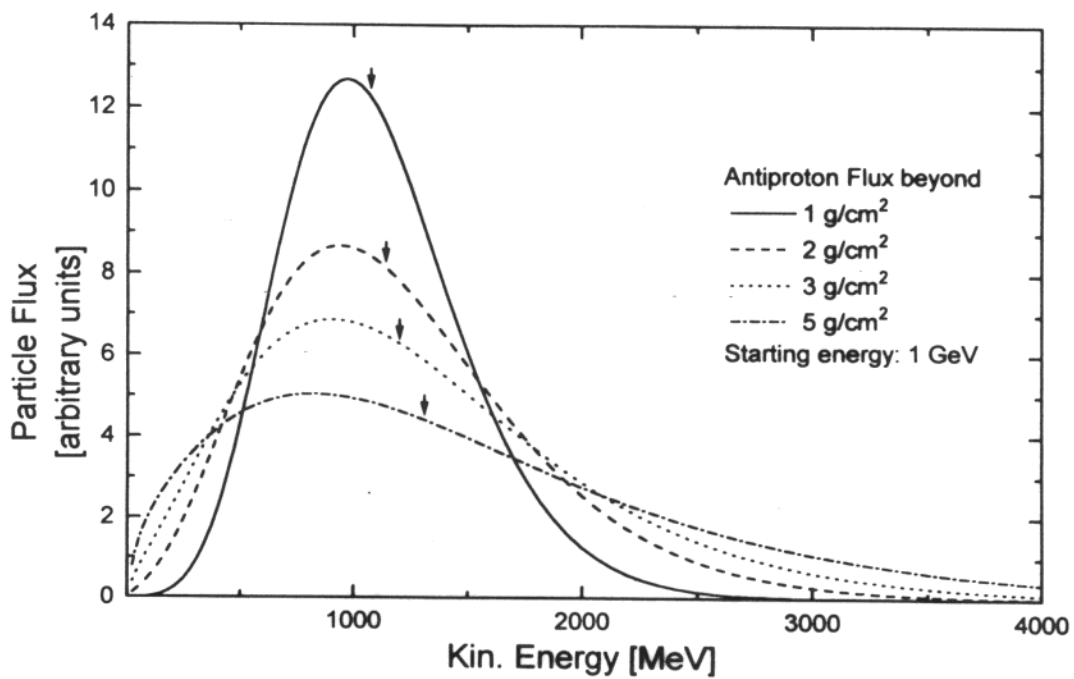
\bar{p} loses $T_{\bar{p}}/2$



$$\frac{dn}{dt} = - \frac{dbn}{dE}$$

Figure 5 $b = \dot{E} = - \frac{T_{\bar{p}}}{2} \sigma_{HP} n v_{\bar{p}}$

Figure 7: IS secondary antiproton spectra as a function of the \bar{p} kinetic energy. Solid, dotted and dashed lines denote the fluxes obtained from the median, minimal and maximal IS primary proton fluxes. The dot-dashed line denotes the median \bar{p} spectrum, when the \bar{p} energy losses are neglected.



Conclusions

1. \bar{p} data & Secondary source
OK
2. Data precision : Susy sensitive
3. However, large systematic
uncertainties for secondary \bar{p}
Hope!
4. \bar{D} : smoking gun
Clear signal - Low flux