

# Is cosmology compatible with Sterile (light) neutrinos ?

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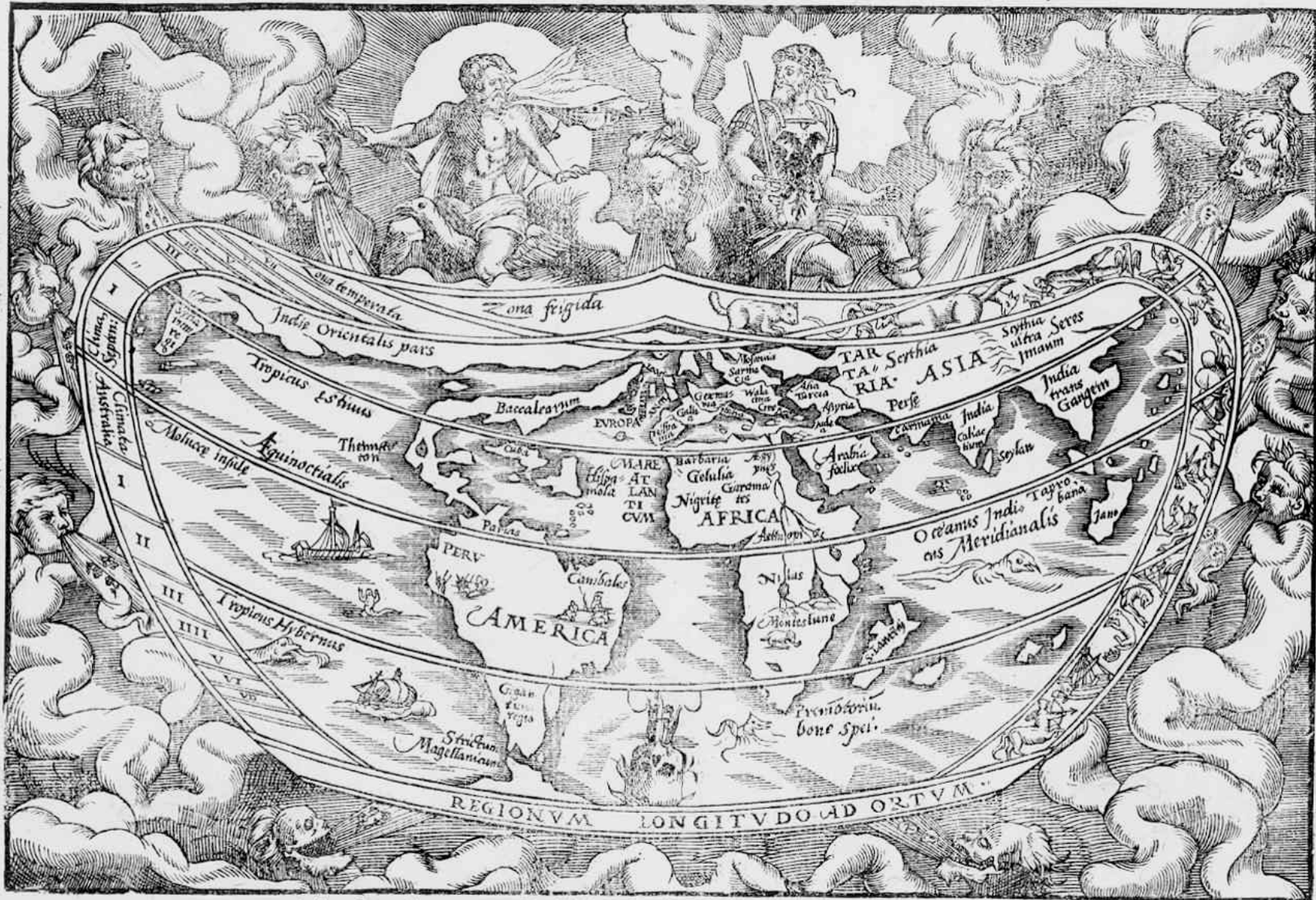
Venezia, 8th March, 2007



CHARTA COSMOGRAPHICA, CVM VENTORVM PROPRIA NATVRA ET OPERATIONE.  
 Circius, Noorde noordwest. SEPTENTRIONALIS, Noordt. I RIO Aquilo, Noorde noorde oost.

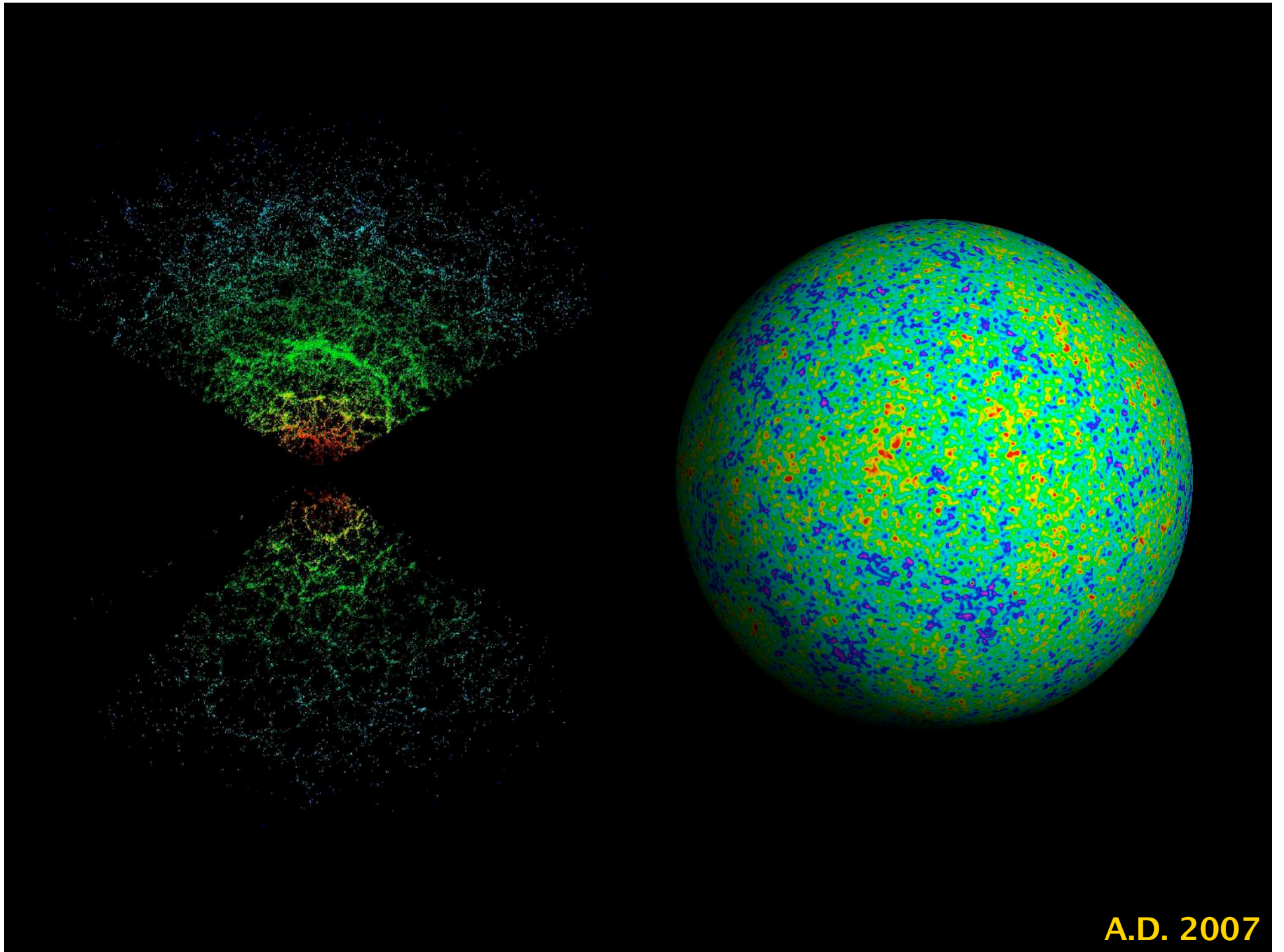
Argellus, West noordwest.  
 OCCI Zephyrus DENS, West.  
 Libs West zuidwest.

Hellepontius Oost noordooft.  
 ORI Subolanus ENS, Oost.  
 Vulturius, Oost zuidooft.



Austroafricus, Zuidt zuidwest. MELI Auster, Zuidt DIE S. Euroauster, Zuidt zuidooft. I

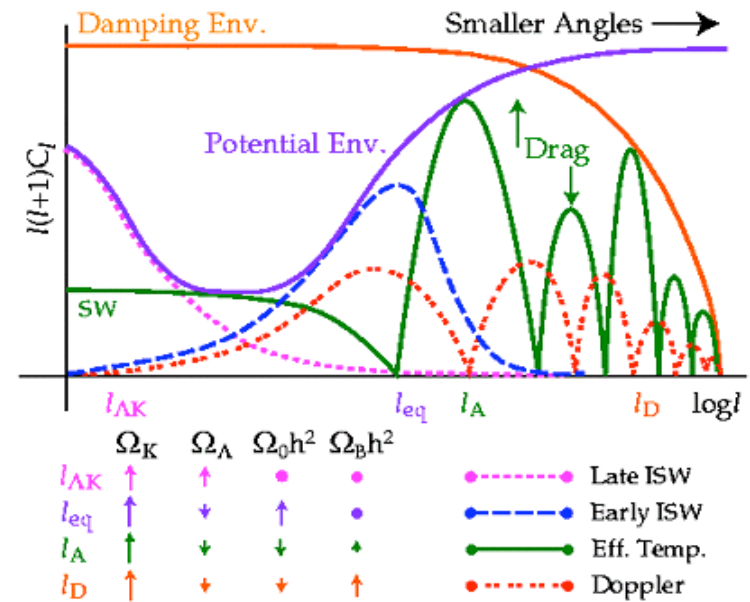
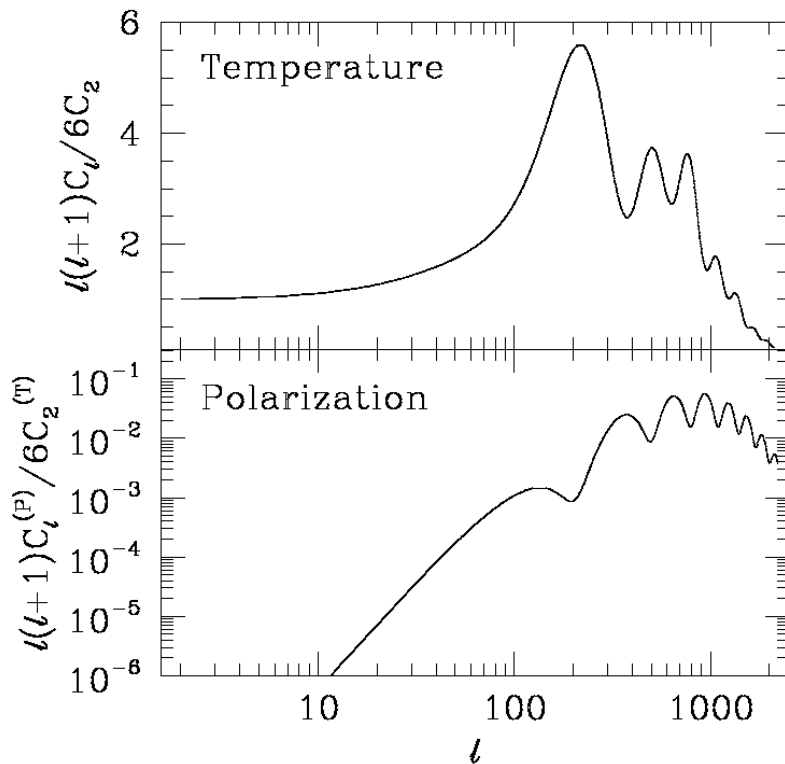
A.D. 1544



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# CMB: Theory

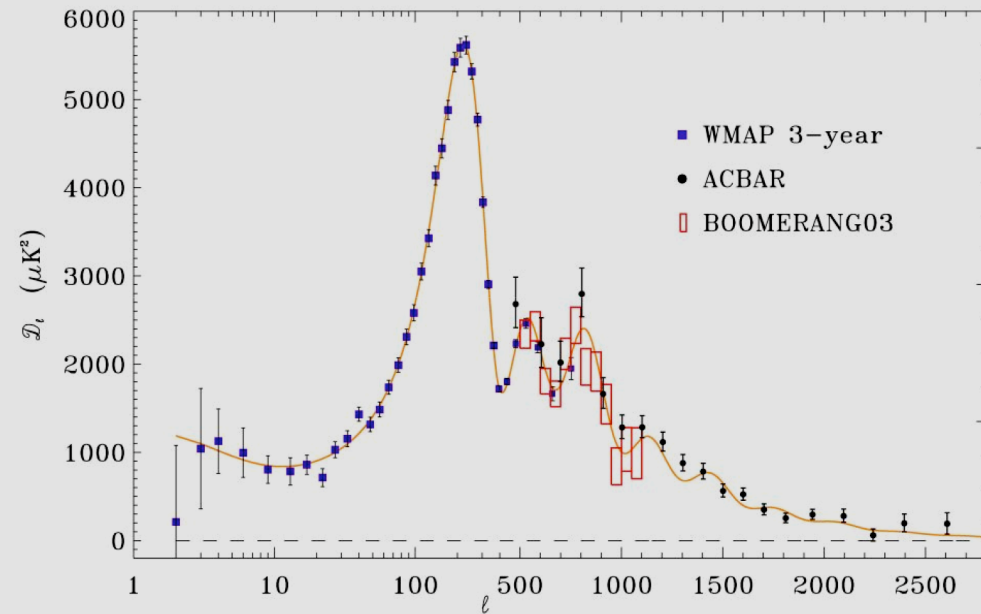
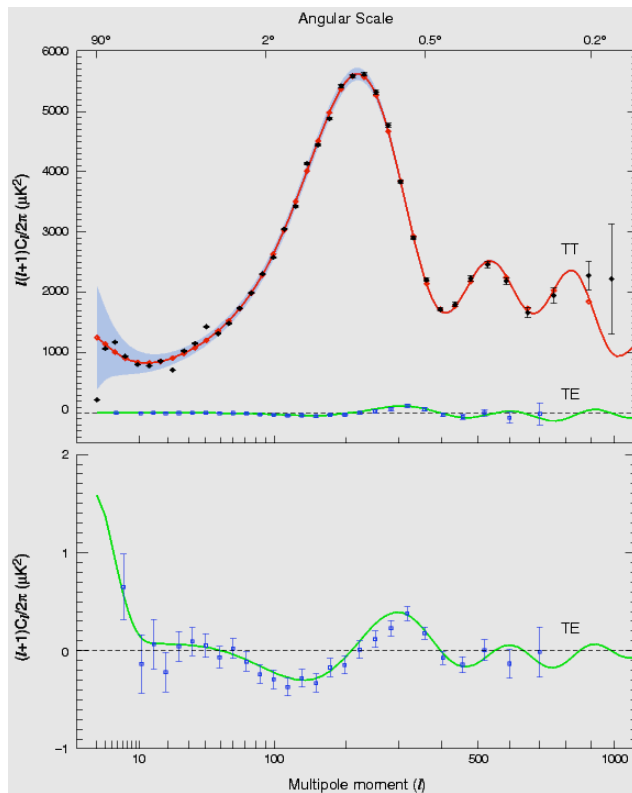
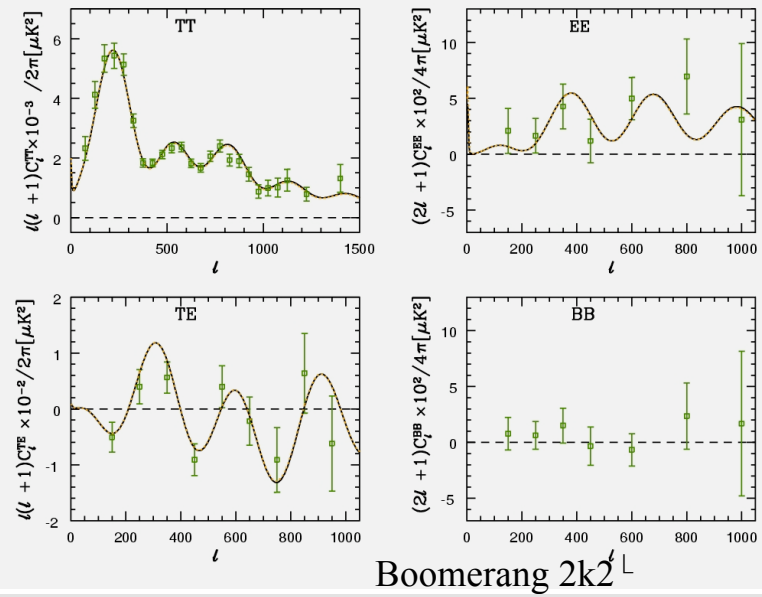
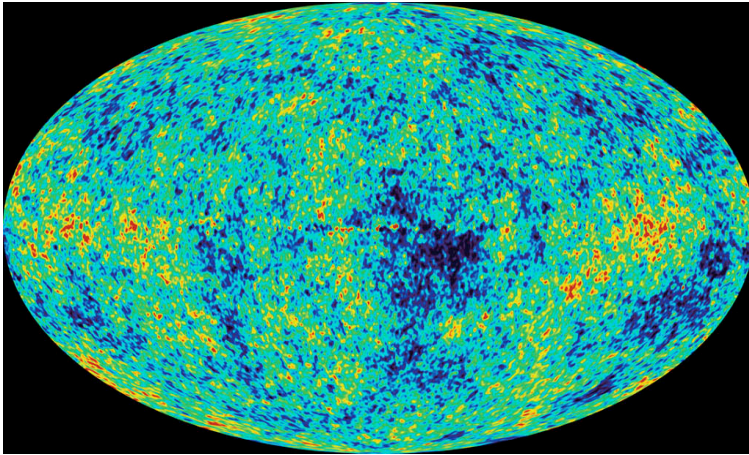
$$\left\langle \frac{\Delta T}{T}(\hat{n}_1) \frac{\Delta T}{T}(\hat{n}_2) \right\rangle = \frac{1}{2\pi} \sum_{\ell} (2\ell + 1) C_{\ell} P_{\ell}(\hat{n}_1 \cdot \hat{n}_2)$$



1/(Angular Scale)



# CMB: Data

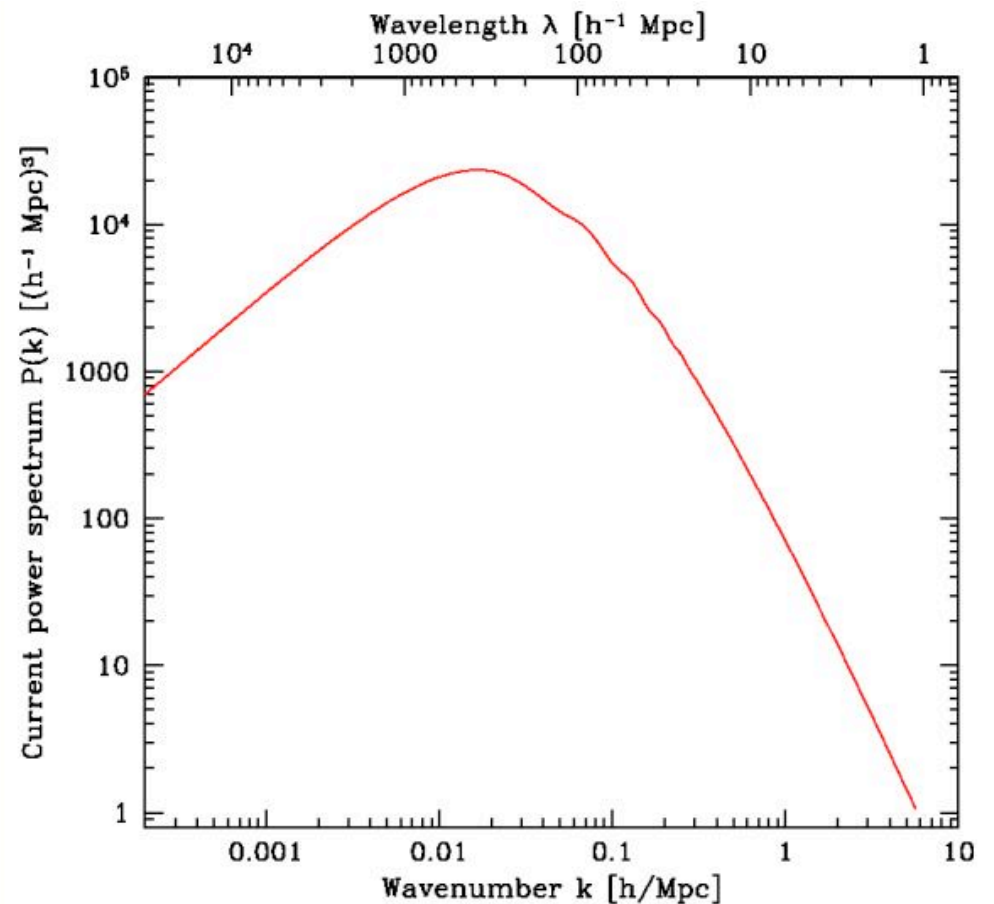
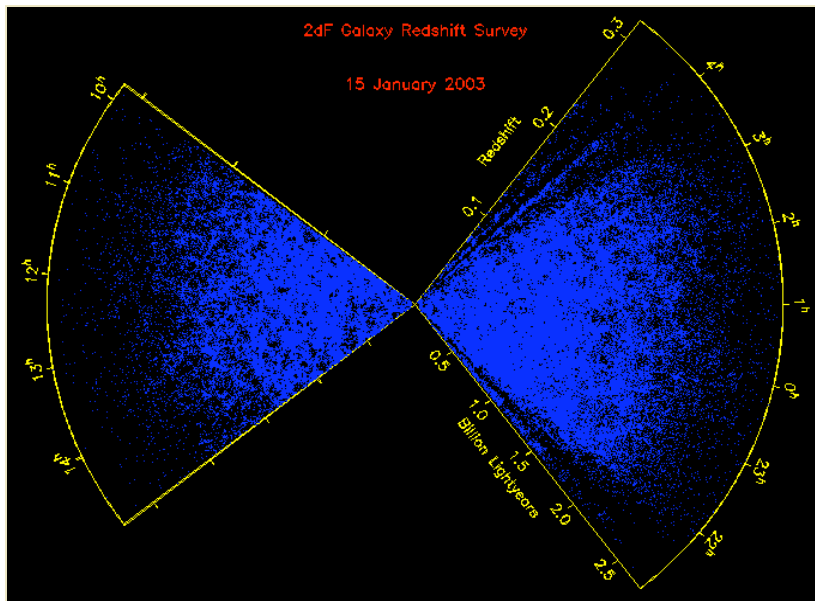


# Galaxy Clustering: Theory

$$\bar{\rho}(r, t) = \langle \bar{\rho}(\vec{x}, t) \bar{\rho}(\vec{x} + \vec{r}, t) \rangle$$

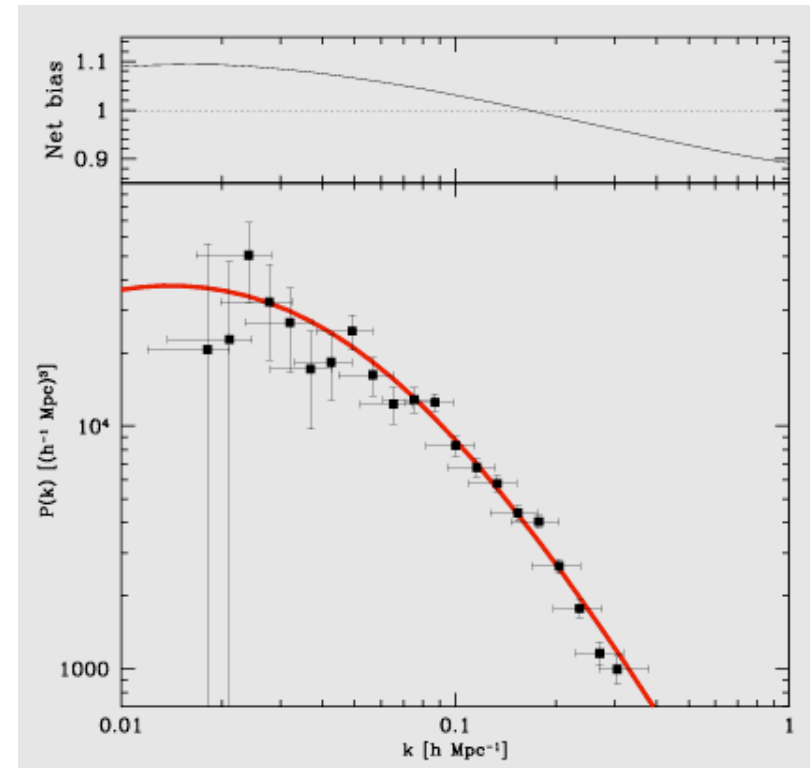
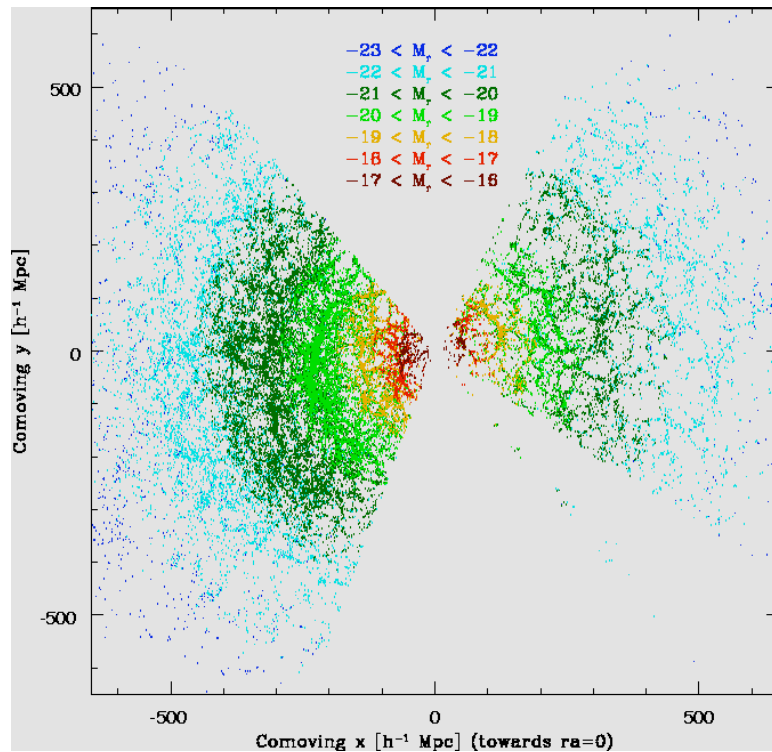
$$\bar{\rho}_{\text{galaxies}}(r, t) = b^2 \bar{\rho}(r, t)$$

$$P(k, t) = \int d^3r \bar{\rho}(r, t) e^{i\vec{k} \cdot \vec{r}}$$





# Galaxy Clustering: Data



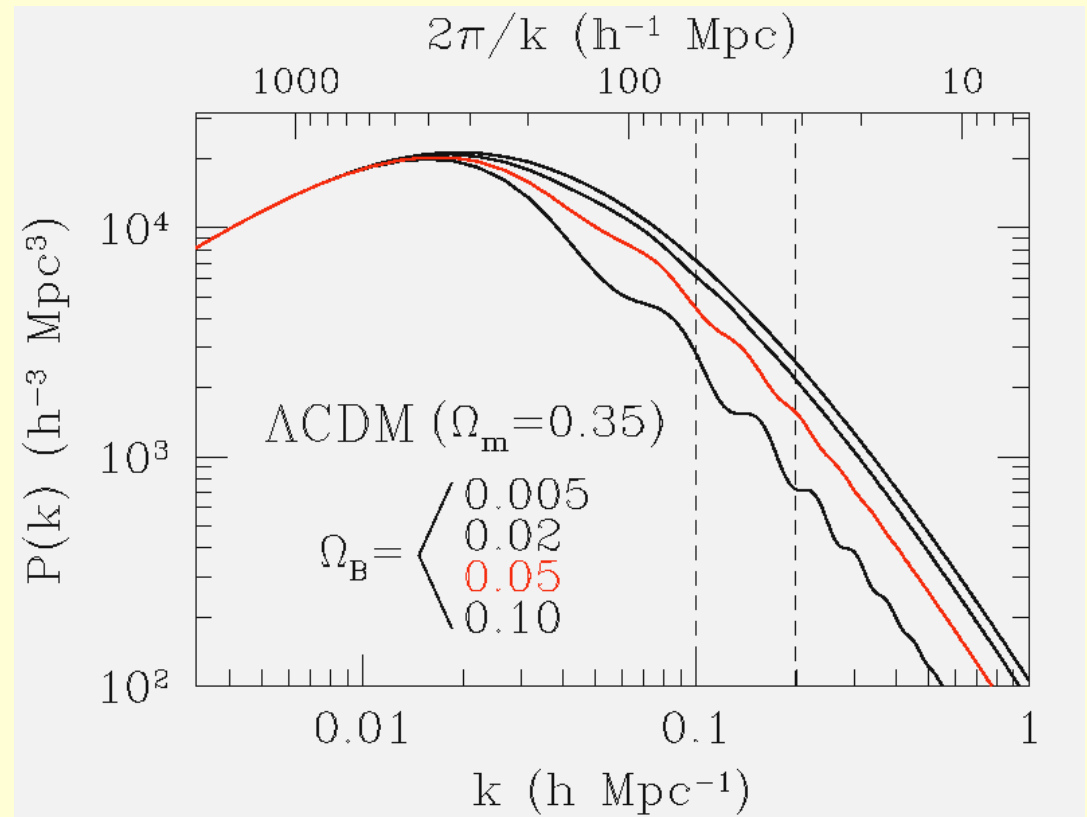
Again, perfect agreement with (low density) L-CDM model...

# LSS as a cosmic yardstick

Imprint of oscillations  
less clear in LSS  
spectrum unless high  
baryon density

Detection much more  
difficult:

- Survey geometry
- Non-linear effects
- Biasing



Big pay-off:

Potentially measure  $d_A(z)$  at many redshifts!

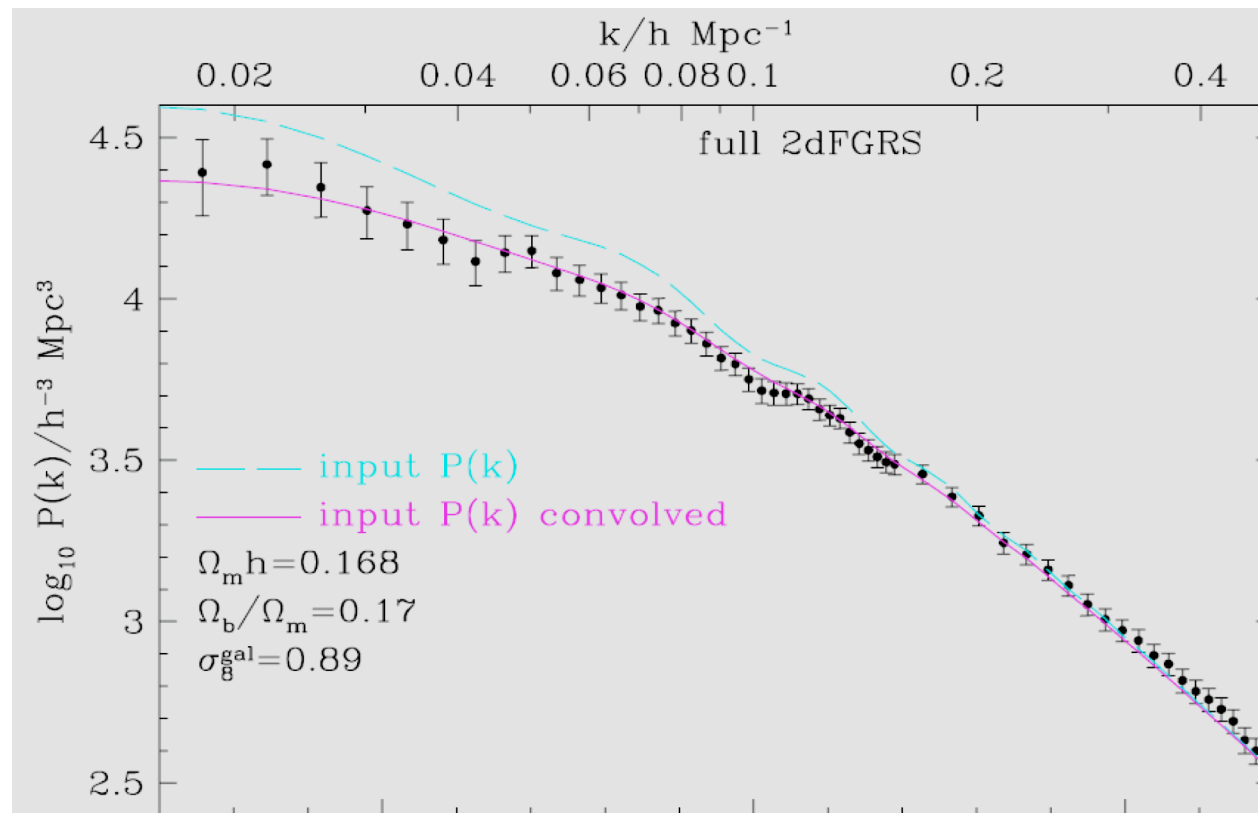


# Recent detections of the baryonic signature

- ◆ Cole et al
  - 221,414 galaxies,  $b_J < 19.45$
  - (final 2dFGRS catalogue)
- ◆ Eisenstein et al
  - 46,748 luminous red galaxies (LRGs)
  - (from the Sloan Digital Sky Survey)

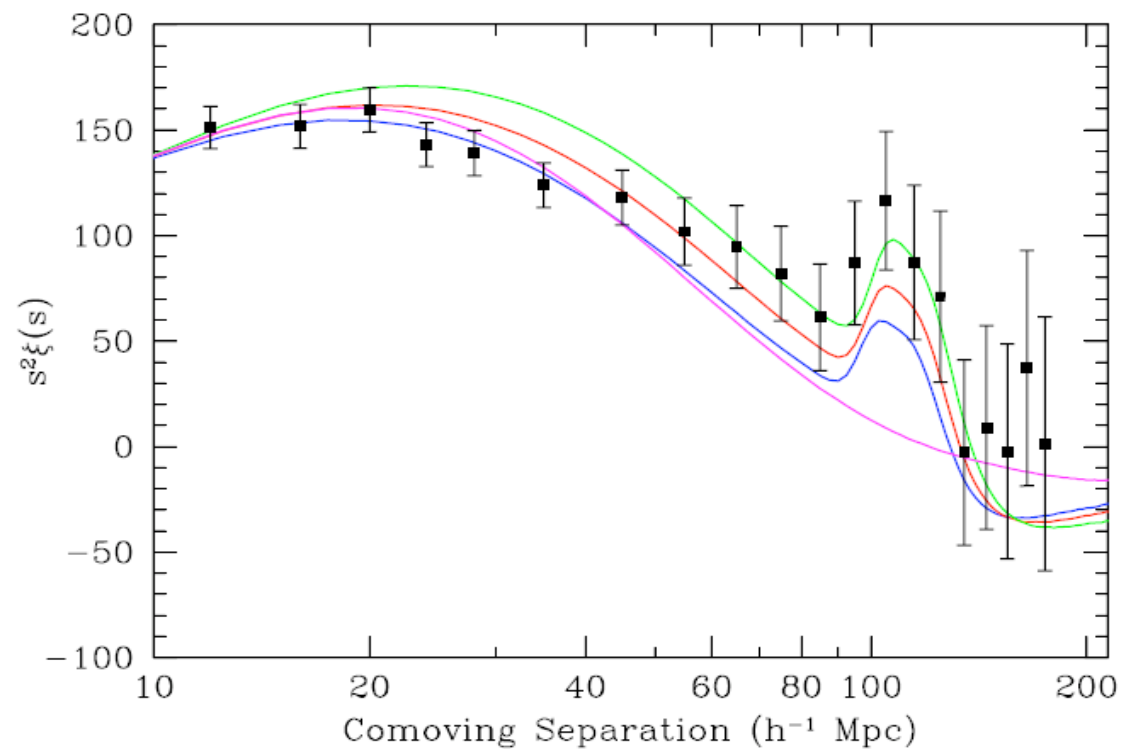


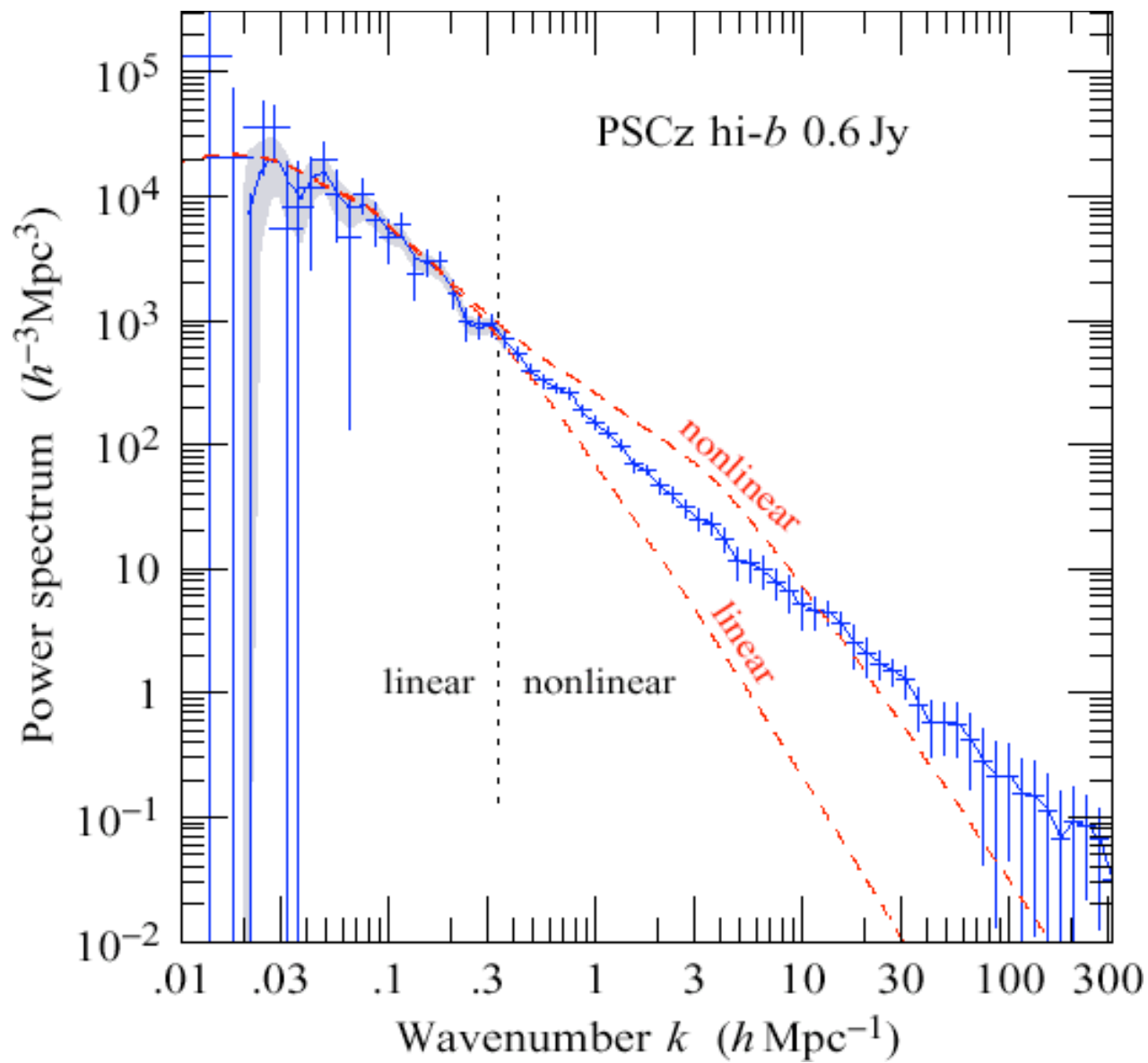
# The 2dFGRS power spectrum





# The SDSS LRG correlation function

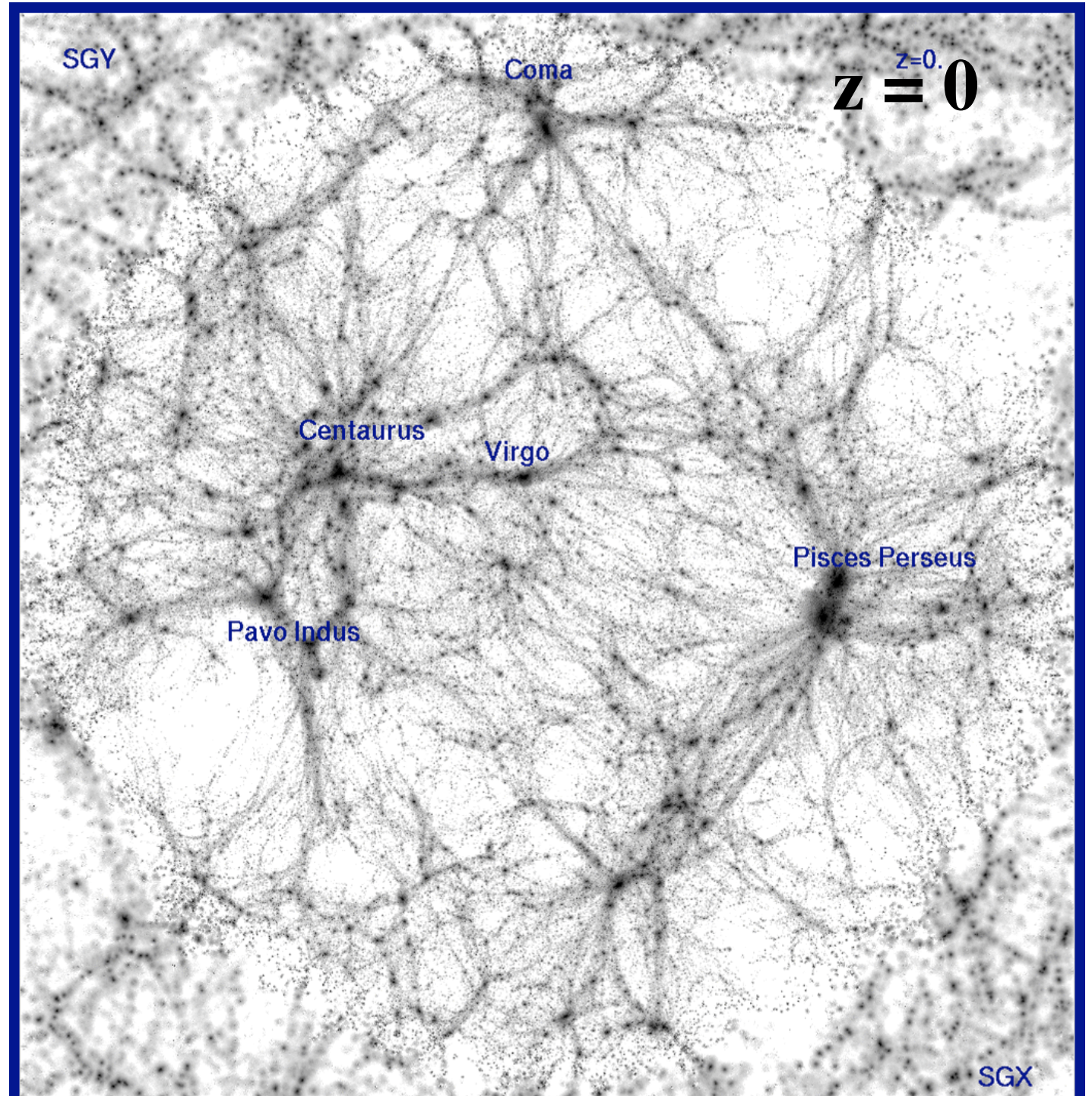




We want to go to  
smaller scales!!!  
(and be linear)



Mathis, Lemson, Springel, Kauffmann, White & Dekel 2001

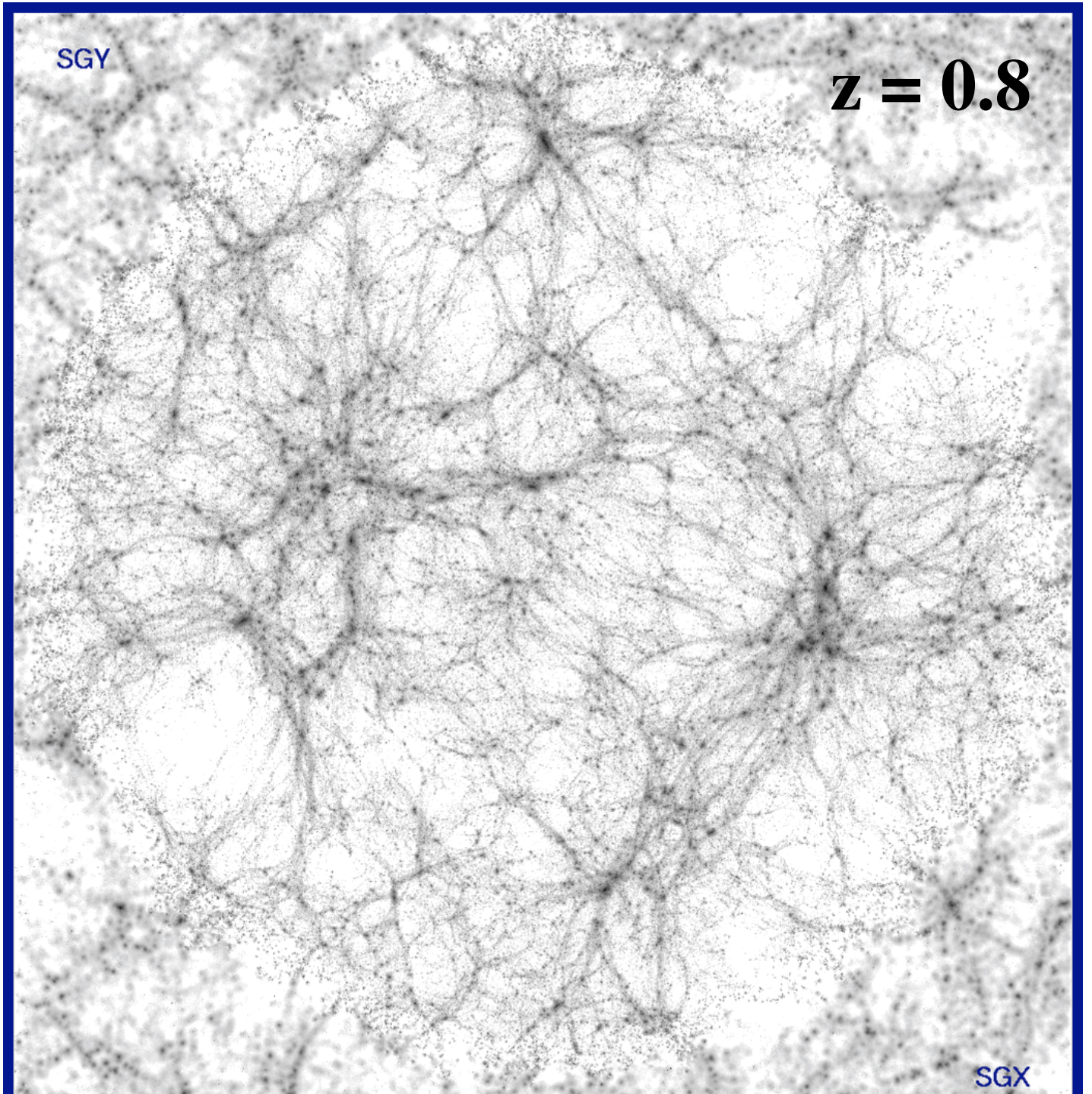




SGY

**$z = 0.8$**

SGX

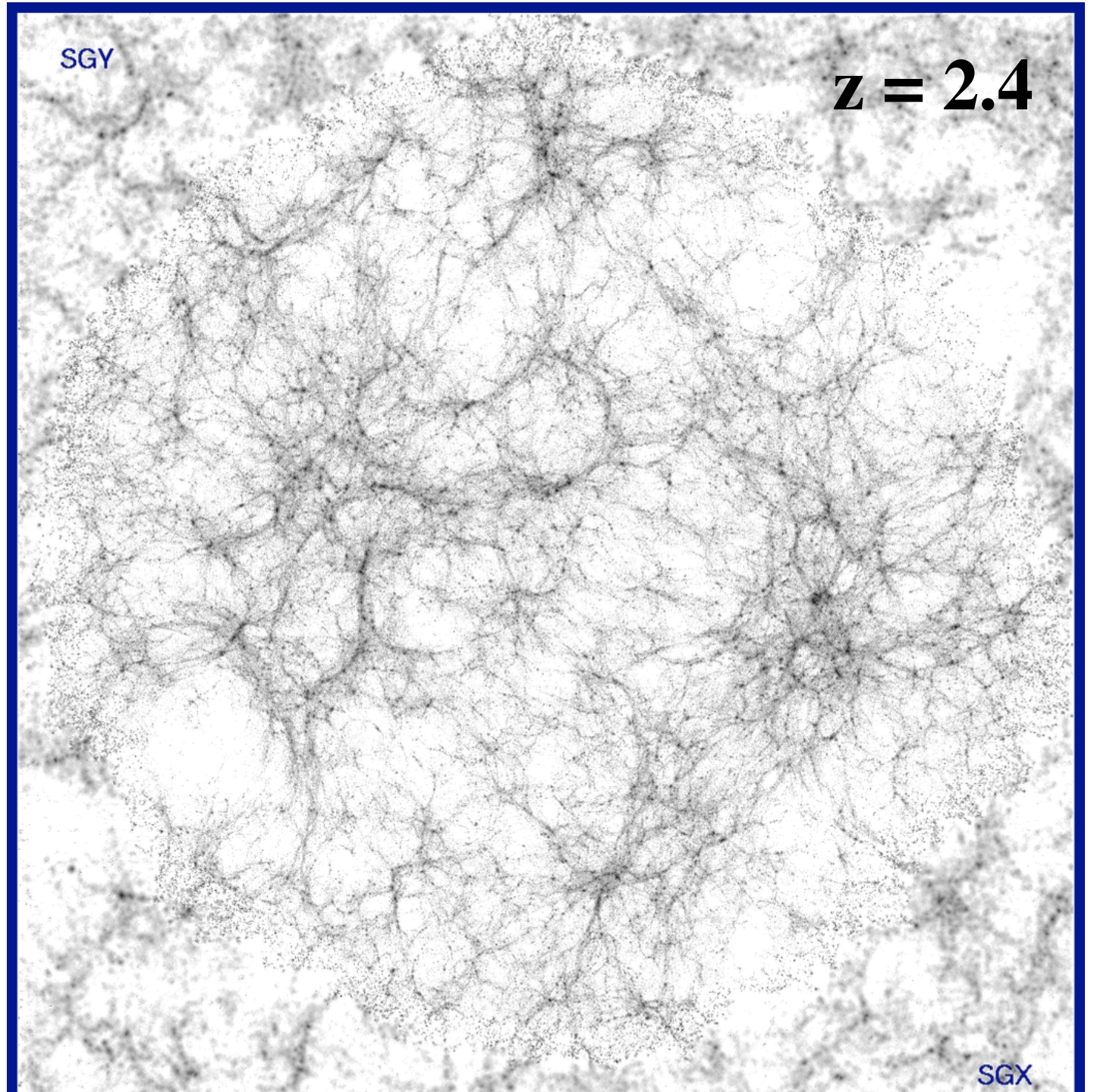




SGY

**$z = 2.4$**

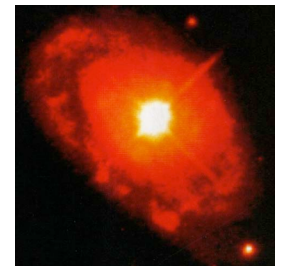
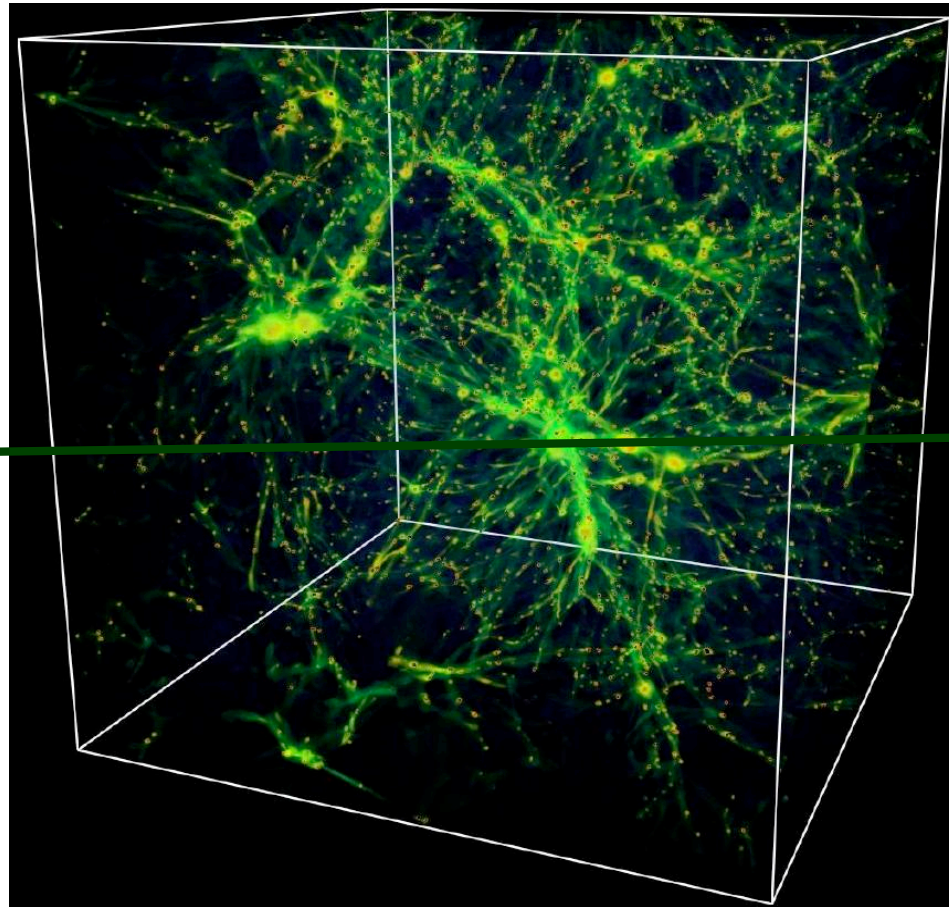
SGX



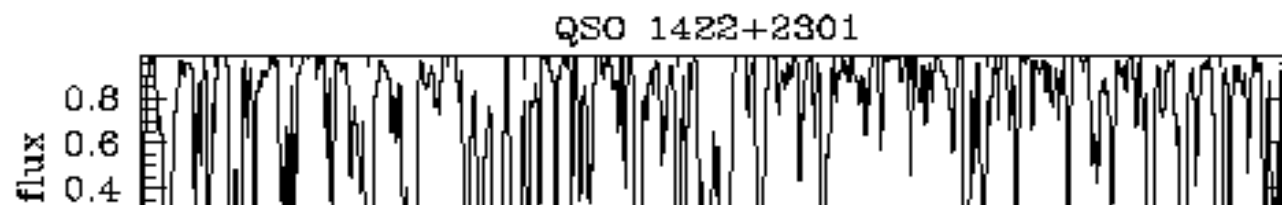
# Lyman Alpha Forest Simulation: Cen et al 2001



You

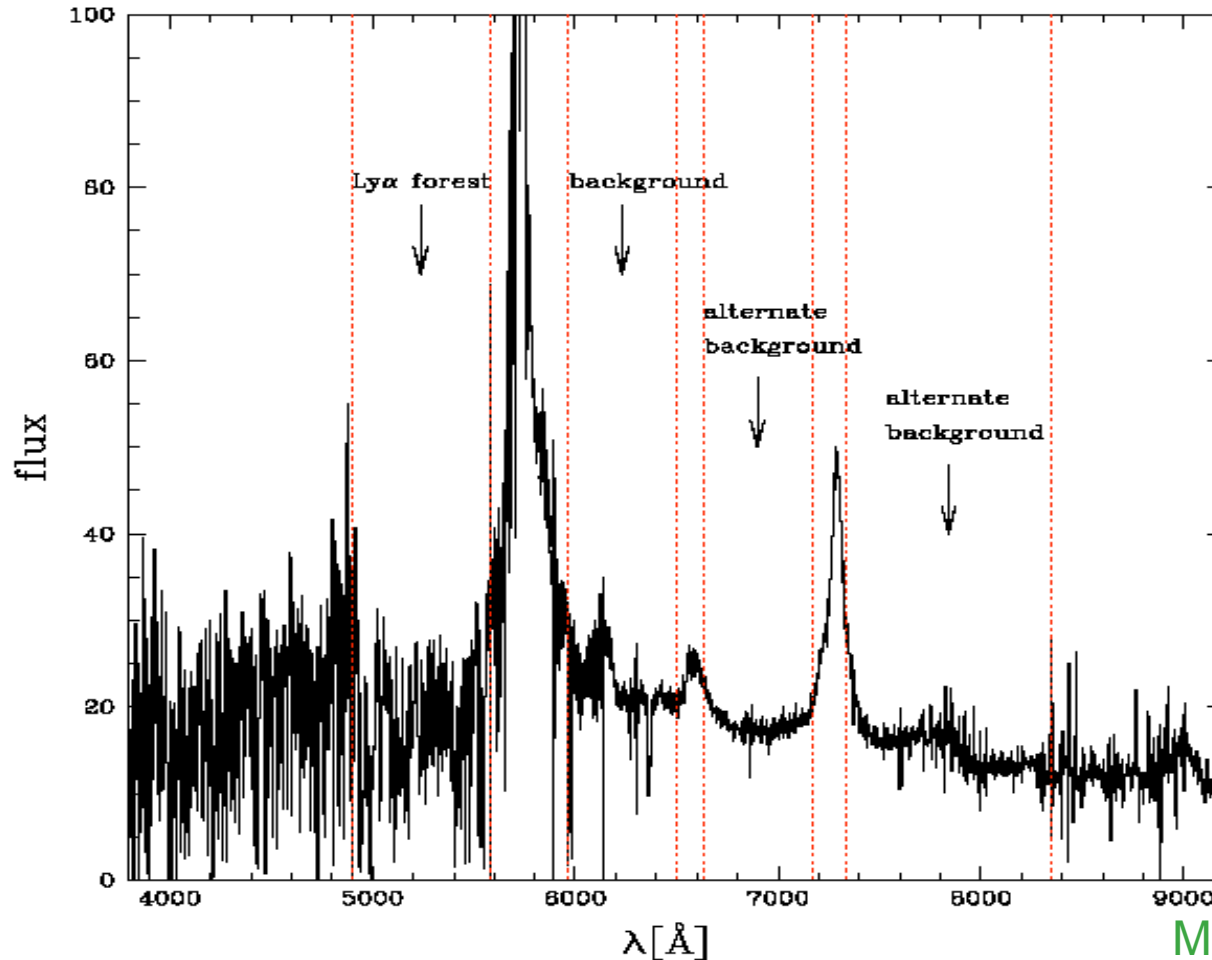


Quasar



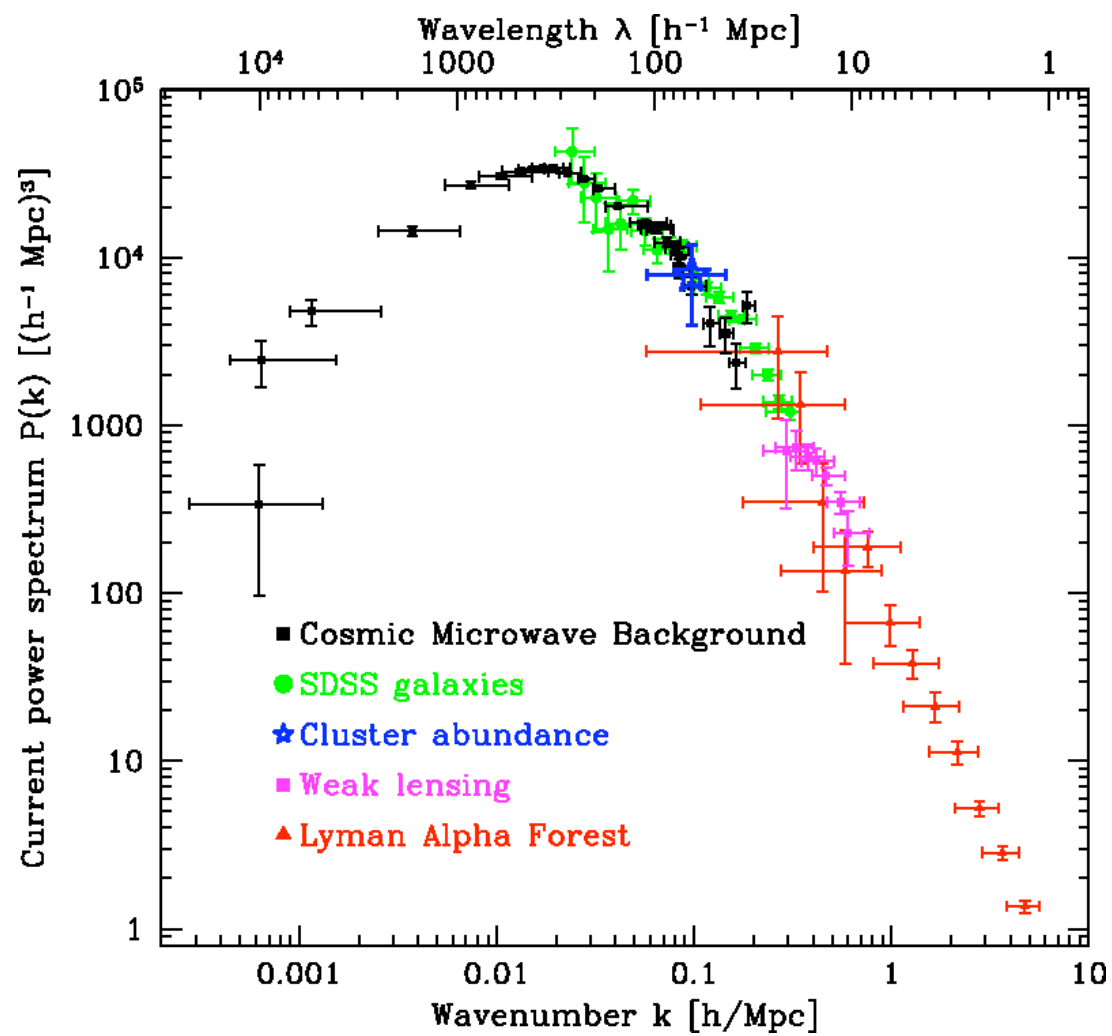


# Lyman alpha forest



McDonald et al. 02

Photons with energy  $>$  ( $n=1$  to  $n=2$  transition energy) get absorbed along the line of sight as they lose energy due to cosmic redshift. Every absorption line corresponds to *cloud* of neutral hydrogen.



# Cosmological (Active) Neutrinos

Neutrinos are in equilibrium with the primeval plasma through weak interaction reactions. They decouple from the plasma at a temperature

$$T_{dec} \approx 1\text{MeV}$$

We then have today a Cosmological Neutrino Background at a temperature:

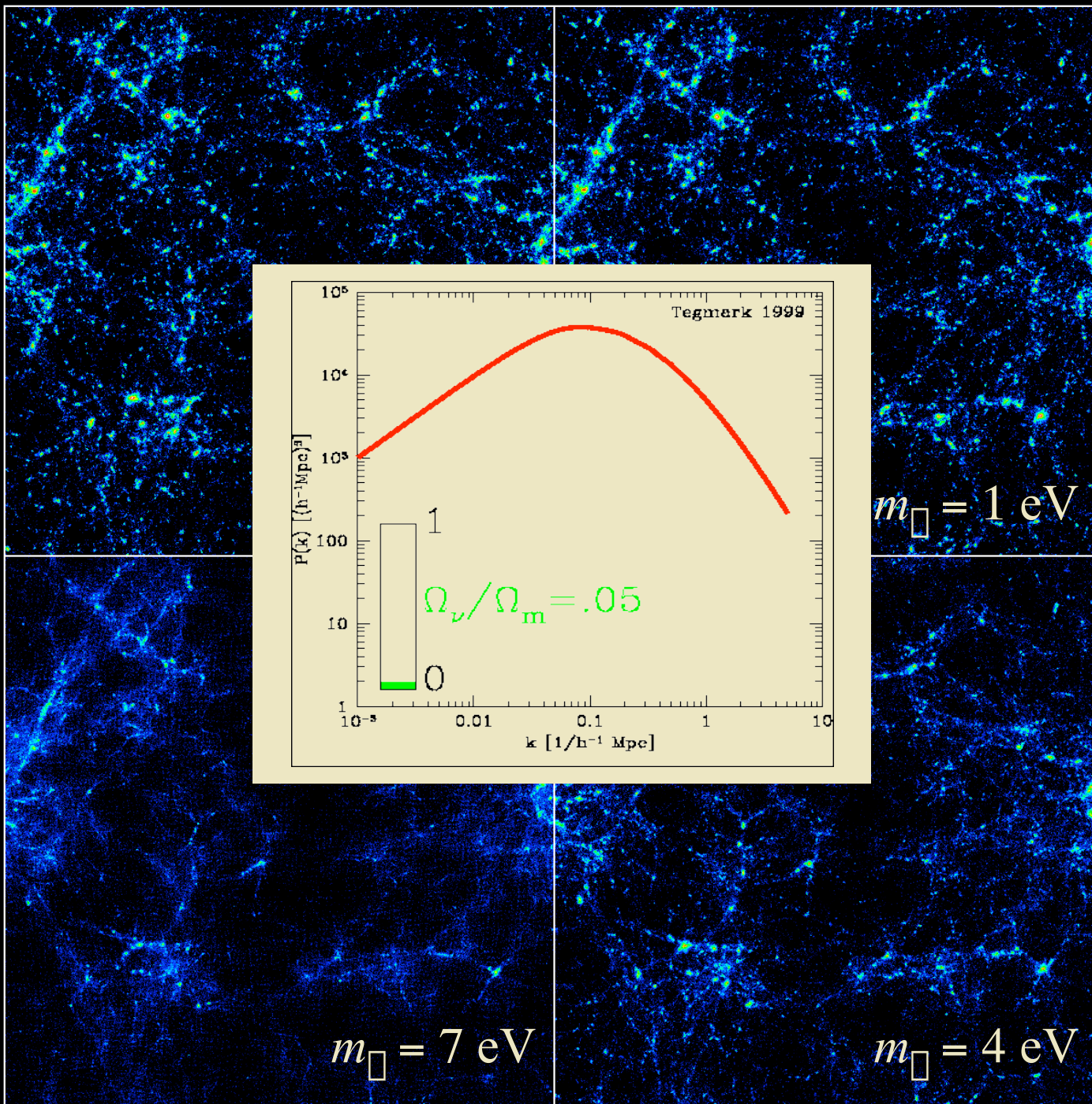
$$T_\nu = \left(\frac{4}{11}\right)^{1/3} T_\gamma \approx 1.945\text{K} \quad kT_\nu \approx 1.68 \cdot 10^{-4} \text{eV}$$

With a density of:

$$n_f = \frac{3}{4} \frac{\pi^2(3)}{\pi^2} g_f T_f^3 \quad n_{\nu_k \bar{\nu}_k} \approx 0.1827 \cdot T_\nu^3 \approx 112 \text{cm}^{-3}$$

That, for a massive neutrino translates in:

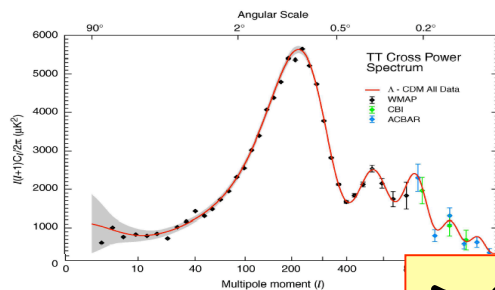
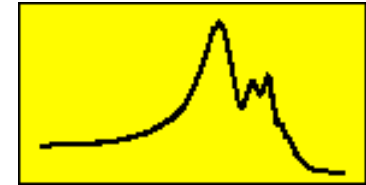
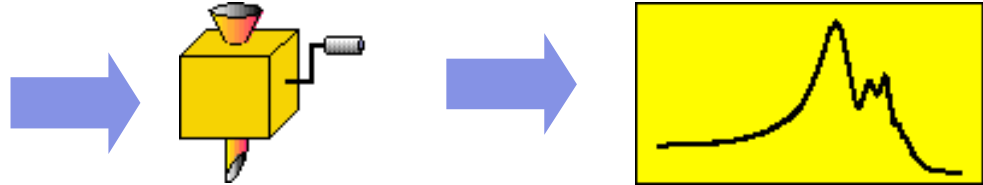
$$\Omega_k = \frac{n_{\nu_k \bar{\nu}_k} m_k}{\rho_c} \approx \frac{1}{h^2} \frac{m_k}{92.5 \text{eV}} \approx \Omega_\nu h^2 = \frac{\sum_k m_k}{92.5 \text{eV}}$$



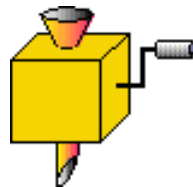
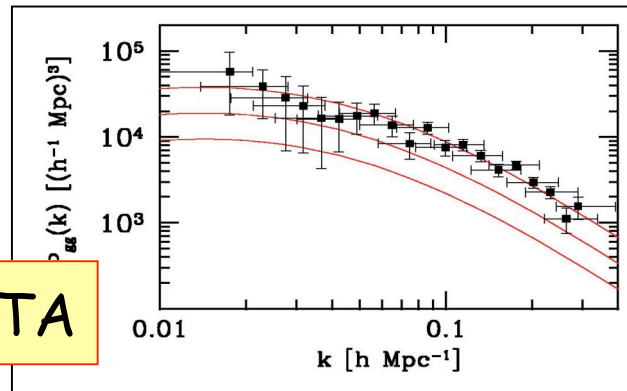


# How to get a bound (measurement) of neutrino masses from Cosmology

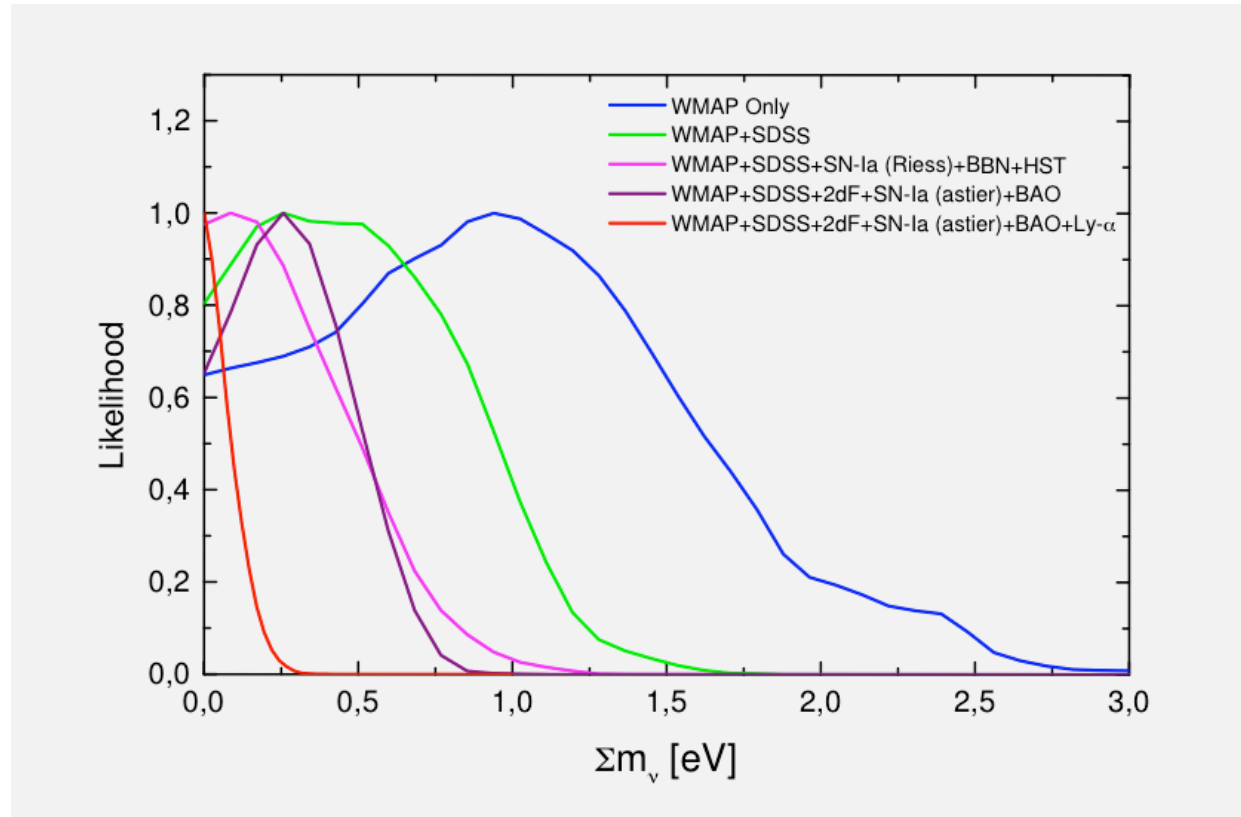
Fiducial cosmological model:  
 $(\Omega_b h^2, \Omega_m h^2, h, n_s, \tau, \Sigma m_\nu)$



DATA



PARAMETER  
ESTIMATES

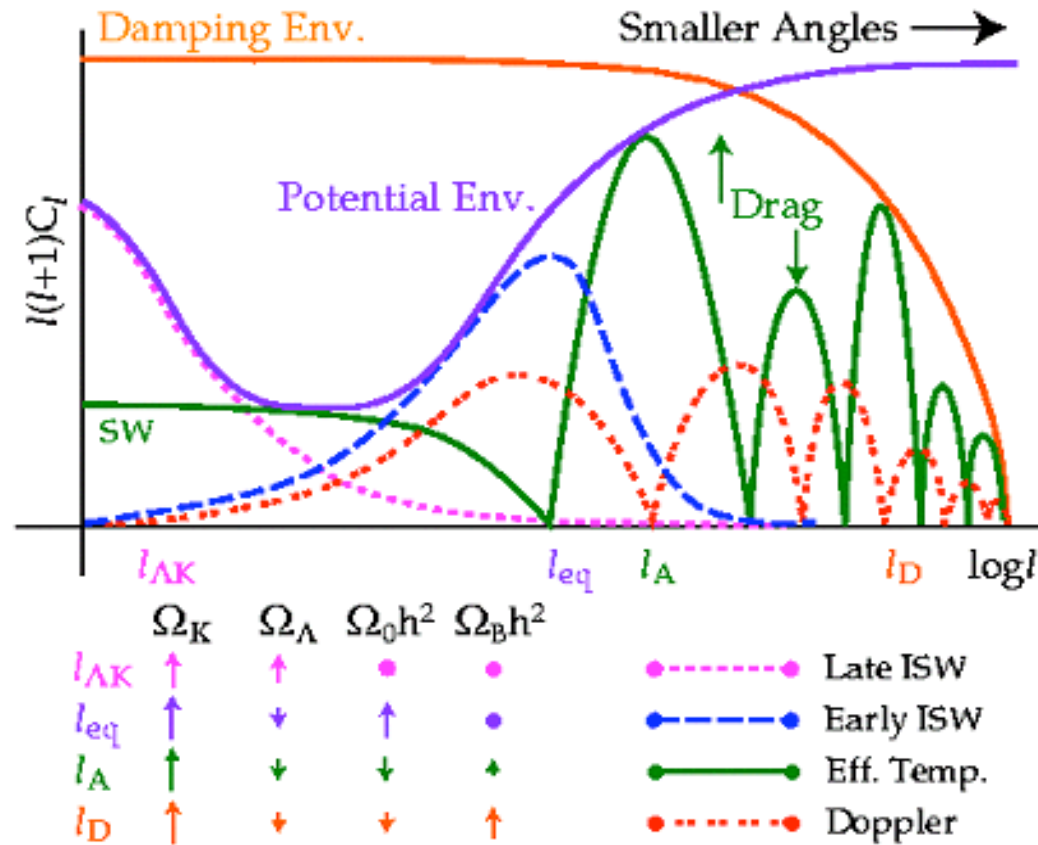


Bounds on  $\Sigma m_\nu$  for increasingly rich data sets (assuming 3 Active Neutrino model):

Case	Cosmological data set	$\Sigma$ bound ( $2\sigma$ )
1	WMAP	$< 2.3$ eV
2	WMAP + SDSS	$< 1.2$ eV
3	WMAP + SDSS + $SN_{\text{Riess}}$ + HST + BBN	$< 0.78$ eV
4	CMB + LSS + $SN_{\text{Astier}}$	$< 0.75$ eV
5	CMB + LSS + $SN_{\text{Astier}}$ + BAO	$< 0.58$ eV
6	CMB + LSS + $SN_{\text{Astier}}$ + Ly- $\alpha$	$< 0.21$ eV
7	CMB + LSS + $SN_{\text{Astier}}$ + BAO + Ly- $\alpha$	$< 0.17$ eV

What about  $N > 3$  ?

# Extra neutrino light component: effects on the CMB



Hu, Sugiyama, Silk, Nature 1997, astro-ph/9604166



# Integrated Sachs-Wolfe effect

while most cmb anisotropies arise on the last scattering surface, some may be induced by passing through a time varying gravitational potential:

$$\frac{\delta T}{T} = -2 \int d\tau \dot{\Phi}(\tau)$$

linear regime - integrated Sachs-Wolfe (ISW)  
non-linear regime - Rees-Sciama effect

when does the linear potential change?

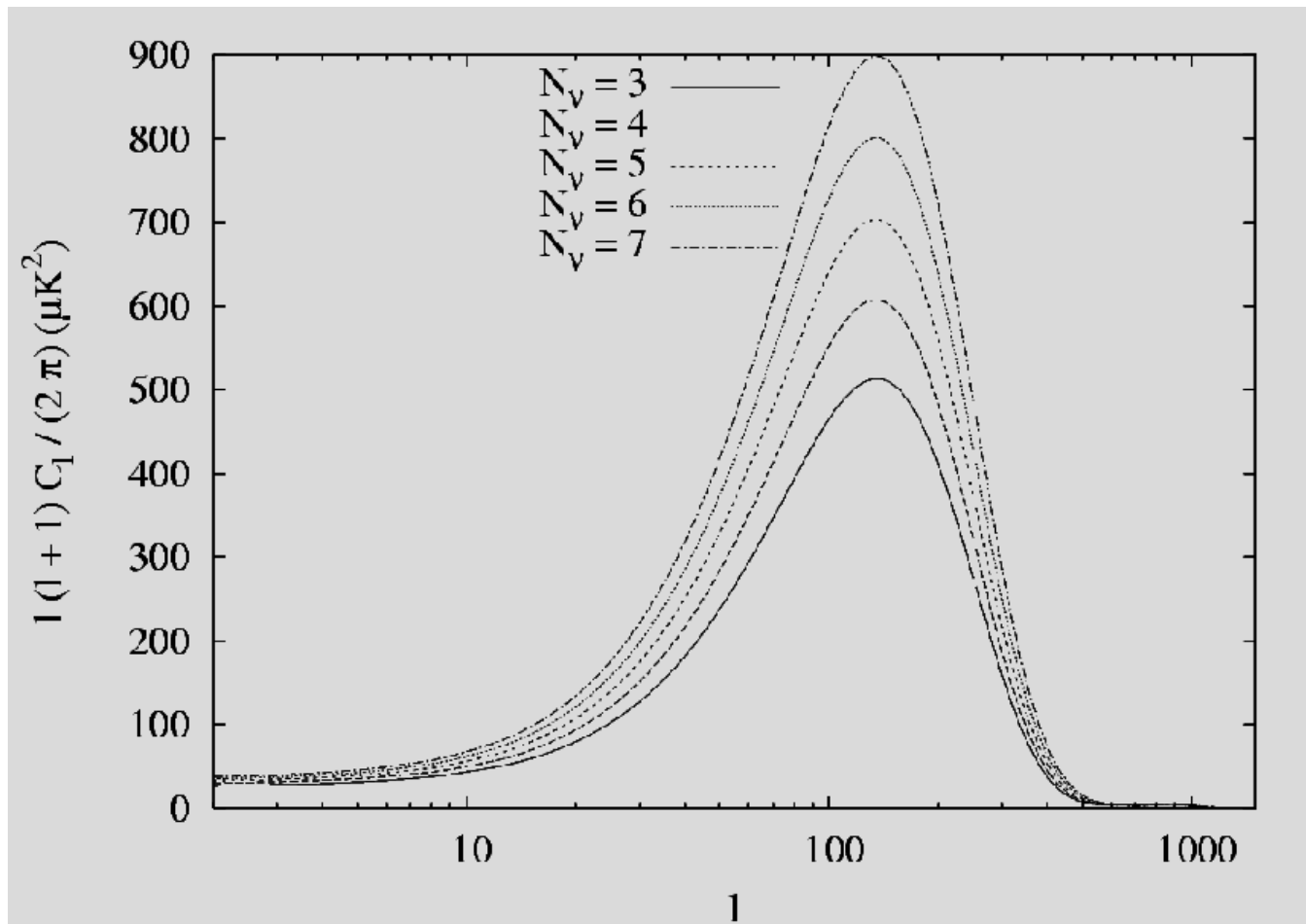
$$\nabla^2 \Phi = 4\pi G a^2 \bar{\rho} \delta$$

Poisson's equation

- changes during **radiation** domination
- decays after curvature or dark energy come to dominate ( $z \sim 1$ )

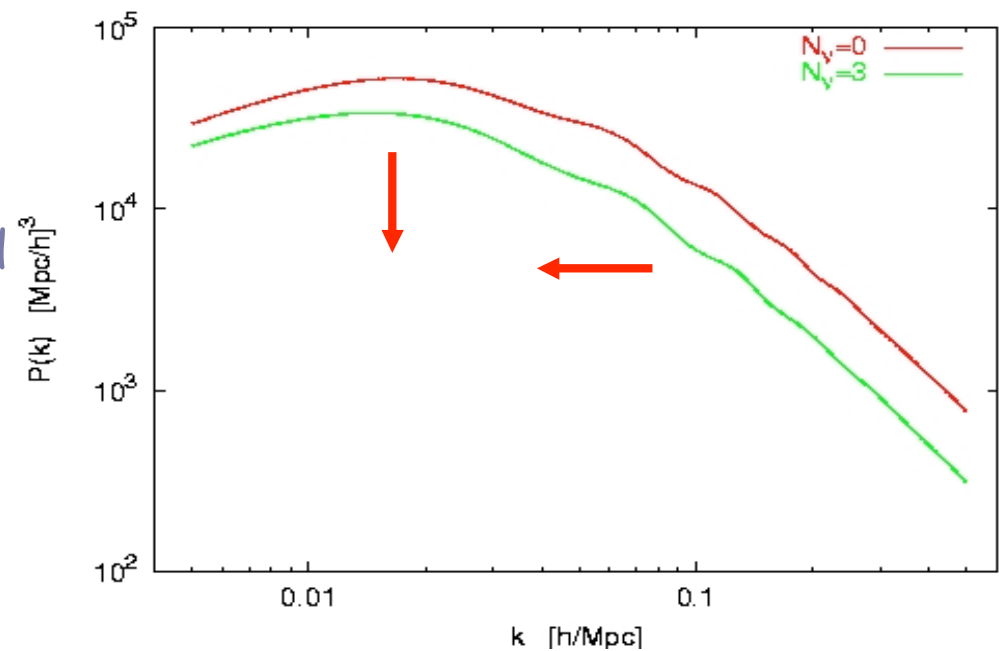
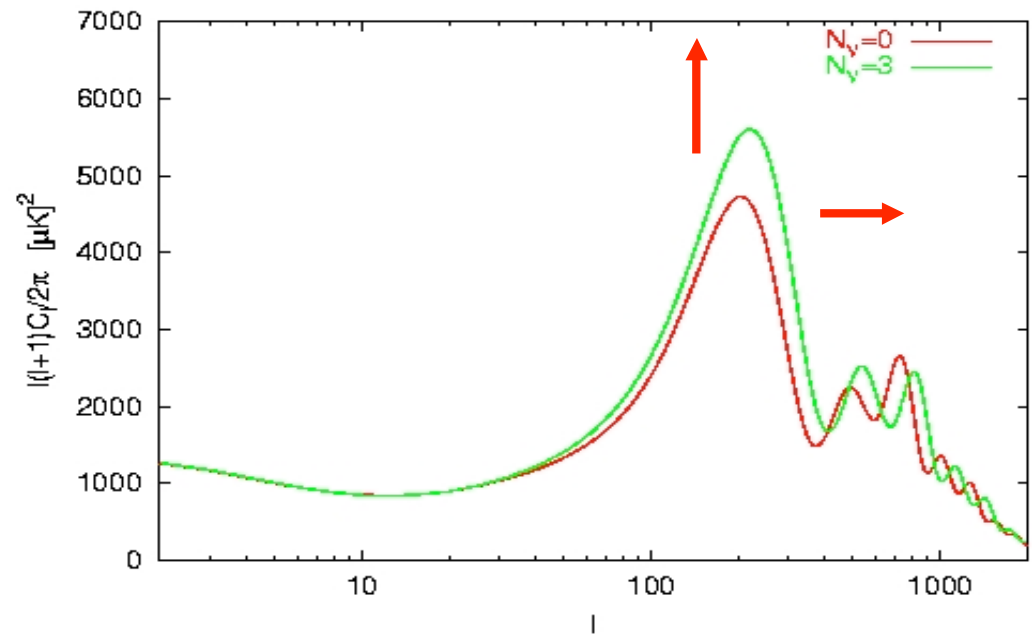
# Effect of Neutrinos in the CMB: ISW

Changing the number of neutrinos (assuming them as massless) shifts the epoch of equivalence, affecting the ISW:

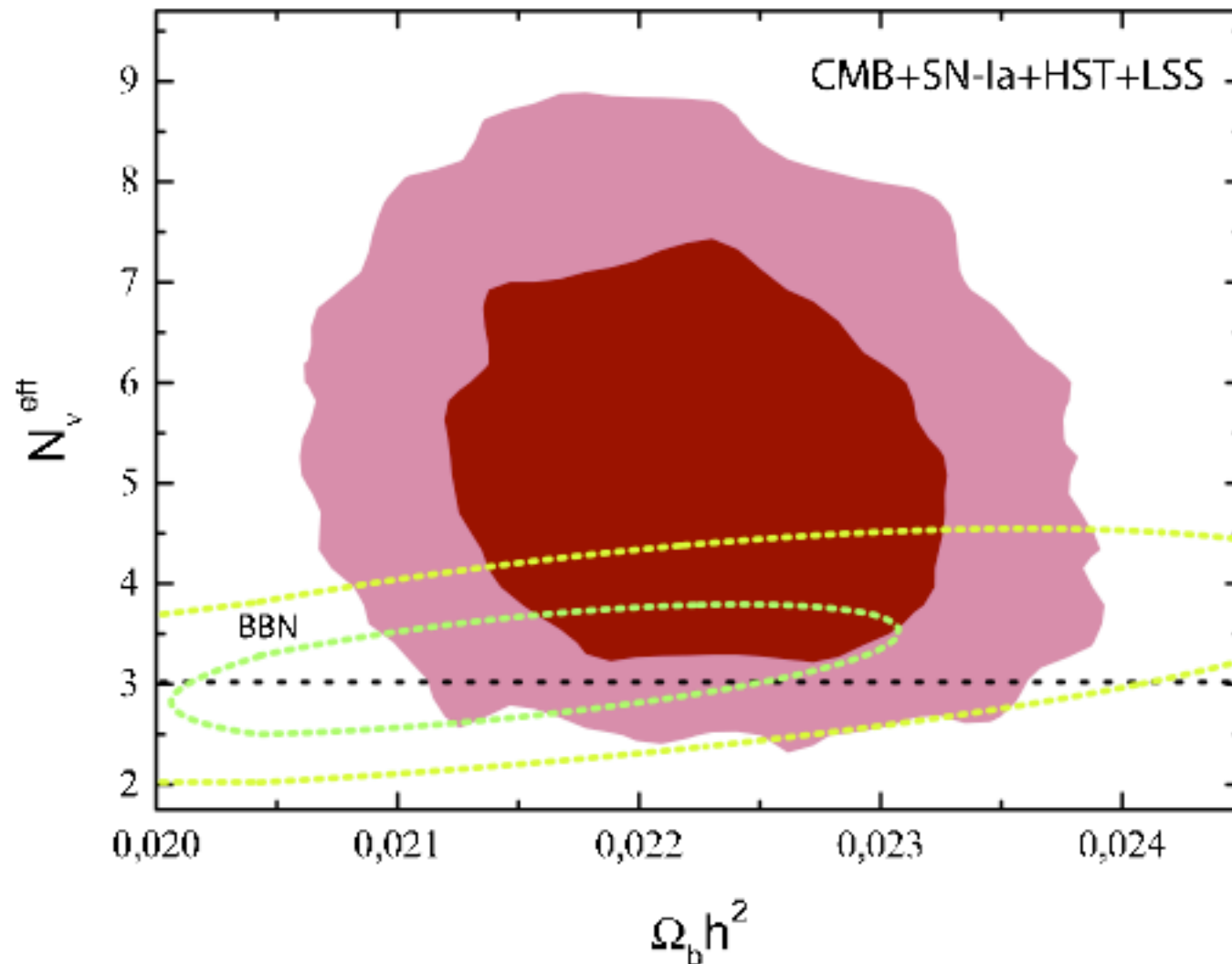


Increasing the Neutrino Massless number postpone the equivalence (while keeping constant the time of decoupling).

This produces a shift in the CMB power spectra since changes the sound horizon at decoupling. The height of the first peak is also increased thanks to the Early Integrated Sachs-Wolfe. The LSS matter power spectrum is also shifted since the size of the horizon at equivalence is now larger. There is less growth of perturbations in the MD regime.

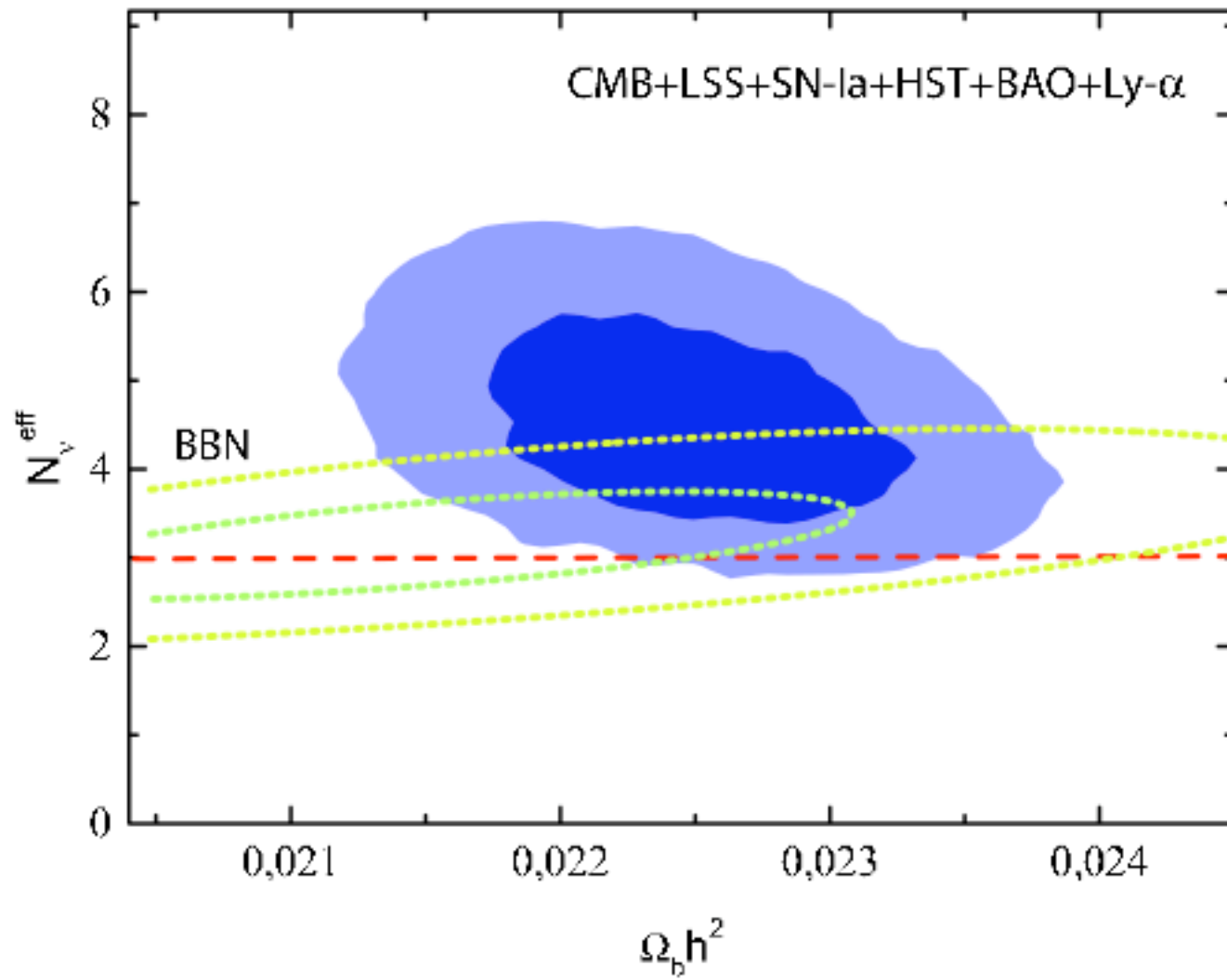


## Latest Analysis: Indication for $N > 3$ from Cosmology ?



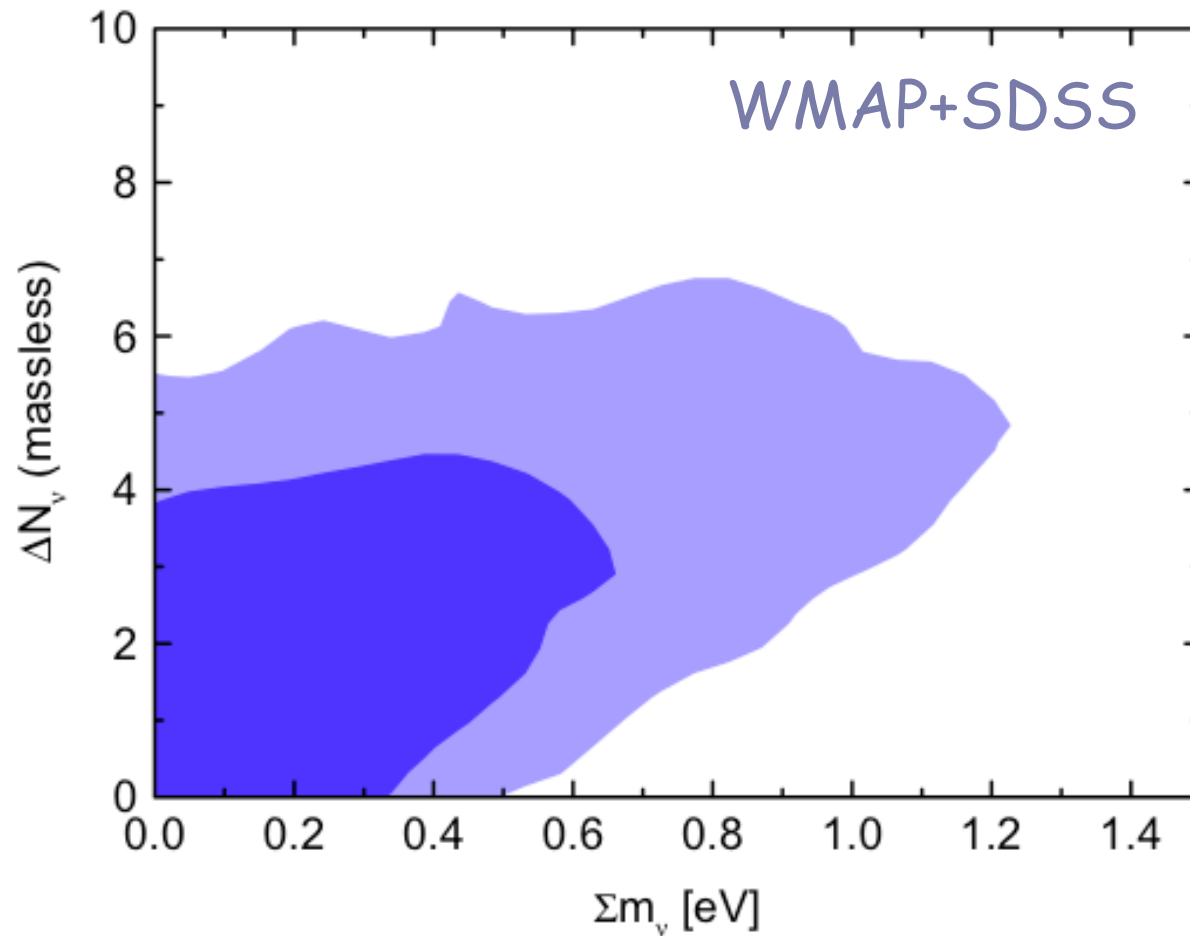
Mangano, Melchiorri, Mena, Miele, Slosar JCAP03(2007)006





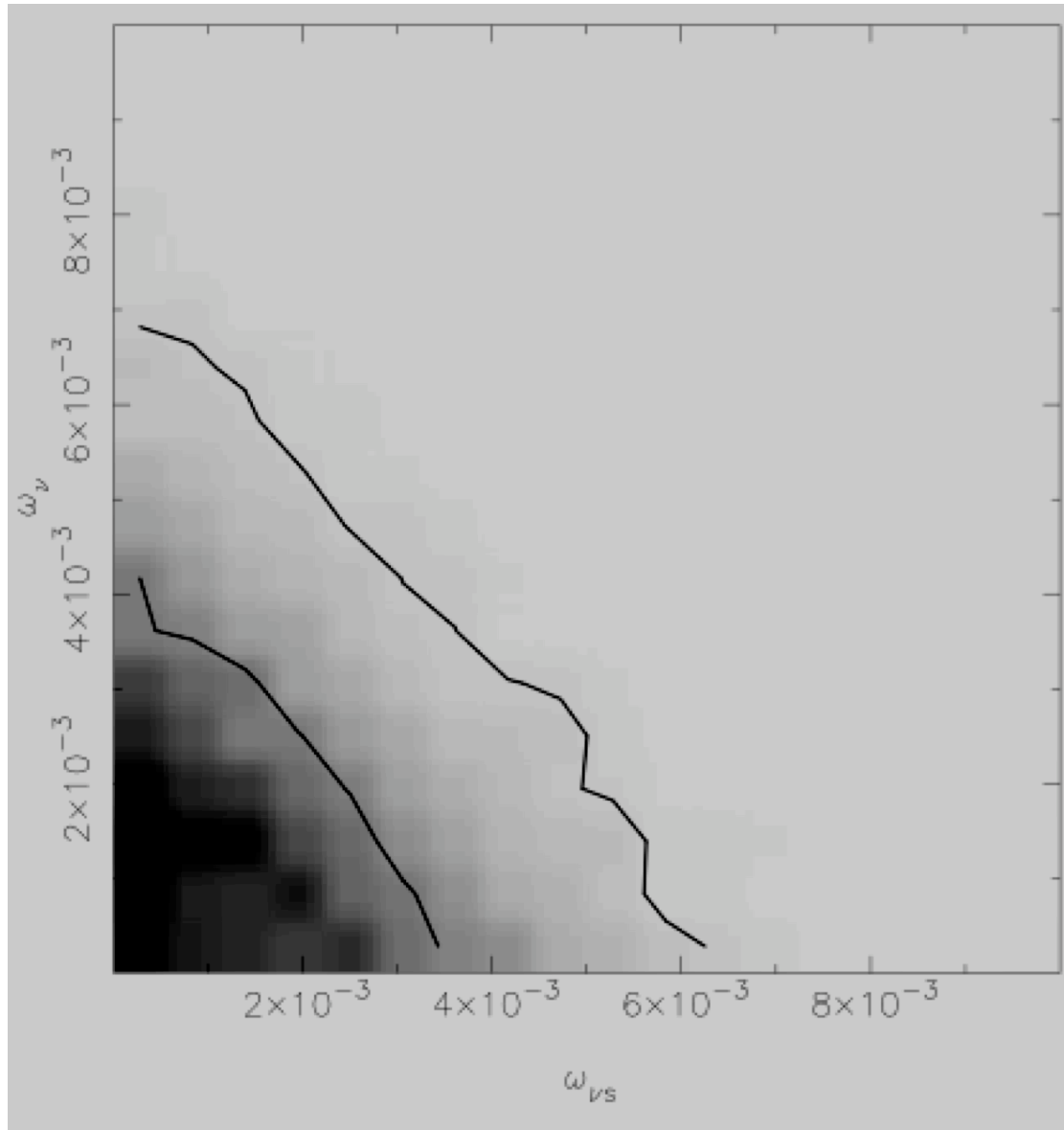
Mangano, Melchiorri, Mena, Miele, Slosar JCAP03(2007)006

## Massless Neutrino Number vs Active Neutrino Masses



Adding an extra relativistic component change the bound by 10-20% per specie (See e.g. Melchiorri, Serra PRD 2006)

## What about a fourth massive sterile neutrino ?



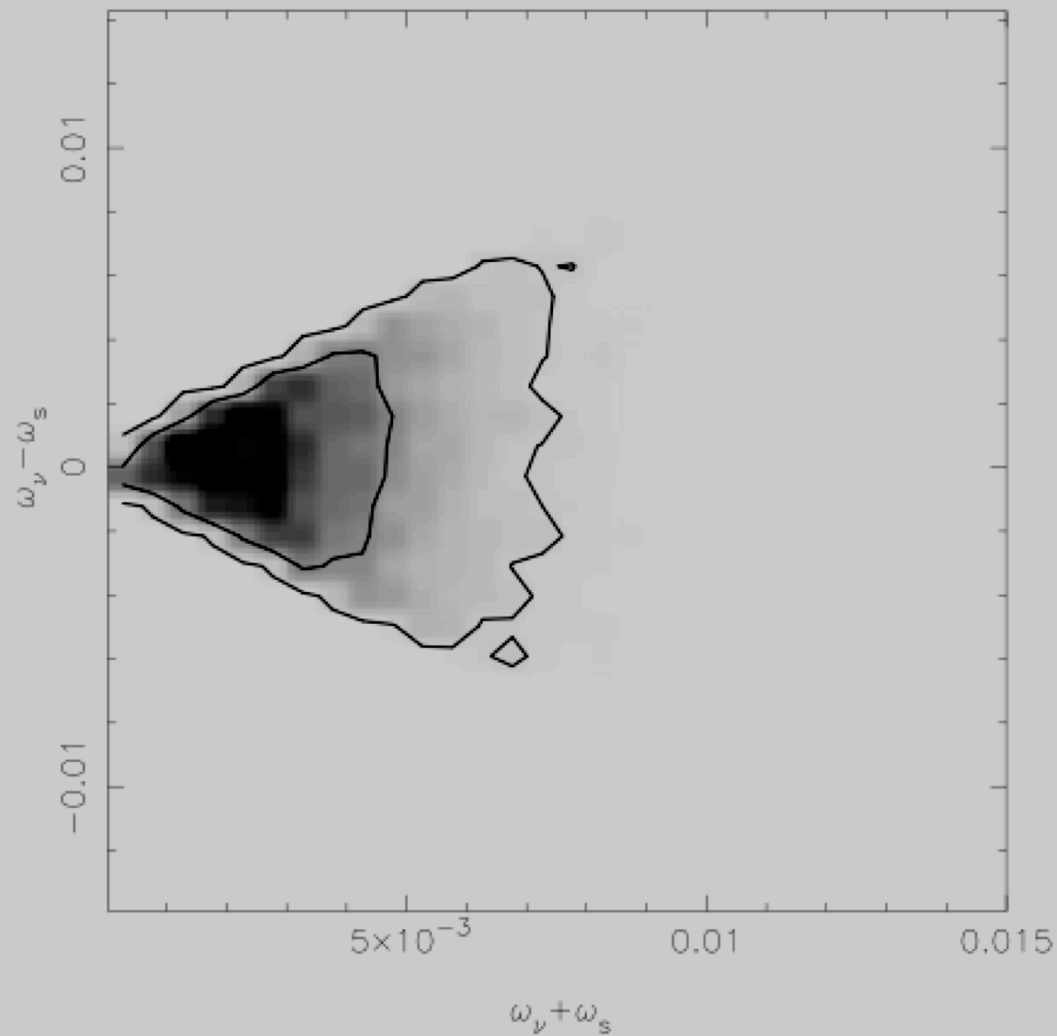
CMB+2df+  
Sloan+Ly- $\alpha$

$$\Omega_s = 0.0106 \frac{m_s}{\text{eV}}$$

$$\Omega_{\bar{\nu}} = 0.0106 \frac{3m_{\bar{\nu}}}{\text{eV}}$$

$m_s < 0.23 \text{ eV}$  at  
95% c.l.

Dodelson,  
Melchiorri,  
Slosar,  
Phys.Rev.Lett.  
97 (2006) 04301



Cosmology tests  
only the sum  
of the neutrino  
masses  
(see also Slosar 2006)

Howver sterile neutrino can be non-thermal.

Thermalization occurs if:

$$\left[ m^2 \sin^4 \theta \right] > 3 \times 10^{16} eV^2$$

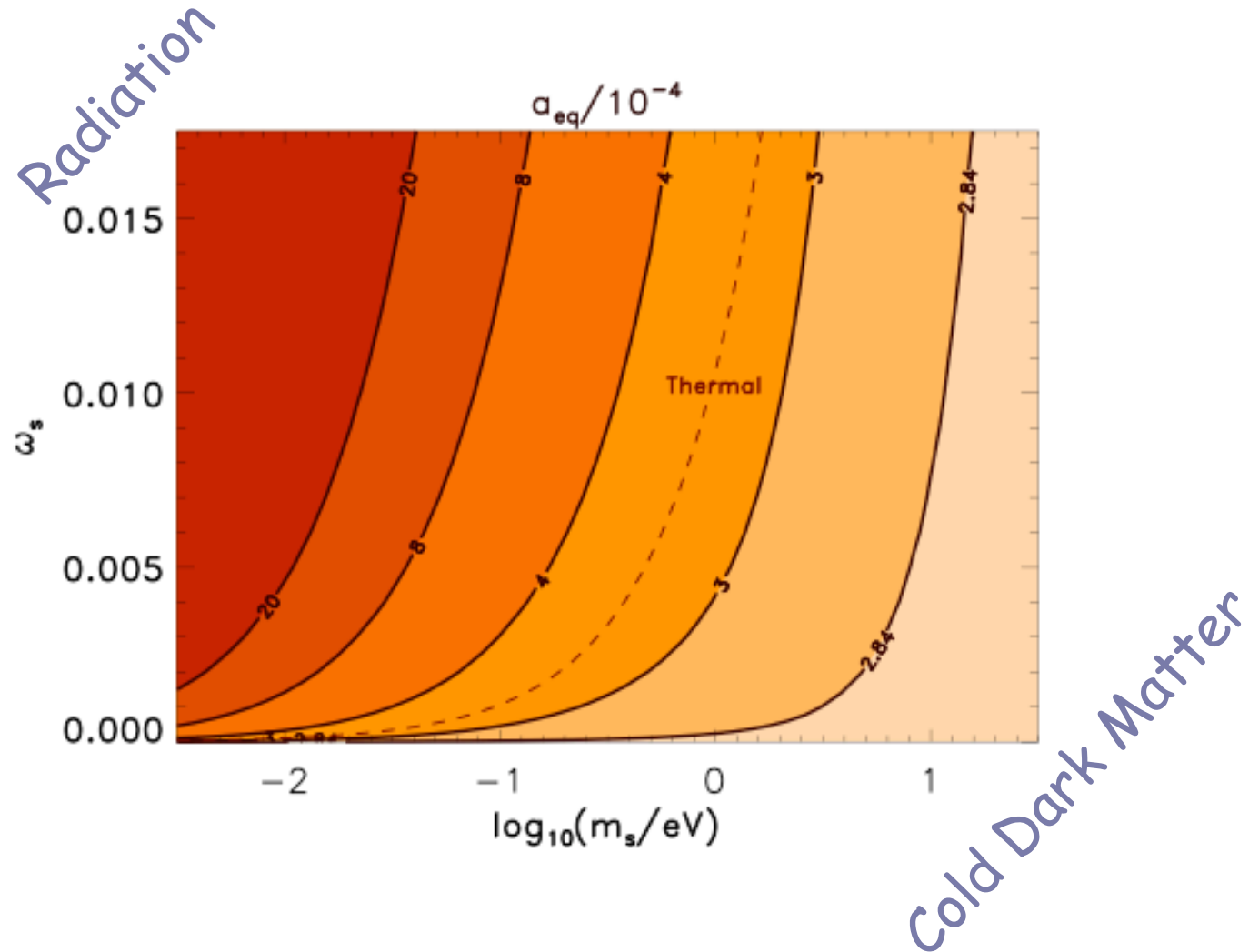
In the simplest models with one sterile neutrino this Condition is satisfied bu there are many ways of evading thermalization (see e.g. Abazajian, 2003).

In practice:

$$\theta_s \neq 0.0106 \frac{m_s}{eV}$$

Mass and cosmological energy density should be considered as independent parameters !

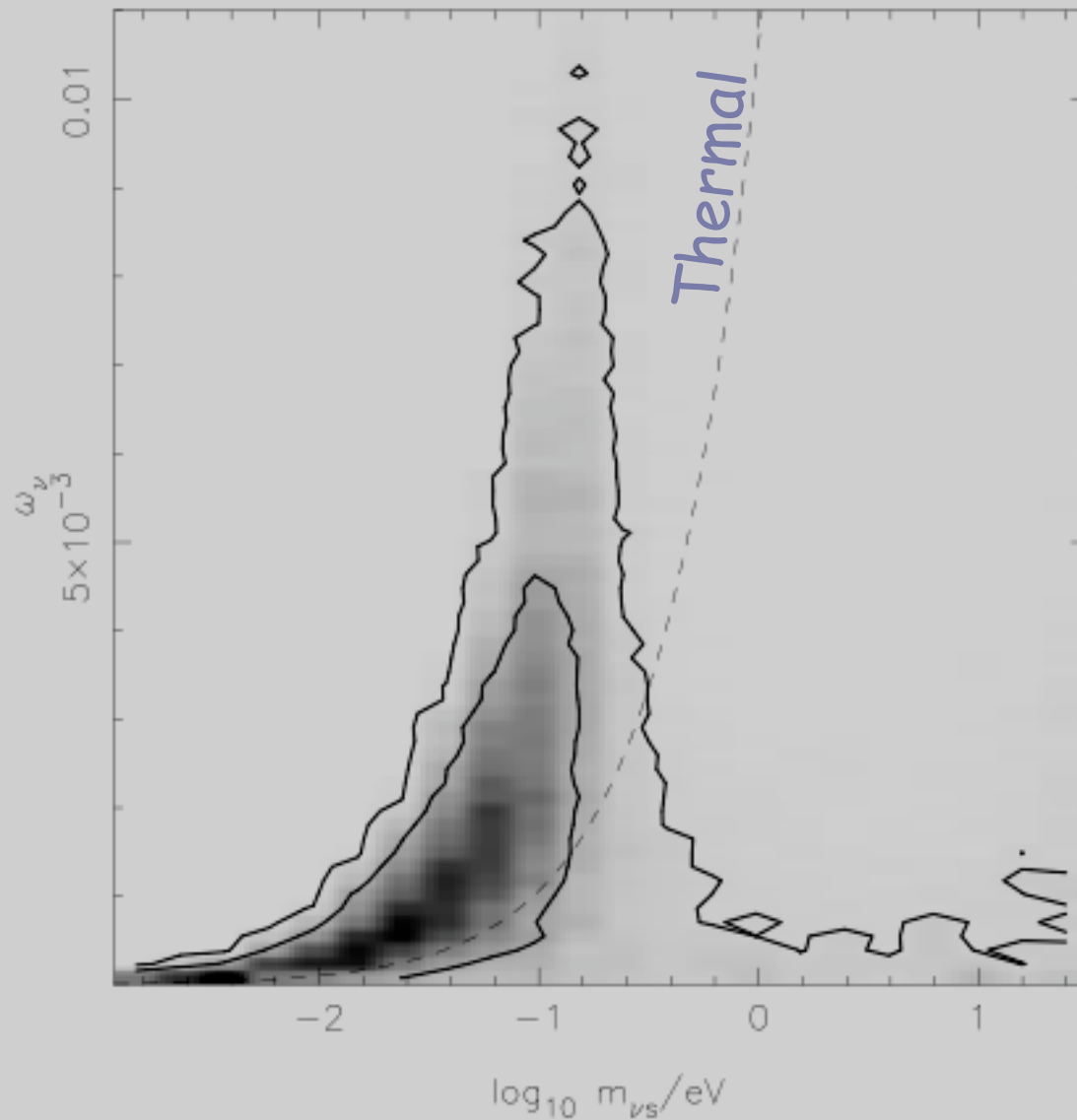
Effects on the scale of equality:



Dodelson, Melchiorri, Slosar, PRL 2006



## Constraints on non-thermalized sterile neutrino

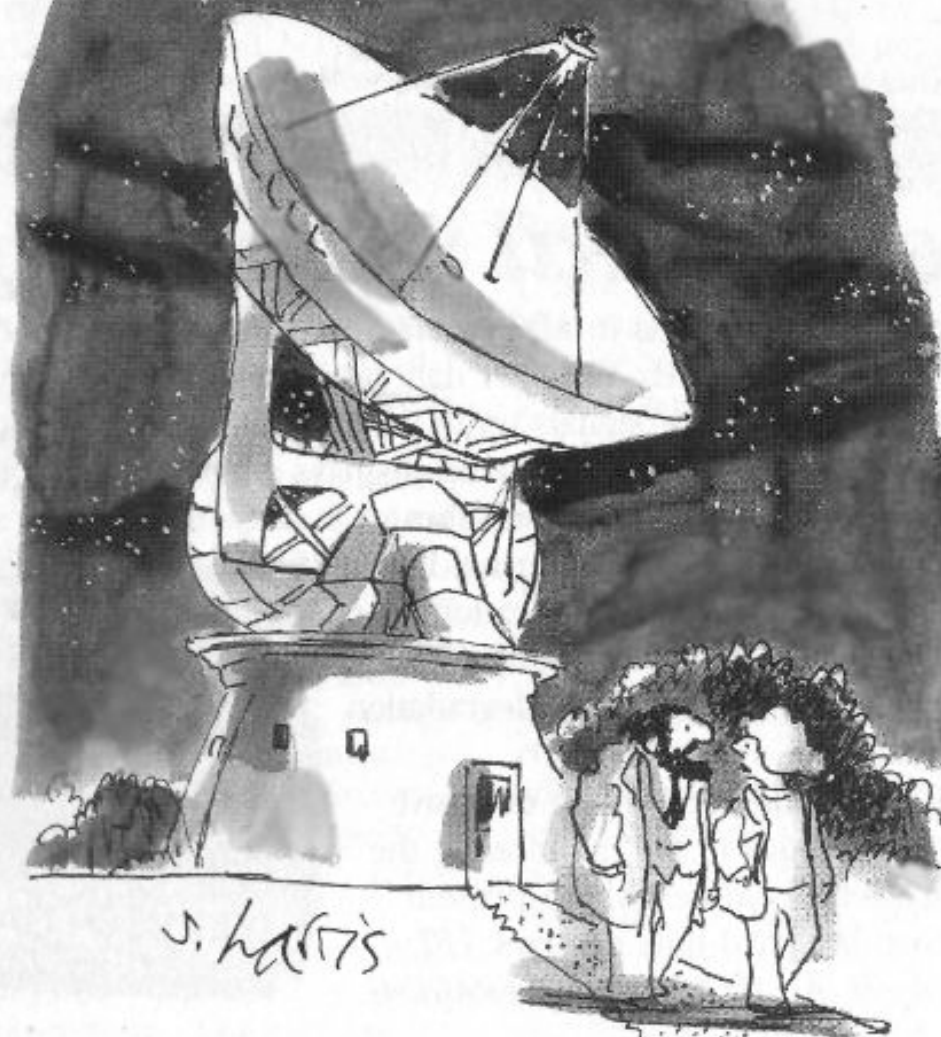


Energy density  
Can be higher  
For smaller masses

You may have large  
masses but in this  
case they are not  
cosmologically  
relevant.

# Conclusions

- ◆ Current CMB and LSS data are in very good agreement with the standard scenario. Limits on  $N_{\text{eff}}$  are still weak, Sensitivity comparable to BBN is possible in the very near future. If Lyman-alpha are included there is some indication that  $N > 3$ .
- ◆ Cosmological constraints on neutrino mass are rapidly improving. If one includes Ly-alpha then  $\Sigma m_\nu < 0.17$  eV. Tension with the  $0 < \Delta \rho < 0.1$  results. Fourth sterile neutrino mass (if thermal constrained to be  $m_s < 0.25$  eV). LSND,  $0 < \Delta \rho < 0.1$  and cosmology all incompatible. Neutrino mass detection up to  $\Sigma m_\nu = 0.05$  eV is possible in the very near future.
- ◆ The constraints are model dependent (quite common in physics...)



"I'LL BE WORKING ON THE LARGEST AND SMALLEST  
OBJECTS IN THE UNIVERSE—SUPERCLUSTERS AND  
NEUTRINOS. I'D LIKE YOU TO HANDLE EVERYTHING IN BETWEEN."