

1) Neutrino - Electron Elastic Cross Section

T. Hooft *Phys Lett B37(1974)195*
 John N. Bahcall *Rev. Mod. Phys. 52(1980)52*

T. Hooft Form

$$\frac{d\sigma}{dE_e} = \frac{G_F^2 m_e}{2\pi} \left[(g_V + g_A)^2 + \left(1 - \frac{E_e}{E_\nu}\right)^2 (g_V - g_A)^2 + \frac{m_e E_e}{E_\nu^2} (g_A^2 - g_V^2) \right]$$

*This is also done K. E. Gribben
 see E. Commins + P. H. Bucksbaum
 Weak Interactions of Leptons + Quarks
 Error in T. Hooft paper
 will not be*

~~scribble~~

$$g_V = \frac{1}{2} + 2 \sin^2 \theta_w$$

$$g_A = \frac{1}{2}$$

$$\frac{g_V + g_A}{2} = \frac{1 + 2 \sin^2 \theta_w}{2} = \frac{1}{2} + \sin^2 \theta_w$$

$$\frac{g_V - g_A}{2} = \sin^2 \theta_w$$

$$\frac{g_A^2 - g_V^2}{4} = \frac{\frac{1}{4} - (\frac{1}{2} + 2 \sin^2 \theta_w + 4 \sin^4 \theta_w)}{4} = -\frac{1}{2} \sin^2 \theta_w + \sin^4 \theta_w = -\sin^2 \theta_w \left(\frac{1}{2} + \sin^2 \theta_w \right)$$

Taking $g_L = \frac{1}{2} + \sin^2 \theta_w = \frac{g_V + g_A}{2}$

+ for ν_e
 $g_R = \sin^2 \theta_w = \frac{g_V - g_A}{2}$

- for ν_μ
 $g_L g_R = \sin^2 \theta_w \left(\frac{1}{2} + \sin^2 \theta_w \right) = -\frac{g_A^2 - g_V^2}{4}$

Bahcall Notation

$$\frac{d\sigma}{dE_e} = \frac{4G_F^2 m_e}{2\pi} \left[\left(\frac{g_V + g_A}{2}\right)^2 + \left(\frac{g_V - g_A}{2}\right)^2 \left(1 - \frac{E_e}{E_\nu}\right)^2 + \frac{m_e E_e}{E_\nu^2} \left(\frac{g_A^2 - g_V^2}{4}\right) \right]$$

$$\frac{d\sigma}{dE_e} = \frac{2m_e E_e^2}{\pi} \left[g_L^2 + g_R^2 \left(1 - \frac{E_e}{E_\nu}\right)^2 + \frac{m_e E_e}{E_\nu^2} g_L g_R \right]$$

Let $E_e = m_e T$ $E_\nu = m_e q$

$$\frac{d\sigma}{dT} = \left(\frac{2m_e^2 G_F^2}{\pi}\right) \left[g_L^2 + g_R^2 \left(1 - \frac{T}{q}\right)^2 + \frac{T}{q^2} g_L g_R \right] \leftarrow \text{Bahcall form}$$

$\sigma_0 = \frac{2m_e^2 G_F^2}{\pi} \times (\hbar c)^2$ $G_F = 1.166 \times 10^{-5} \text{ GeV}^{-2} = \text{Fermi } \beta \text{ decay constant.}$

$$= \left(\frac{2}{\pi}\right) (5.11 \times 10^{-3} \text{ GeV})^2 (1.166 \times 10^{-5} \text{ GeV}^{-2})^2 (1.97 \times 10^{-14} \text{ GeV-cm})^2 = 8.77 \times 10^{-45} \text{ cm}^2$$

Integrating over energy transfers $T_{max} = \frac{2q^2}{1+q}$ to T_{min}

$$\sigma_{\text{tot}, q}(T_{min}) = \sigma_0 \left[(g_L^2 + g_R^2)(T_{max} - T_{min}) - \left[\frac{g_R^2}{q} + \frac{g_L g_R}{2q^2} \right] (T_{max}^2 - T_{min}^2) + \frac{g_R^2}{3q^2} (T_{max}^3 - T_{min}^3) \right]$$

Example $E_\nu = 860 \text{ KeV}$ $q = 1.683 = ①$ $g_L = \frac{.72}{.11}$ $g_L^2 + g_R^2 = .5668 = ③$
~~scribble~~ $g_R = \frac{.11}{.11}$ $\frac{g_R^2}{q} + \frac{g_L g_R}{2q^2} = .0567 = ④$
 $T_{max} = 1.2975 = ②$ $\frac{g_R^2}{3q^2} = .0057 = ⑤$
 $T_{min} = 0$

$T = \frac{2 E_\nu^2 m_e^2}{1 + 2 E_\nu / m_e} T_{max}$
 $= \frac{2 E_\nu}{m_e} \frac{2 m_e E_\nu}{1 + 2 E_\nu / m_e}$
 $T = \frac{2 E_\nu}{m_e + 2 E_\nu}$

$$\sigma_{\text{tot}, q}(860 \text{ KeV}) = 8.77 \times 10^{-45} \text{ cm}^2 \left[.5668 - (.0567 + .0125) \right] = 8.77 \times 10^{-45} \text{ cm}^2 \times .4976 = 4.36 \times 10^{-45} \text{ cm}^2$$

$$\sigma_{\text{tot}, q}(0) = 8.77 \times 10^{-45} \text{ cm}^2 \left[.5668 - .0567 + .0125 \right] = 8.77 \times 10^{-45} \text{ cm}^2 \times .5226 = 4.59 \times 10^{-45} \text{ cm}^2$$

$$\sigma(100 \text{ KeV}) = 8.77 \times 10^{-45} \text{ cm}^2 \left[.6245 - .0937 + .0124 \right] = 8.77 \times 10^{-45} \text{ cm}^2 \times .5432 = 4.77 \times 10^{-45} \text{ cm}^2$$

A HIGH RATE SOLAR NEUTRINO DETECTOR WITH ENERGY DETERMINATION*

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- b) World Laboratory, Geneva, Switzerland
- c) CERN, Geneva, Switzerland
- d) INFN- Gran Sasso National Laboratory, Italy

Abstract:

A real time (non radiochemical) experiment is proposed which can detect the low energy solar neutrinos from the reactions $pp \rightarrow e^+ + d + \nu_e$ and $e^- + {}^7\text{Be} \rightarrow {}^7\text{Li} + \nu_e$, identify the individual spectral components, is sensitive to ν_μ , ν_τ and ν_e with an event rate of 20,000 per year.

Neutrino-electron (νe) elastic scattering, like Compton scattering, results in transfer of kinetic energy (T) to the electron with a characteristic Compton edge. The proposed detector is a time projection chamber (TPC) filled with 20 tons of He gas as the electron target with 6.3 ktons of liquid nitrogen as the shield. The TPC method allows determination of the direction of the initial recoil track and its energy from the total track length. From these, and knowledge of the sun's position, the energy of the initiating neutrino can be calculated thus allowing monoenergetic line sources to be identified. A parallel method to determine the spectral composition of the sources is by observation of Compton edges. The energy resolution of this detector from range ($\sigma_T/T \leq 5\%$) is sufficiently good to keep these edges sharp. An event rate of $2 \cdot 10^4/\text{year}$ can be attained, mainly from the pp and ${}^7\text{Be}$ channels, with 20 tons of He gas target. To keep the TPC size reasonable, the helium gas and the liquid nitrogen shield are pressurized to 25 bar at 77°K . The advantages of a He target and N_2 shield are: i) they have no long lived radioactive isotopes, ii) they freeze out potential radioactive impurities iii) they are cheap, easily available and safe iv) no muon induced radiogenesis in helium.

Detector calibration can be made with the pure EC radioactive source e.g. $e^- + {}^{37}\text{Ar} \rightarrow {}^{37}\text{Cl} + \nu_e$ ($E_\nu = 814\text{ keV}$, $t_{1/2} = 35\text{ days}$) placed just outside the shield. An activity of 1 MCi will produce ≈ 70 $\nu_e e$ elastic scatters/day, slightly more than the Standard Solar Model rate of $\approx 55/\text{day}$. Apart from calibration, events from the low T region of the electron recoil spectrum can determine an upper limit for the ν_e magnetic moment. This scattering, due to a purely electromagnetic interaction between the neutrino magnetic moment and the electron charge, does not require the sun's magnetic field to produce transversely polarized neutrinos.

* Presented by T. Ypsilantis at the Neutrino Telescopes Conference, Venice, Italy, 10-13/ 3/ 92.

Résumé:

Une expérience (non radio-chimique) en temps réel est proposée pour détecter les neutrinos solaires de basses énergies des réactions $pp \rightarrow e^+ + d + \nu_e$ et $e^- + {}^7\text{Be} \rightarrow {}^7\text{Li} + \nu_e$, et identifier les composantes du spectre. L'expérience est sensible aux neutrinos ν_μ , ν_τ et ν_e avec un taux de $2 \cdot 10^4$ événements détectés par année.

Le diffusion élastique neutrino-électron (νe) résulte, comme pour la diffusion Compton, au transfert d'une énergie cinétique T à l'électron avec une coupure caractéristique de l'effet Compton.

Solar Neutrino TPC in the Gran Sasso Laboratory

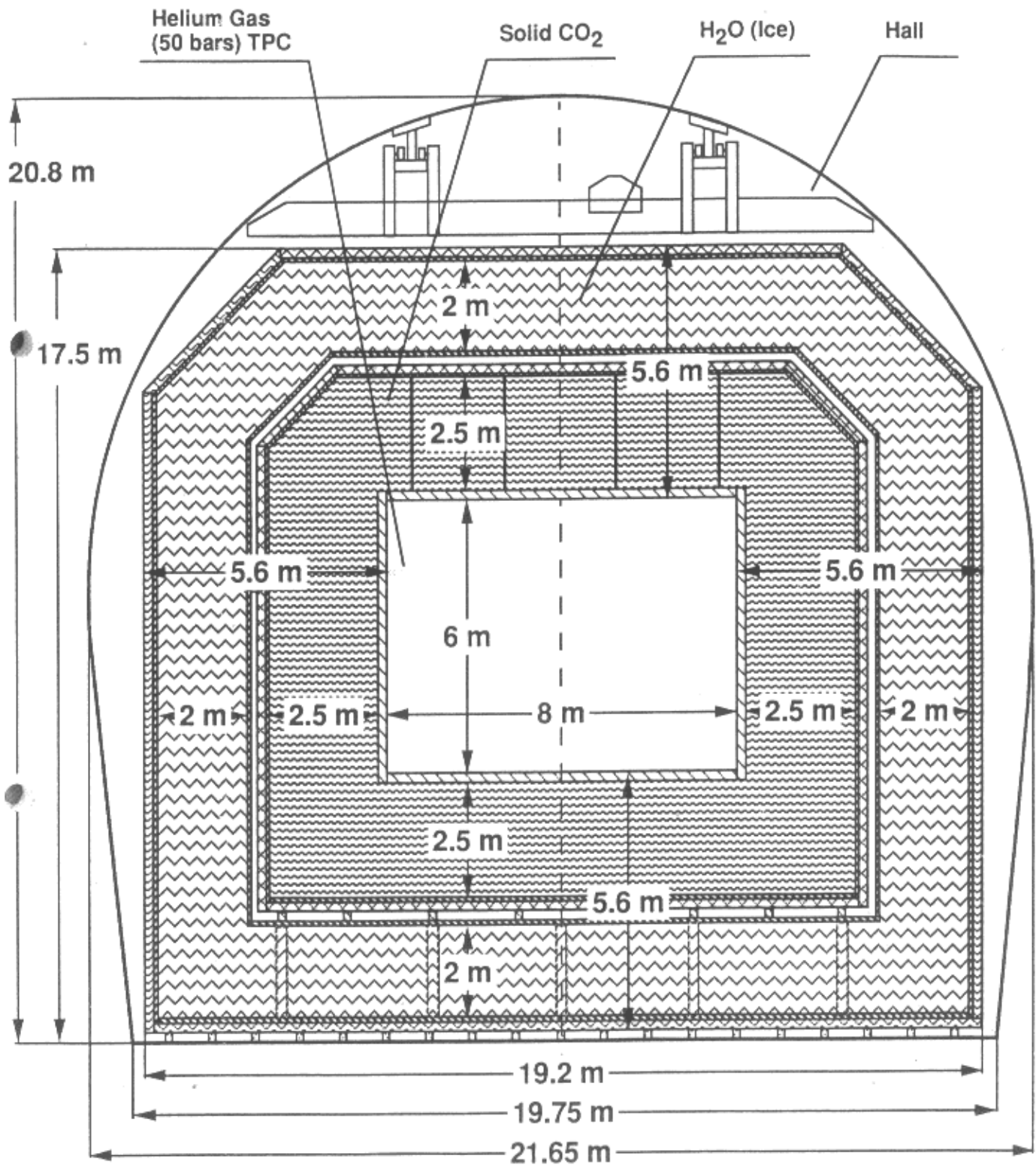
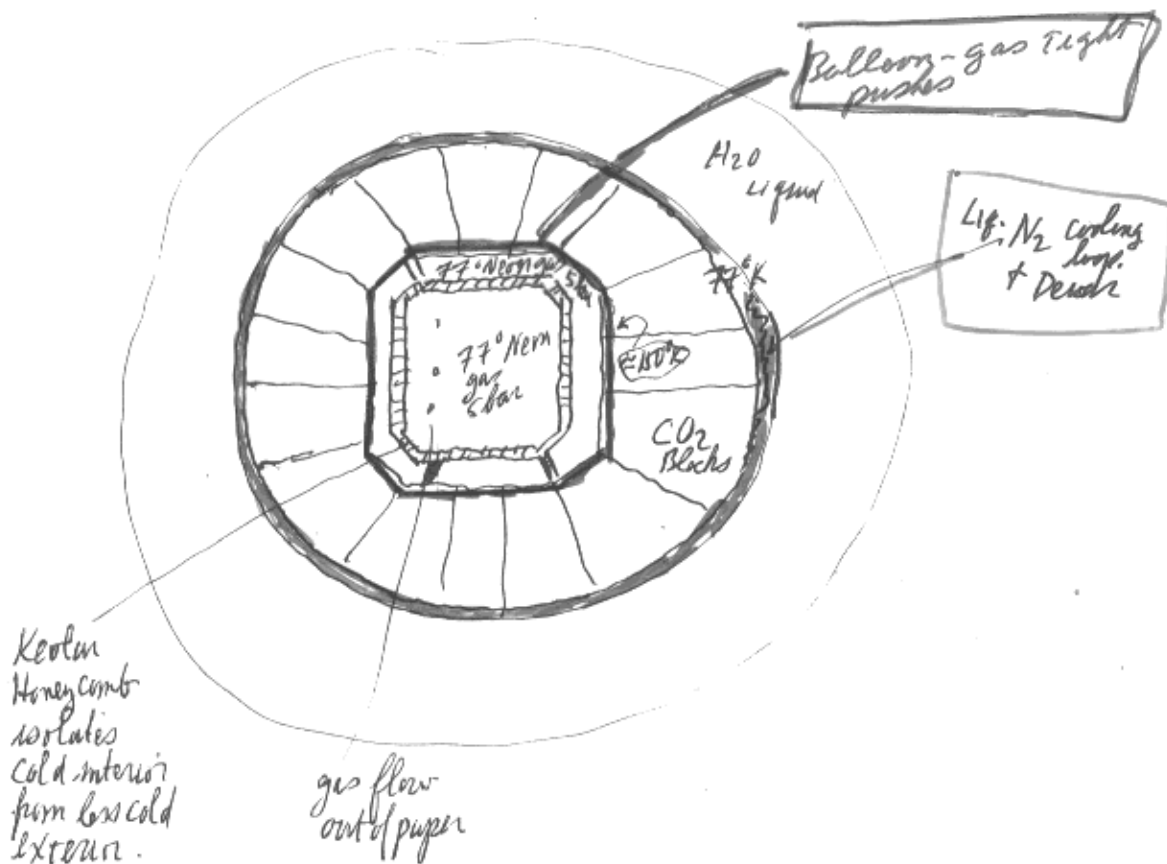


Fig. 14. A transverse view the TPC and shield assembly placed in a Gran Sasso underground experimental hall.

Meeting 9/9/92 with Laurenti, Binocini, Segunit & Ypsilantis
on SND at Bldg 54-R012.

- Decided that best strategy for SND proposal is



- Advantages:
- 77° K Nem gas entered
 - CO₂ blocks take their eventual temp. but because of low thermal conductivity it changes with time
 - Temperature gradient is assured by honeycomb + gas layer between honeycomb + balloon (CO₂ wall)
 - Balloon pushes CO₂ wall outward + assure they stay in place even if wall contracts with temperature.

C.S. CHEMICALLY SEPARABLE

Looks Good if C.S.

BAD

Good if Recent Chemical Separation certainly OK for containment.

BAD

(2)

	52 Cr(0.3%)	53 Cr(0.3%)	54 Cr(0.3%)	55 Cr(0.3%)	56 Mn(0.01%)	57 Fe(0.5%)	58 Fe(0.1%)	59 Fe(0.3%)	60 Co(0.1%)	61 Ni(0.3%)
N, X	52 Cr(0.3%)	53 Cr(0.3%)	54 Cr(0.3%)	55 Cr(0.3%)	56 Mn(0.01%)	57 Fe(0.5%)	58 Fe(0.1%)	59 Fe(0.3%)	60 Co(0.1%)	61 Ni(0.3%)
N, P	52 Cr(0.3%)	53 Cr(0.3%)	54 Cr(0.3%)	55 Cr(0.3%)	56 Mn(0.01%)	57 Fe(0.5%)	58 Fe(0.1%)	59 Fe(0.3%)	60 Co(0.1%)	61 Ni(0.3%)
N, d	52 Cr(0.3%)	53 Cr(0.3%)	54 Cr(0.3%)	55 Cr(0.3%)	56 Mn(0.01%)	57 Fe(0.5%)	58 Fe(0.1%)	59 Fe(0.3%)	60 Co(0.1%)	61 Ni(0.3%)
N, t	52 Cr(0.3%)	53 Cr(0.3%)	54 Cr(0.3%)	55 Cr(0.3%)	56 Mn(0.01%)	57 Fe(0.5%)	58 Fe(0.1%)	59 Fe(0.3%)	60 Co(0.1%)	61 Ni(0.3%)
N, alpha	52 Cr(0.3%)	53 Cr(0.3%)	54 Cr(0.3%)	55 Cr(0.3%)	56 Mn(0.01%)	57 Fe(0.5%)	58 Fe(0.1%)	59 Fe(0.3%)	60 Co(0.1%)	61 Ni(0.3%)
N, 2N	52 Cr(0.3%)	53 Cr(0.3%)	54 Cr(0.3%)	55 Cr(0.3%)	56 Mn(0.01%)	57 Fe(0.5%)	58 Fe(0.1%)	59 Fe(0.3%)	60 Co(0.1%)	61 Ni(0.3%)
N, 3N	52 Cr(0.3%)	53 Cr(0.3%)	54 Cr(0.3%)	55 Cr(0.3%)	56 Mn(0.01%)	57 Fe(0.5%)	58 Fe(0.1%)	59 Fe(0.3%)	60 Co(0.1%)	61 Ni(0.3%)

Nachel OK if C.S. KALLI.B

Good IF C.S.

BAD even with C.S.

	62 Ni(0.1%)	63 Cu(0.2%)	64 Ni(0.1%)	65 Ni(0.1%)	66 Zn(0.01%)	67 Zn(0.1%)	68 Zn(0.1%)	69 Zn(0.1%)	70 Zn(0.1%)
N, X	62 Ni(0.1%)	63 Cu(0.2%)	64 Ni(0.1%)	65 Ni(0.1%)	66 Zn(0.01%)	67 Zn(0.1%)	68 Zn(0.1%)	69 Zn(0.1%)	70 Zn(0.1%)
N, P	62 Ni(0.1%)	63 Cu(0.2%)	64 Ni(0.1%)	65 Ni(0.1%)	66 Zn(0.01%)	67 Zn(0.1%)	68 Zn(0.1%)	69 Zn(0.1%)	70 Zn(0.1%)
N, d	62 Ni(0.1%)	63 Cu(0.2%)	64 Ni(0.1%)	65 Ni(0.1%)	66 Zn(0.01%)	67 Zn(0.1%)	68 Zn(0.1%)	69 Zn(0.1%)	70 Zn(0.1%)
N, t	62 Ni(0.1%)	63 Cu(0.2%)	64 Ni(0.1%)	65 Ni(0.1%)	66 Zn(0.01%)	67 Zn(0.1%)	68 Zn(0.1%)	69 Zn(0.1%)	70 Zn(0.1%)
N, alpha	62 Ni(0.1%)	63 Cu(0.2%)	64 Ni(0.1%)	65 Ni(0.1%)	66 Zn(0.01%)	67 Zn(0.1%)	68 Zn(0.1%)	69 Zn(0.1%)	70 Zn(0.1%)
N, 2N	62 Ni(0.1%)	63 Cu(0.2%)	64 Ni(0.1%)	65 Ni(0.1%)	66 Zn(0.01%)	67 Zn(0.1%)	68 Zn(0.1%)	69 Zn(0.1%)	70 Zn(0.1%)
N, 3N	62 Ni(0.1%)	63 Cu(0.2%)	64 Ni(0.1%)	65 Ni(0.1%)	66 Zn(0.01%)	67 Zn(0.1%)	68 Zn(0.1%)	69 Zn(0.1%)	70 Zn(0.1%)

EXCELLENT BUT MP 400°C

BAD

SEMI-GOOD

	72 Ga(0.1%)	73 Ge(0.1%)	74 Ge(0.1%)	75 As(0.1%)	76 Se(0.1%)	77 Se(0.1%)
N, X	72 Ga(0.1%)	73 Ge(0.1%)	74 Ge(0.1%)	75 As(0.1%)	76 Se(0.1%)	77 Se(0.1%)
N, P	72 Ga(0.1%)	73 Ge(0.1%)	74 Ge(0.1%)	75 As(0.1%)	76 Se(0.1%)	77 Se(0.1%)
N, d	72 Ga(0.1%)	73 Ge(0.1%)	74 Ge(0.1%)	75 As(0.1%)	76 Se(0.1%)	77 Se(0.1%)
N, t	72 Ga(0.1%)	73 Ge(0.1%)	74 Ge(0.1%)	75 As(0.1%)	76 Se(0.1%)	77 Se(0.1%)
N, alpha	72 Ga(0.1%)	73 Ge(0.1%)	74 Ge(0.1%)	75 As(0.1%)	76 Se(0.1%)	77 Se(0.1%)
N, 2N	72 Ga(0.1%)	73 Ge(0.1%)	74 Ge(0.1%)	75 As(0.1%)	76 Se(0.1%)	77 Se(0.1%)
N, 3N	72 Ga(0.1%)	73 Ge(0.1%)	74 Ge(0.1%)	75 As(0.1%)	76 Se(0.1%)	77 Se(0.1%)

GOOD

	78 Se(0.1%)	79 Br(0.1%)	80 Kr(0.1%)	81 Kr(0.1%)	82 Kr(0.1%)	83 Kr(0.1%)	84 Kr(0.1%)	85 Kr(0.1%)
N, X	78 Se(0.1%)	79 Br(0.1%)	80 Kr(0.1%)	81 Kr(0.1%)	82 Kr(0.1%)	83 Kr(0.1%)	84 Kr(0.1%)	85 Kr(0.1%)
N, P	78 Se(0.1%)	79 Br(0.1%)	80 Kr(0.1%)	81 Kr(0.1%)	82 Kr(0.1%)	83 Kr(0.1%)	84 Kr(0.1%)	85 Kr(0.1%)
N, d	78 Se(0.1%)	79 Br(0.1%)	80 Kr(0.1%)	81 Kr(0.1%)	82 Kr(0.1%)	83 Kr(0.1%)	84 Kr(0.1%)	85 Kr(0.1%)
N, t	78 Se(0.1%)	79 Br(0.1%)	80 Kr(0.1%)	81 Kr(0.1%)	82 Kr(0.1%)	83 Kr(0.1%)	84 Kr(0.1%)	85 Kr(0.1%)
N, alpha	78 Se(0.1%)	79 Br(0.1%)	80 Kr(0.1%)	81 Kr(0.1%)	82 Kr(0.1%)	83 Kr(0.1%)	84 Kr(0.1%)	85 Kr(0.1%)
N, 2N	78 Se(0.1%)	79 Br(0.1%)	80 Kr(0.1%)	81 Kr(0.1%)	82 Kr(0.1%)	83 Kr(0.1%)	84 Kr(0.1%)	85 Kr(0.1%)
N, 3N	78 Se(0.1%)	79 Br(0.1%)	80 Kr(0.1%)	81 Kr(0.1%)	82 Kr(0.1%)	83 Kr(0.1%)	84 Kr(0.1%)	85 Kr(0.1%)

Radimetric (M) 10^-10, 10^-11

BAD

BAD

	85 Rb(0.1%)	86 Sr(0.1%)	87 Sr(0.1%)	88 Sr(0.1%)	89 Y(0.1%)	90 Zr(0.1%)	91 Zr(0.1%)	92 Zr(0.1%)	93 Zr(0.1%)
N, X	85 Rb(0.1%)	86 Sr(0.1%)	87 Sr(0.1%)	88 Sr(0.1%)	89 Y(0.1%)	90 Zr(0.1%)	91 Zr(0.1%)	92 Zr(0.1%)	93 Zr(0.1%)
N, P	85 Rb(0.1%)	86 Sr(0.1%)	87 Sr(0.1%)	88 Sr(0.1%)	89 Y(0.1%)	90 Zr(0.1%)	91 Zr(0.1%)	92 Zr(0.1%)	93 Zr(0.1%)
N, d	85 Rb(0.1%)	86 Sr(0.1%)	87 Sr(0.1%)	88 Sr(0.1%)	89 Y(0.1%)	90 Zr(0.1%)	91 Zr(0.1%)	92 Zr(0.1%)	93 Zr(0.1%)
N, t	85 Rb(0.1%)	86 Sr(0.1%)	87 Sr(0.1%)	88 Sr(0.1%)	89 Y(0.1%)	90 Zr(0.1%)	91 Zr(0.1%)	92 Zr(0.1%)	93 Zr(0.1%)
N, alpha	85 Rb(0.1%)	86 Sr(0.1%)	87 Sr(0.1%)	88 Sr(0.1%)	89 Y(0.1%)	90 Zr(0.1%)	91 Zr(0.1%)	92 Zr(0.1%)	93 Zr(0.1%)
N, 2N	85 Rb(0.1%)	86 Sr(0.1%)	87 Sr(0.1%)	88 Sr(0.1%)	89 Y(0.1%)	90 Zr(0.1%)	91 Zr(0.1%)	92 Zr(0.1%)	93 Zr(0.1%)
N, 3N	85 Rb(0.1%)	86 Sr(0.1%)	87 Sr(0.1%)	88 Sr(0.1%)	89 Y(0.1%)	90 Zr(0.1%)	91 Zr(0.1%)	92 Zr(0.1%)	93 Zr(0.1%)

The Hellaz Neutrino Detector

T. Ypsilantis

Collège de France, Paris

Simulations show that unlike other existing and proposed detectors, a track projection chamber based on 2000 cubic metres of a cold helium-methane gas mixture can help solve the solar neutrino problem by determining the direction and energy of solar neutrinos.

The most interesting solar neutrinos are pp neutrinos, with a continuum extending up to an energy of 420 keV, and the monoenergetic 862 keV ${}^7\text{Be}$ neutrinos formed in the $p+p \rightarrow e+d+\nu_e$ and $e'+{}^7\text{Be} \rightarrow {}^7\text{Li}+\nu_e$ reactions. All neutrinos have a known flavour (electron-neutrino ν_e) at the sun and energy production in the sun mostly transits through the first reaction, so the pp neutrino flux is essentially determined by the solar luminosity. Since the neutrino-electron (ν_e) elastic scattering cross-section is given by the well tested electroweak theory, any observed differences in flux, flavour or spectral shape with respect to theoretical predictions will require either new physics (such as flavour oscillations whereby one type of neutrino mixes into another type) or a major revamping of the Standard Solar Model that describes the Sun's nuclear physics and hydrodynamics. This is the so-called solar neutrino problem (see insert for a discussion of the latest situation).

A new solar neutrino experiment called Hellaz (Helium at liquid azote - nitrogen - temperature) is being proposed to help resolve the problem using pp solar neutrino data. It consists of 6 tonnes of cold helium gas as a target for neutrino-electron elastic scatters. A recoil electron from a ν_e event leaves a trail of ionization electrons which define a track. Each ionization electron drifts in an applied axial electric field until it reaches a two-dimensional x,y-detector located at each end of a cylindrical target volume (Fig. 1). Following the track projection chamber (TPC) method, the z-coordinate is obtained from the electron arrival time. The x,y,z data determine the recoil electron energy and direction and, given the sun's position, the neutrino energy E_ν .

Tom Ypsilantis, a research physicist in high-energy particle physics based at the Collège de France in Paris, is currently working on neutrino experiments (Hellaz and Long Baseline RICH) and on applications of the RICH technique to B-physics at CERN's future LHC collider. He has worked at the Ecole Polytechnique, Paris, the CE Saclay, CERN, and the University of Southern California, Berkeley.

Hellaz will measure recoil electrons with a kinetic energy T as low as 100 keV corresponding to a range of about 50 mm because it is filled with gases of low atomic number and low density. Simulations demonstrate that T is measured to within 3% and the electron direction to within 35 mrad so the energy of incoming 300 keV neutrinos can be determined to within 2 to 4% - an advantage of Hellaz in contrast to existing and proposed ν_e detectors (see insert). The proposed Borexino can detect recoil electrons with energies above 260 keV, but cannot determine the electron direction, and hence the neutrino energy.

Moreover, the cross-sections $d\sigma/dT$ for $\nu_e e$ electron-neutrino and $\nu_e e$ gamma-neutrino elastic scattering on electrons are sufficiently different to allow a sensitive determination of the flavour and flux of incoming neutrinos independently of solar models or mixing scenarios. Indeed, Hellaz will be able to determine the amount of

flavour mixing for pp and pep neutrinos to permit resolution of all the current neutrino mixing scenarios.

Hellaz is also competitive when compared to other much heavier detectors. The calculated rate of ν_e elastic events from the pp and ${}^7\text{Be}$ solar neutrino fluxes is about 16 events/day for 6 tonnes of helium gas if all recoil electrons ($T \geq 0$ keV) are detected. With the more realistic limit of $T \geq 100$ keV there will still be 11 events/day. By comparison, the Gran Sasso Underground Laboratory's GaCl_3 -based GALLEX detector with the same threshold and twice the mass only detects 0.2 events/day, since although the cross-section of ${}^{71}\text{Ga}$ is 4.5 times the average ν_e cross-section, its nuclear density is only 1/88 of the helium electron density.

Several prototypes of the Hellaz detector have been constructed and are now under test. A 40 m³ device will soon be placed in a tunnel at Gran Sasso with full carbon dioxide (or B_2O_3) shielding to confirm estimates of the radiopurity of materials by neutron activation analysis and to measure the ν_e signal and the gamma-electron (γe) background. It will also serve an ideal tool for searching for dark matter, an aspect that will not be discussed.

Direction and Energy Resolution

The precision with which the direction and energy of the incident neutrinos can be determined in Hellaz has been estimated using accurate simulations of

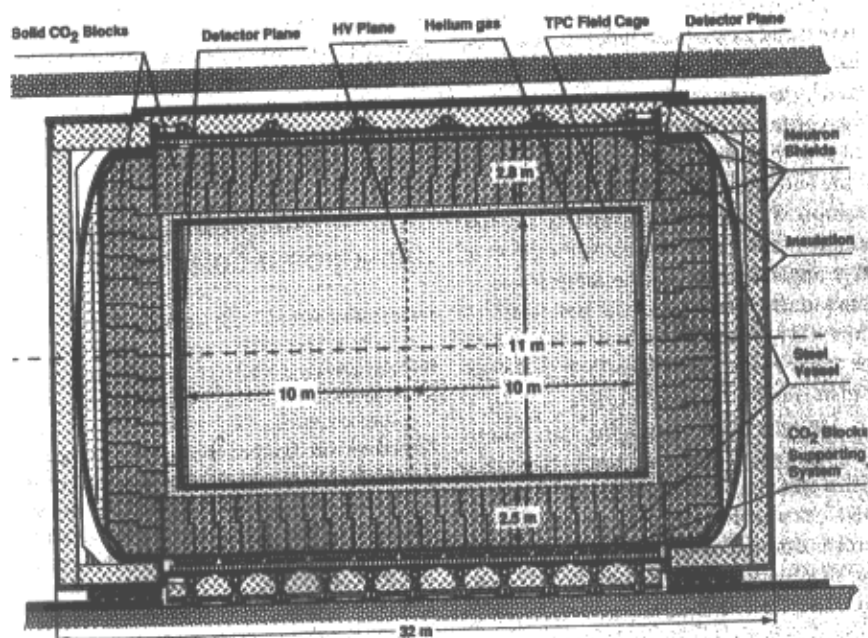


Fig. 1. A longitudinal view of the heavily shielded, cylindrical Hellaz time proportional counter solar-neutrino detector. Ionization tracks drift in an axial field generated by a field cage to modular multiwire detectors located at the endcaps which use the same low-temperature helium-methane gas mixture as the target volume.

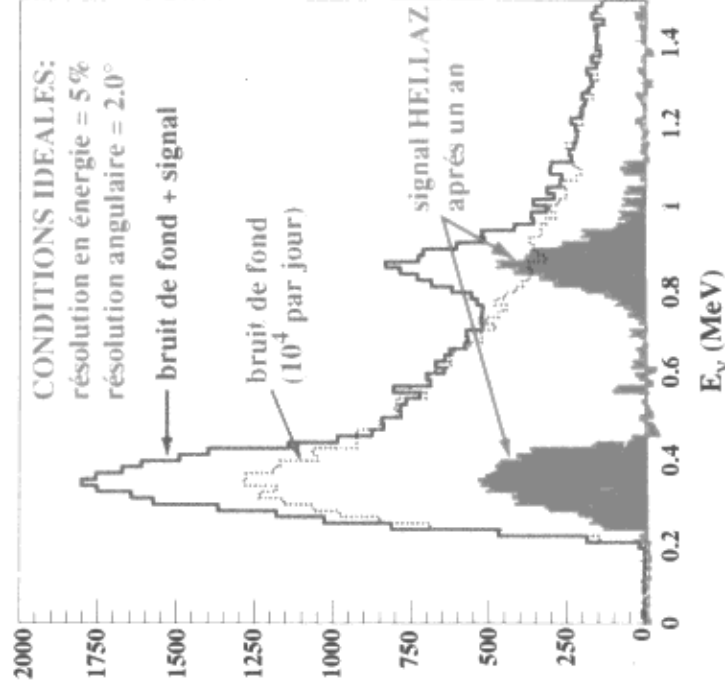
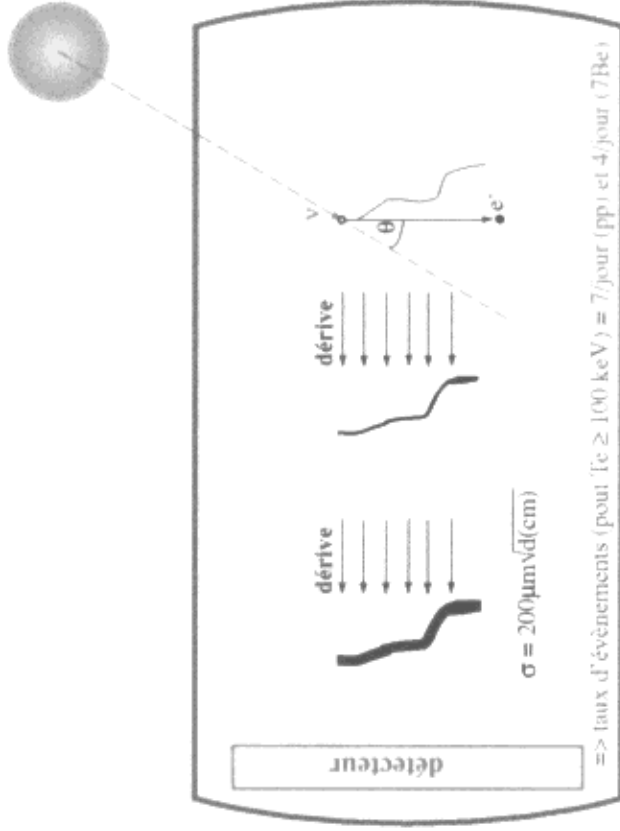
Principe de HELLAZ:

diffusion élastique neutrino - électron dans He - TPC

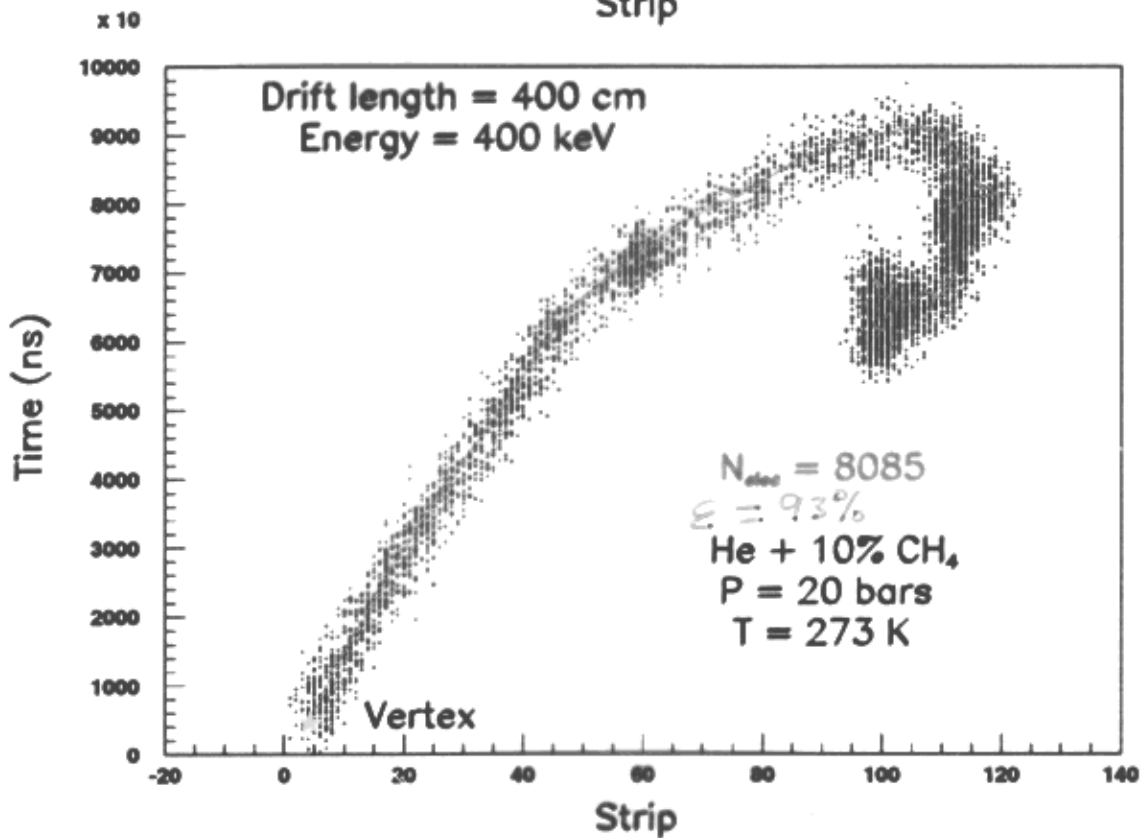
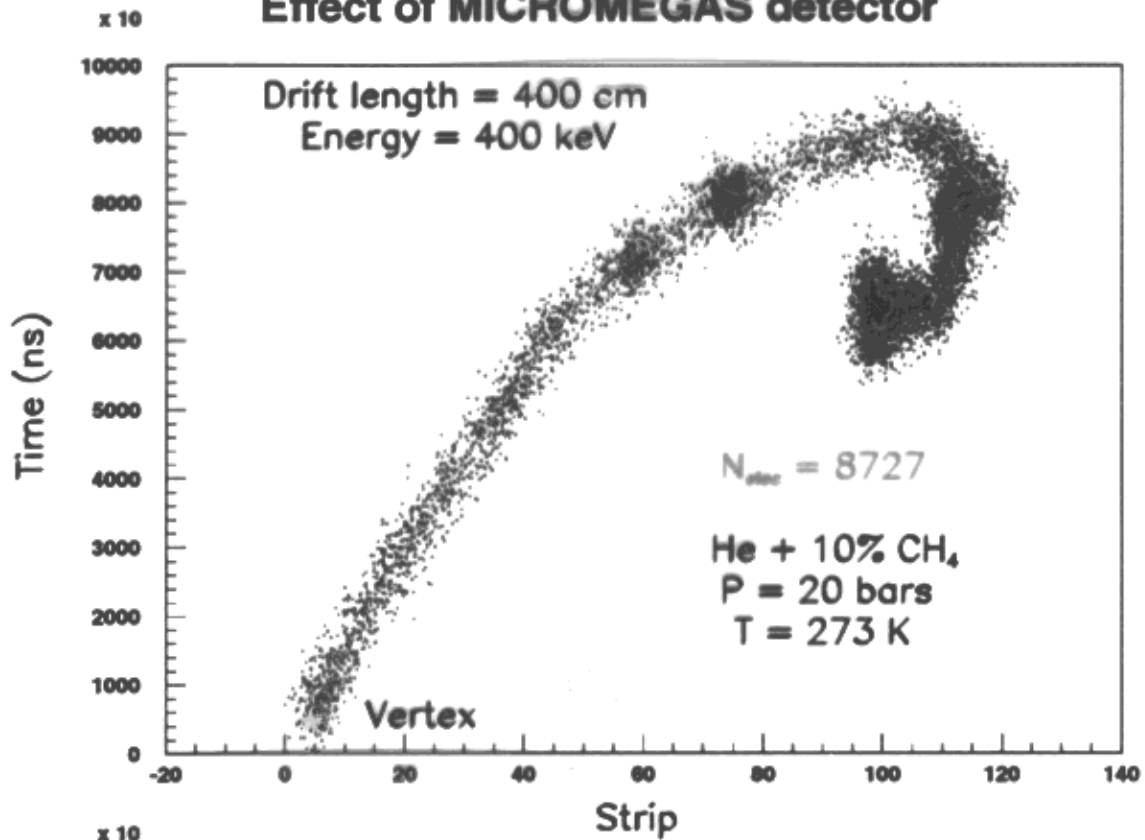
$\rho_{He} = 3.124 \times 10^{-3} \text{ g/cm}^3$, 6 tonnes He $\Rightarrow N_e = 2 \cdot 10^{30}$

• mesure de T_e ($\sigma_T/T \approx 3\%$) et θ ($\sigma_\theta \approx 35 \text{ mrad}$)

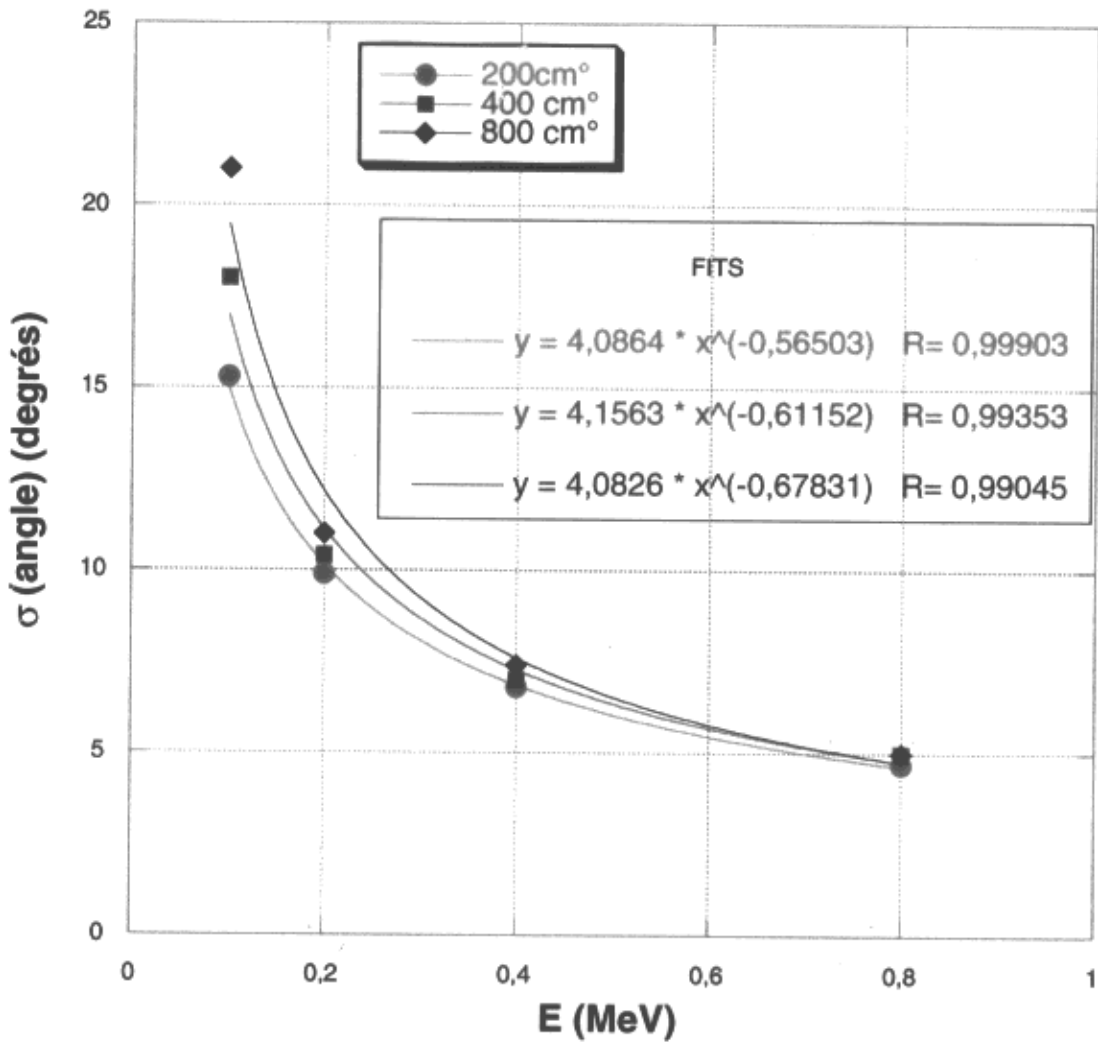
• Énergie du neutrino $E_\nu = \frac{m_e}{\sqrt{1 + \frac{2m_e}{T_e} \cos\theta}}$



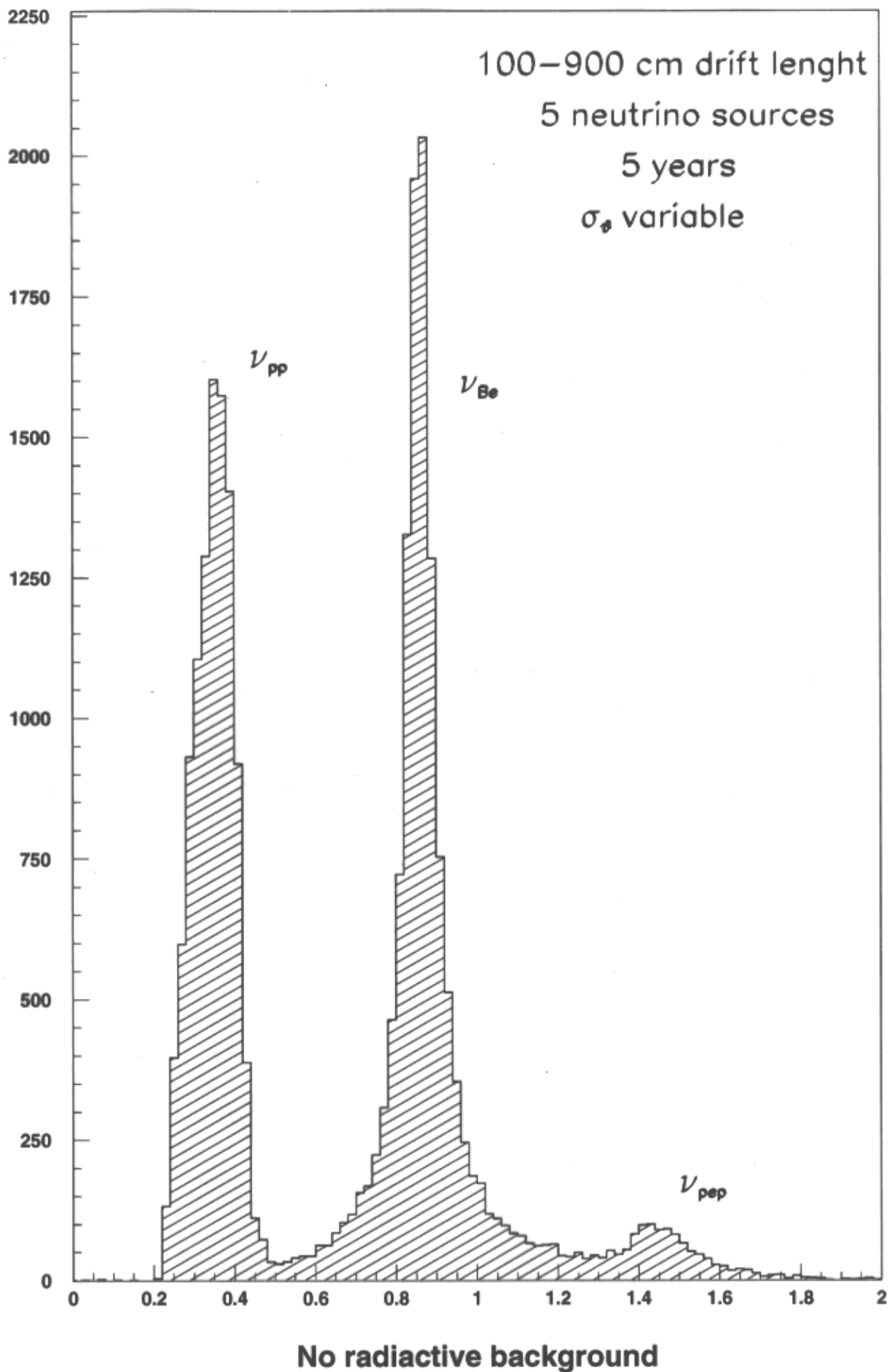
Effect of MICROMEGAS detector



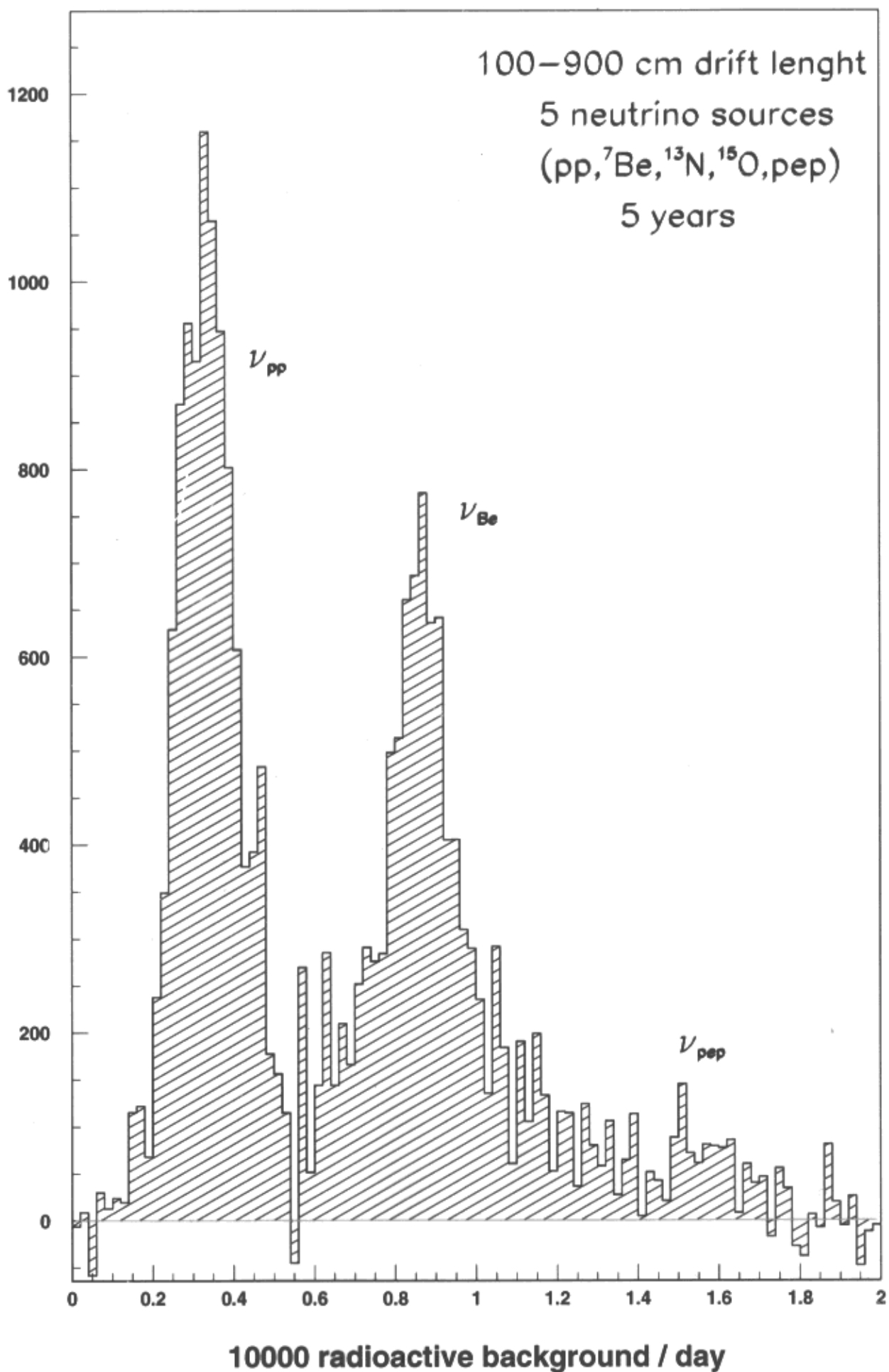
Dolbeau 2/2/01 data



Reconstruct neutrino spectrum

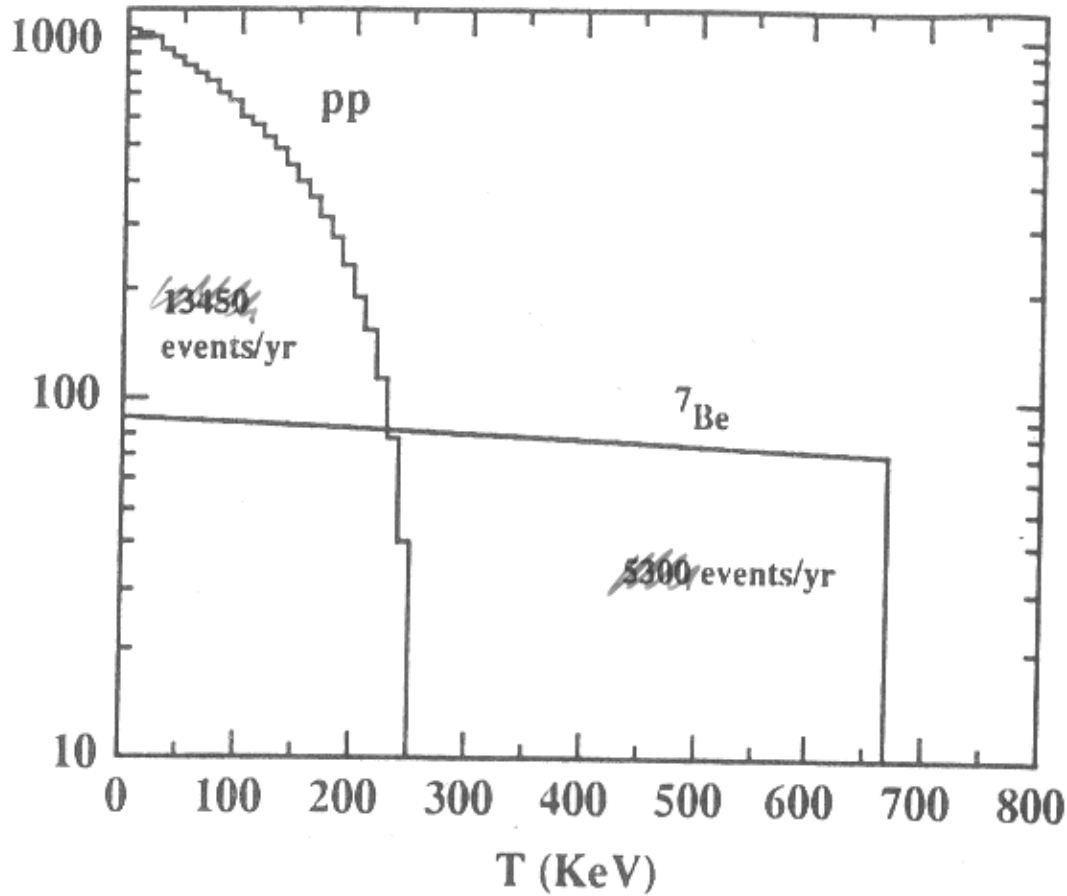


Reconstruct neutrino spectrum



EVENTS / 10 KeV-yr

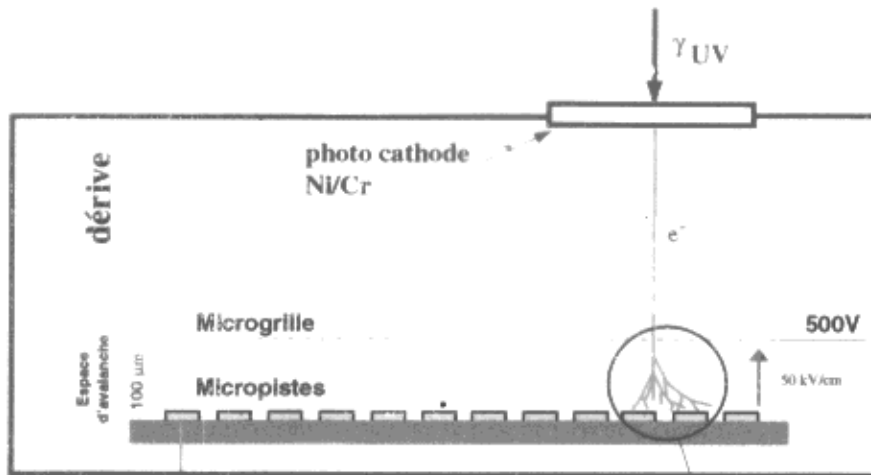
no oscillations



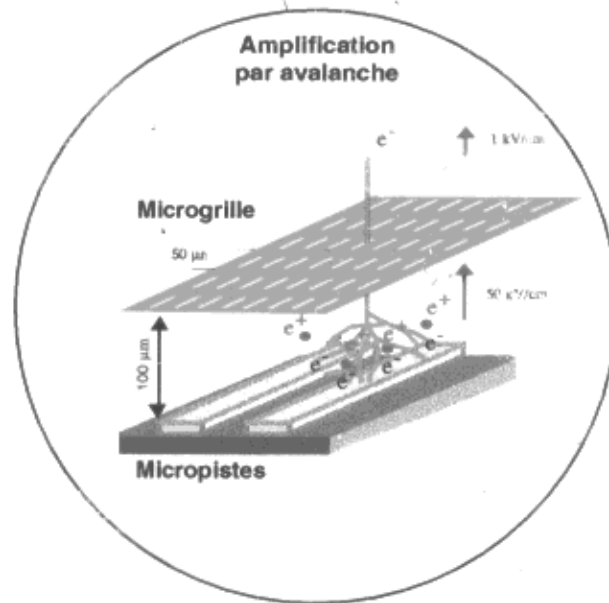
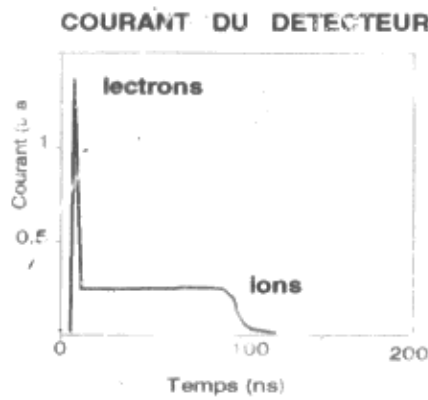
Comparison between E_ν and T_e spectra:

- Lines (${}^7\text{Be}$): obvious (ratio of surfaces)
- β spectra (pp, ${}^8\text{B}$): not so different in shape

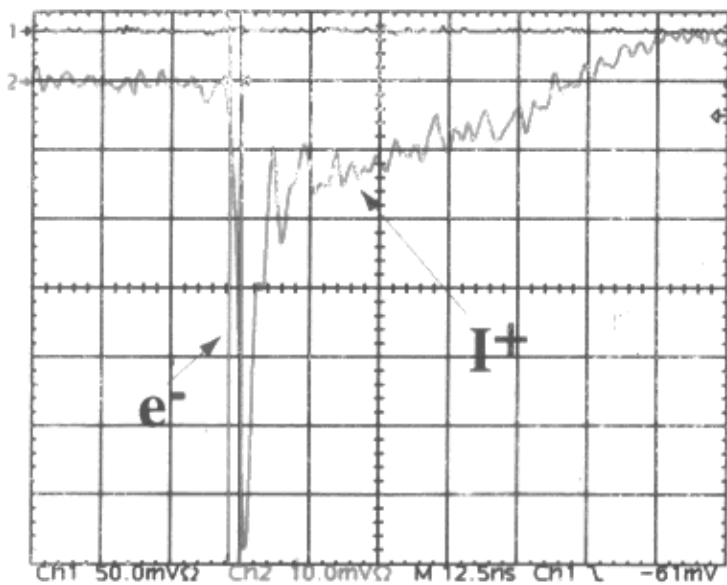
HOWEVER: if oscillations: cannot be seen in T_e .



Amplificateurs
Lecture des pistes

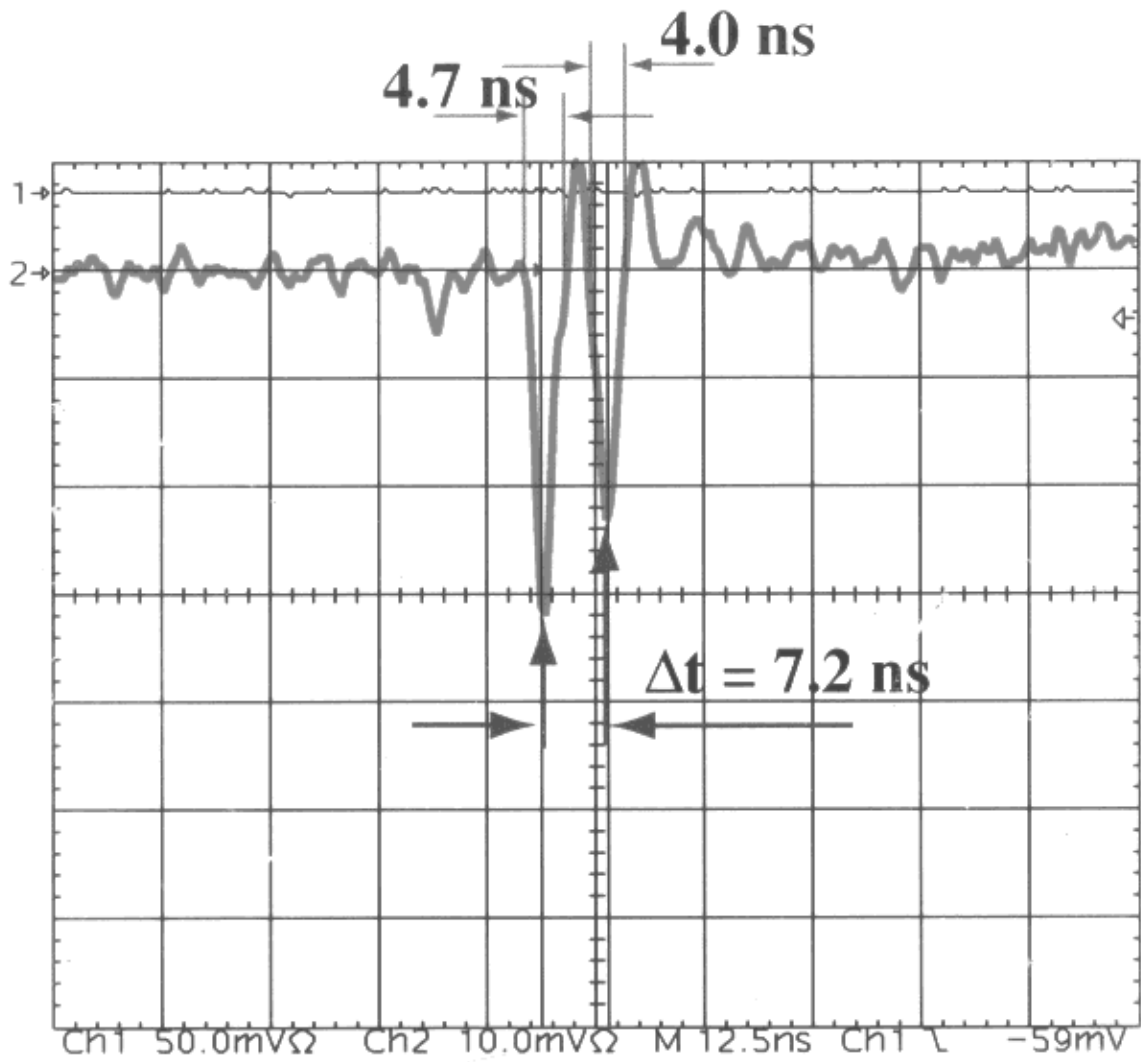


mesure au PCC:

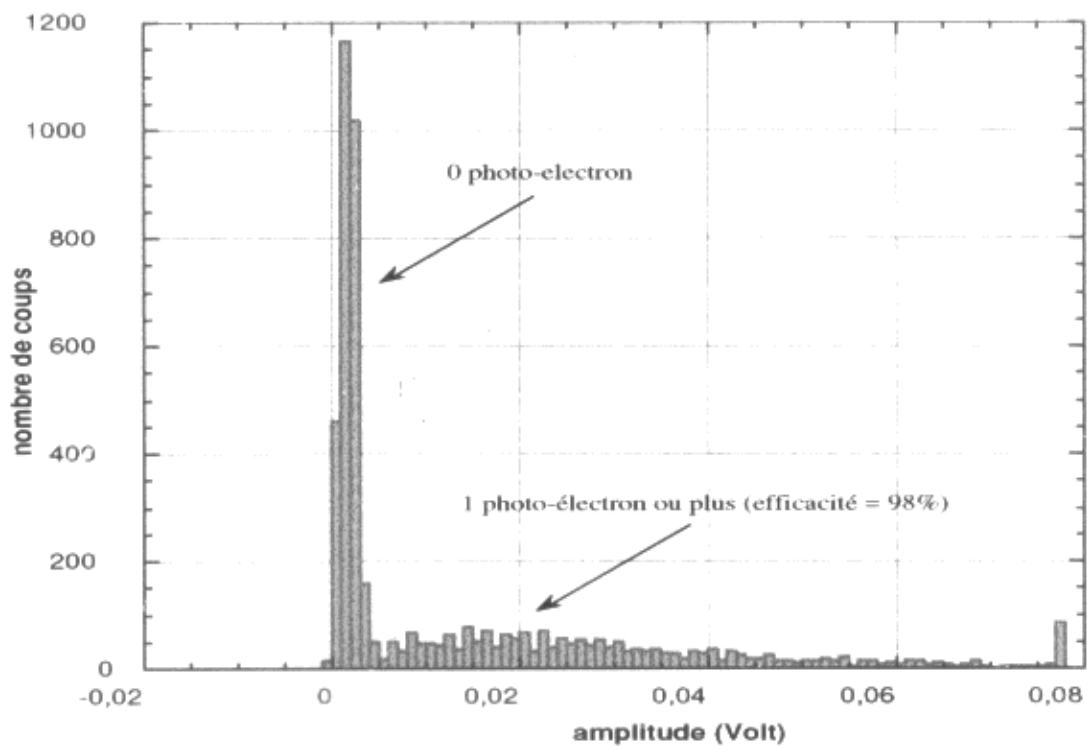


- 100 % He + 6 % Isobutan ($P_{abs} = 1 \text{ bar}$, $T = 310 \text{ K}$)
- distance photo cathode - mesh = 8.5 mm
- distance mesh - cathode = 100 μm
- ampli = ORTEC 120
- source = laser + 1 quartz fibre
- HT(pc) = -1900 V
- HT(mesh) = -570 V

- signal of a single electron
- fall time = 1.80 ns
- $\approx 6.9 \text{ ns}$ on the base of the signal



What about Raether limit? Gain $\approx 10^6$ Rate $\approx 10^8$!

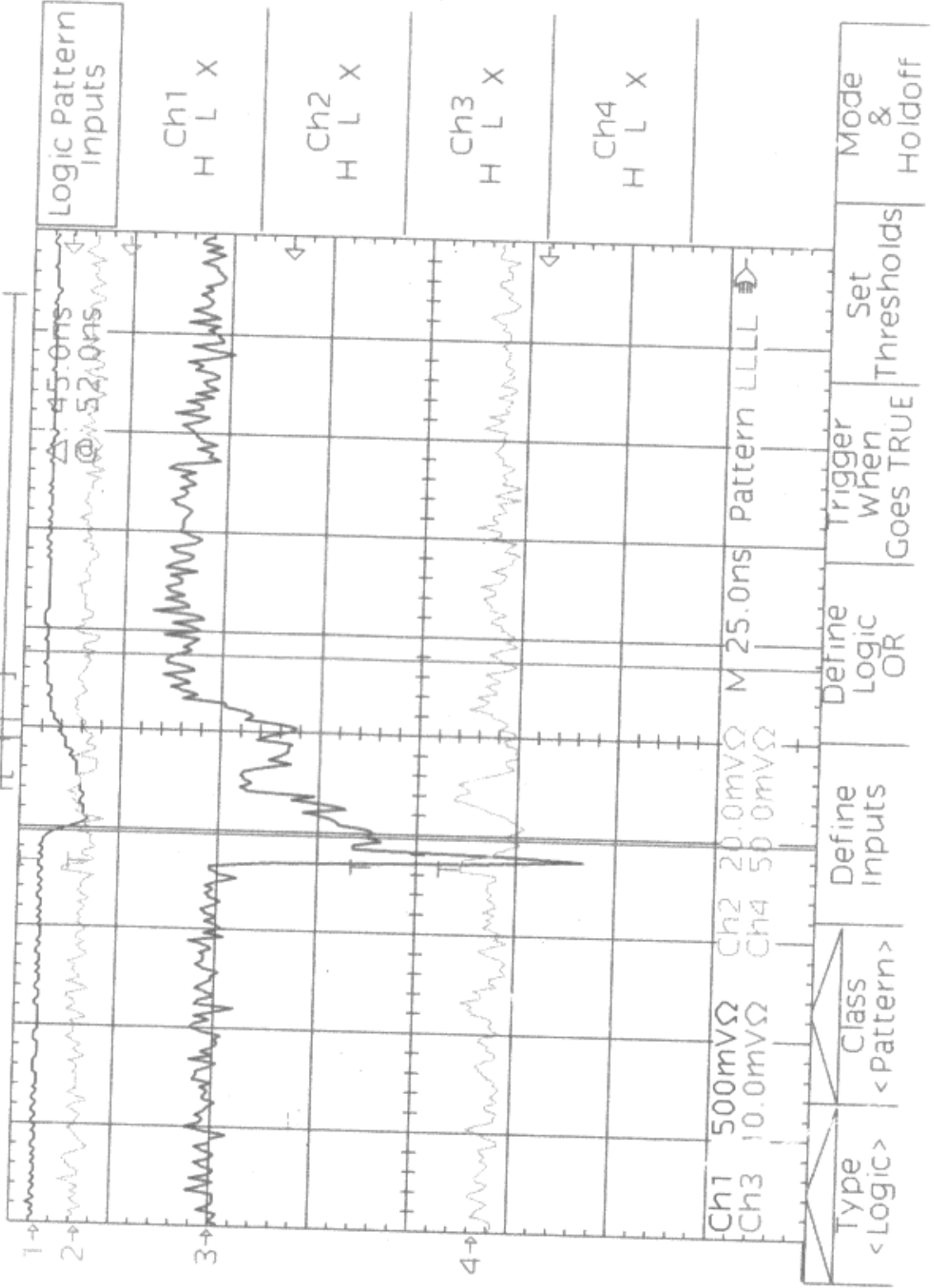


8 octobre 1999

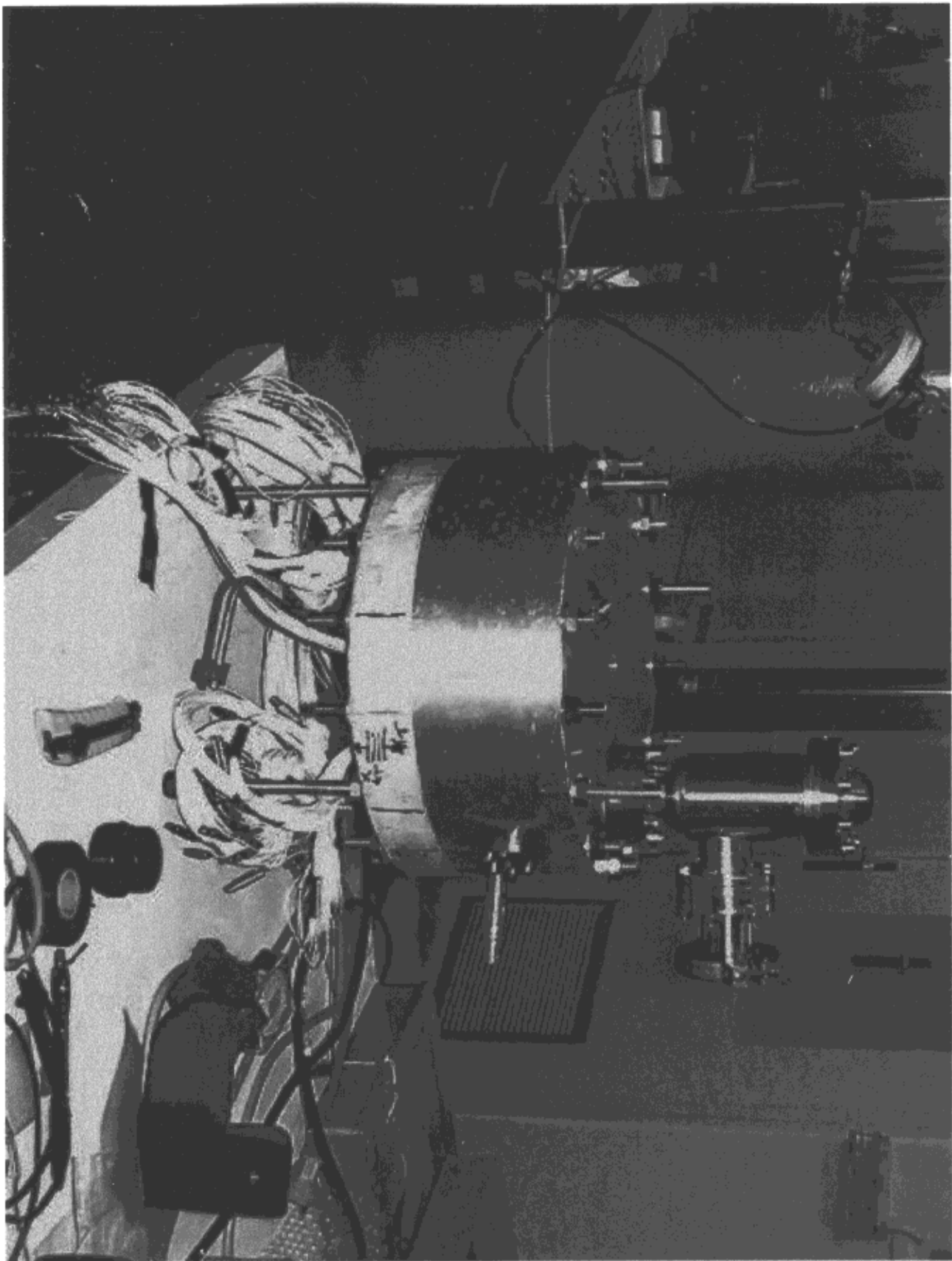
19.94 bar
800 V
100% He
6% CH4

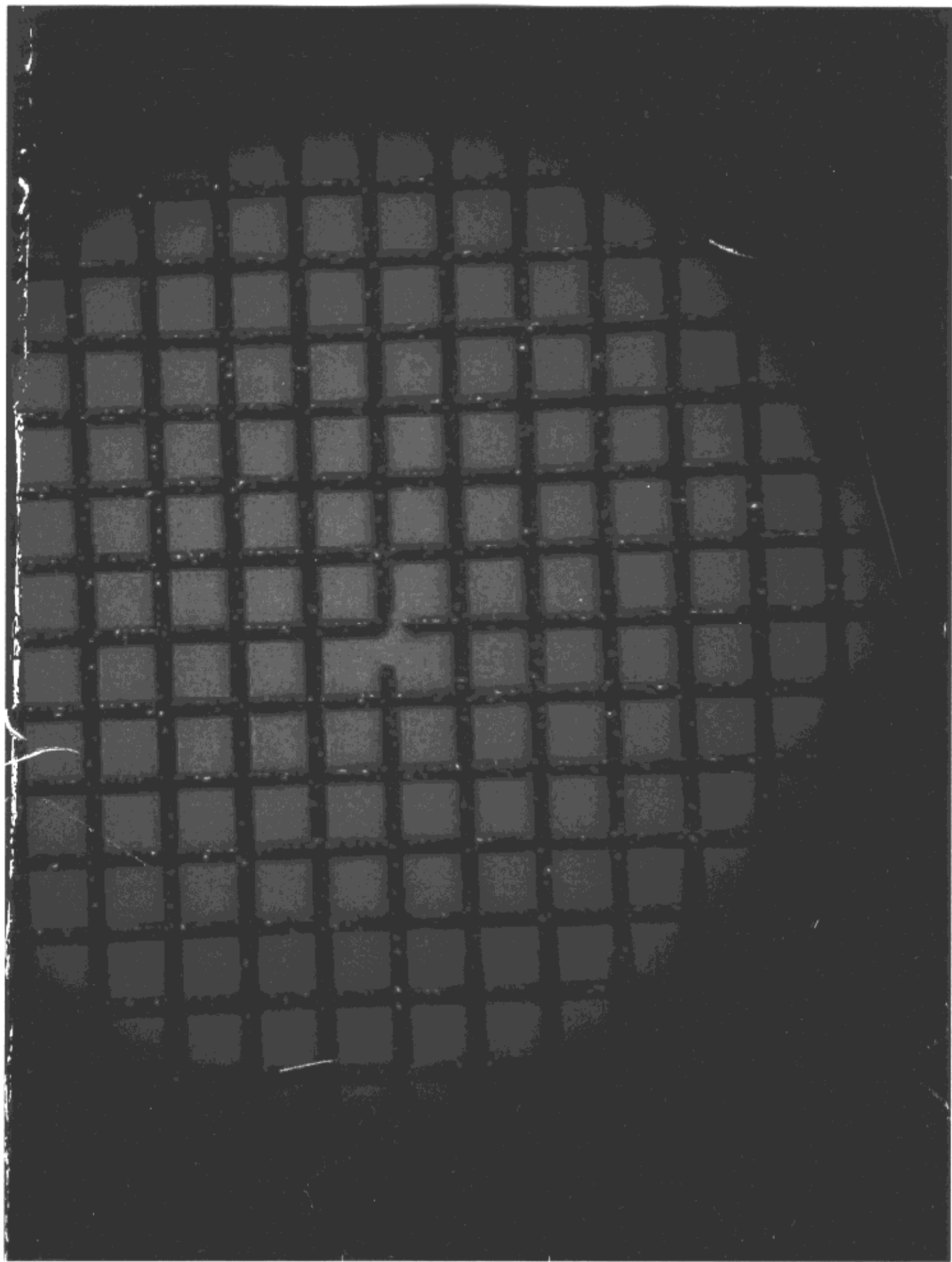
Tek Stop: 1.00GS/s

274 Acqs



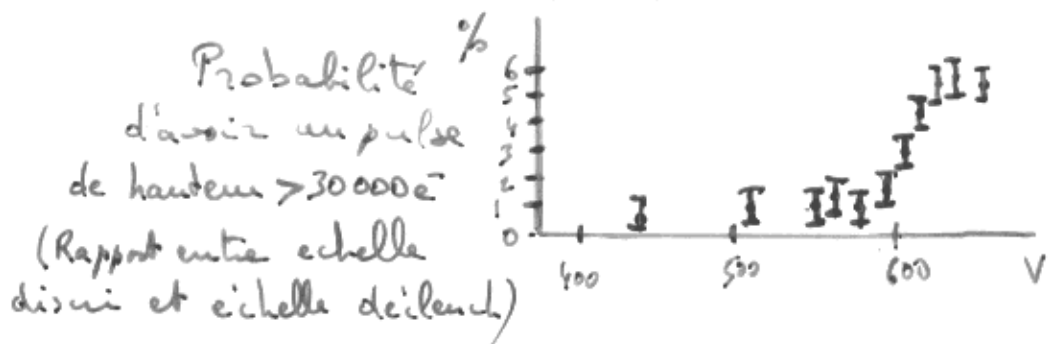
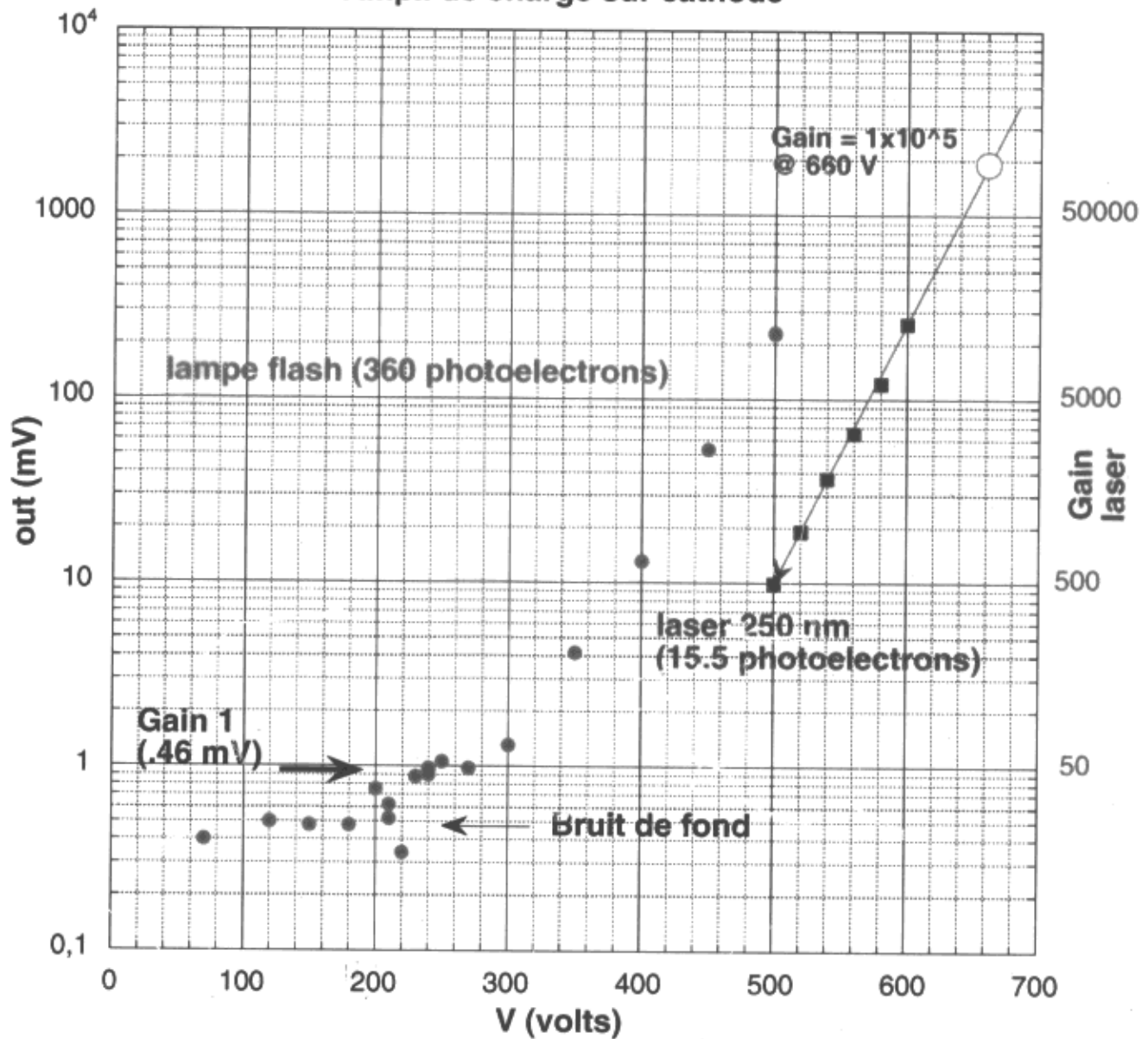
100 kF

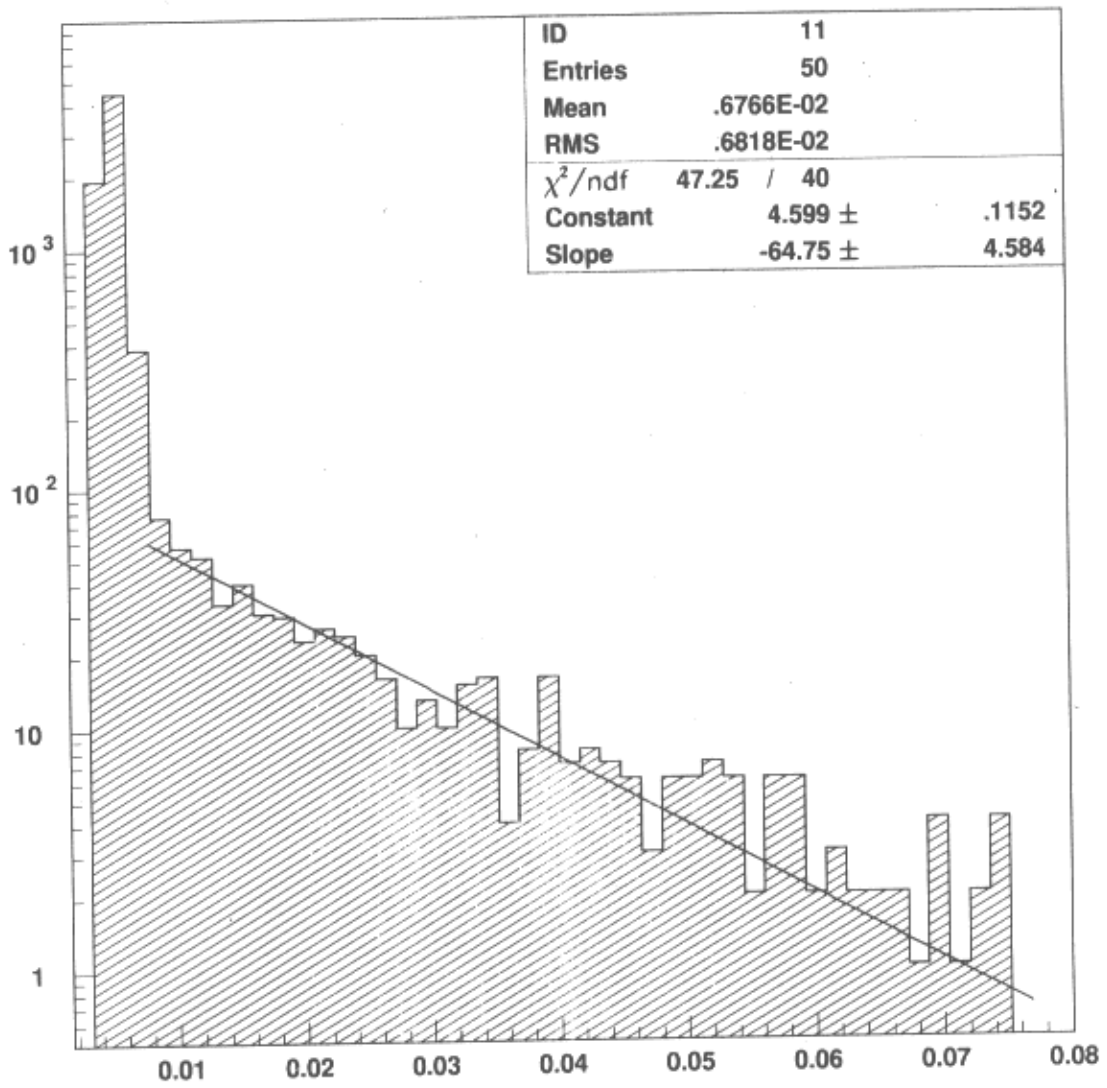




Micromegas 50 μ m
 97.5% He, 2% CH₄, 0.5% isobutane 20 bar

Ampli de charge sur cathode





Have signed the LOI

Wayne State (7 people)

Arizona (2 people)

FSU Tallahassee (5 people)

BNL (2 people)

College de France (4 people)

Have declared an interest:

LBL Berkeley

Chicago

Indiana Cyclotron

New Mexico

Pittsburgh

Alberta (Pinfold)

CEA: S. Turck-Chieze

Discussions with:

Japan

Bologna

Neuchatel

HELLAZ, March 2001