# The impact of solar and atmospheric parameter uncertainties on the measurement of $\theta_{13}$ and $\delta$

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parameter correlations and degeneracies  $\begin{cases} in the \nu_e \to \nu_\mu(\nu_\tau) \ appearance \\ in the \nu_e \to \nu_e, \ \nu_\mu \to \nu_\mu \ disappearance \end{cases}$ 

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## Introduction

#### **Present State of Neutrino Oscillation Parameters**

- solar angle and mass difference are "known"
- atmospheric sector:  $|\Delta m^2_{23}|$  and departure of  $\sin^2 2 heta_{23}$  from maximal mixing are "known"



#### Degeneracies

#### ... in appearance: widely discussed in literature

Fogli96, Burguet-Castell01, Minakata01...

$$\underbrace{N_{\alpha\beta}^{\pm}(\bar{\theta}_{13},\bar{\delta};\bar{s}_{atm},\bar{s}_{oct})}_{true \ parameters} = \underbrace{N_{\alpha\beta}^{\pm}(\theta_{13},\delta;s_{atm}=\pm\bar{s}_{atm};s_{oct}=\pm\bar{s}_{oct})}_{guessed \ parameters}$$

 $N_{\beta}^{\pm}$  = number of charged leptons  $l_{\beta}^{\pm}$  observed in a given detector  $\alpha, \beta =$  neutrino flavours

- four different models (each of them with a definite  $s_{atm}, s_{oct}$  choice) must be used to fit the data on the lhs
- all in all, they form an eightfold degeneracy: the *intrinsic*, the *sign*, the *octant* and the mixed clones

1 large and dangerous nuisance that increases the resulting uncertainty on  $heta_{13}$  and  $\delta$ 

#### ... in disappearance

a similar appoach can be applied to study the  $\begin{cases} & \text{in the } \nu_e \to \nu_e \text{ at the } \beta - \text{Beam} \\ & \text{in the } \nu_\mu \to \nu_\mu \text{ at the SPL Super} - \text{Beam} \end{cases}$ 

The facilities

#### The $\beta$ -Beam

- $\nu_e \ (\bar{\nu}_e)$  from <sup>18</sup>Ne (<sup>6</sup>He) ions accelerated at  $\gamma = 100 \ (\gamma = 60)$
- 10 yrs of  $\nu_e$  and  $\bar{\nu}_e$   $\beta$ -Beam run with a 440 KTon detector
- background and efficiencies

J.	Bouchez et al.,	AIP	Conf.	Proc.	721,	37(2004)	
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appearance :A. Donini et al., hep - ph/0406132

J. Burguet – Castell et al., Nucl. Phys. B695, 217(2004)

disappearance : A. Donini et al., hep - ph/0411402

#### The Super-Beam

- $\nu_{\mu}$   $(\bar{\nu}_{\mu})$  from  $\pi^+$   $(\pi^-)$  (60 m decay tunnel assumed)
- 2 yrs of  $u_{\mu}$  and 8 yrs of  $\bar{
  u}_{\mu}$  Super-Beam run with a 440 KTon detector
- background and efficiencies

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appearance :
J. J. Gomez – Cadenas et al., hep – ph/0105297
A. Donini et al., hep – ph/0406132
disappearance :
A. Donini et al., hep – ph/0411402
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#### plan of the talk

• the  $\beta$ -Beam setup

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-correlation and degeneracies in \nu_{\rm e} disappearance (appendix)
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impact of the parameter uncertainties: 
-full 3-parameter analysis in the measurement of \theta_{13} and \delta
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• the Super-Beam setup

 $\begin{array}{l} \mbox{correlation and degeneracies in $\nu_{\mu}$ disappearance} \\ \mbox{the atmospheric sector}: \delta(\Delta m^2_{\rm atm}), \delta(\theta_{23}) \\ \mbox{combining appearance and disappearance channels to measure $\theta_{13}$ and $\delta$} \end{array}$ 

# **1** Full 3-par analysis of the appearance channel

- we investigate the impact of uncertainties in the atmospheric and solar parameters in the measure of  $(\theta_{13},\delta)$ 

- in order to understand how any single parameter affect the measurement, we perform three-parameters fits in  $\theta_{13}$ ,  $\delta$  and only one of the following parameters in turn  $x = \theta_{12}$ ,  $\Delta m^2_{12}$ ,  $\theta_{23}$  and  $\Delta m^2_{23}$ 

- For each degeneracy
  - we build a three-parameters  $\chi^2$

$$\chi^{2}(\theta_{13},\delta,x) = \sum_{p=\pm} \left[ \frac{N^{p}(\theta_{13},\delta,x,s_{atm},s_{oct}) - N^{p}(\bar{\theta}_{13},\bar{\delta},\bar{x},\bar{s}_{atm},\bar{s}_{oct})}{\delta_{N^{p}}} \right]^{2}$$

$$(\delta N^{p})^{2} = \sigma_{N}^{2} p + (s N^{p})^{2} + (s B^{p})^{2} \qquad s = sys, B^{p} = background$$

parameter	present Bahcall04 (solar)-Gonzalez-Garcia03 (atm)	expected LOI of T2K (hep-ex/0106019)	C.V.
$ an^2  heta_{12}$	<b>0.30–0.54</b> (29° – 36°)	_	0.39
$\Delta m^2_{12}/10^{-5}~eV^2$	7.5–9.1	—	8.2
$ an^2  heta_{23}$	0.53-2.04 (36° - 55°)	<b>0.57–0.88</b> (37° – 43°)	0.7*
$\Delta m^2_{23}/10^{-3}~eV^2$	1.7–3.5	2.3–2.7	2.5
$ an^2  heta_{23}$	new analysis Enrique's talk	<b>0.62–0.85</b> (38° – 43°)	
$\Delta m^2_{23}/10^{-3}~eV^2$	new analysis	2.42-2.61 (2.47-2.64)	

• projection of the 90% CL surface of  $\chi^2(\theta_{13}, \delta, r)$  on the  $(\theta_{13}, \delta)$  plane

- easy to understand
- easy to compare with the usual bidimensional fit

## Full 3-par analysis of the app channel at $\beta$ -Beam

The solar sector

we perform two distinct three-parameters fit in  $\theta_{13}$ ,  $\delta$  and  $x = \theta_{12}$  (for fixed  $\Delta m_{sol}^2$ ) or  $x = \Delta m_{sol}^2$  (for fixed  $\theta_{12}$ )  $\Longrightarrow$  same results  $\Longrightarrow$  only the case  $x = \theta_{12}$  shown



the solar sector does not introduce further uncertainties on the measure of  $\theta_{13}$  and  $\delta$ 

 $\chi^2(\theta_{13}, \delta, x) = \chi^2_{app}(\theta_{13}, \delta, x) + \chi^2_{dis}(\theta_{13}, \delta, x)$ 

The atmospheric sector (I)–current errors

 $x = \theta_{23}$ 

 $(\bar{\theta}_{13}, \bar{\delta}) = (2/7^{\circ}, 45^{\circ})$ 

 $x = \Delta m_{23}^2$ 





- two(dotted) and three-par fit superimposed

-  $\delta(\theta_{13}) \sim 3^{\circ} - 4^{\circ}$ (to compesate a change in  $\theta_{23}$  in the leading term in  $P_{e\mu} : \sin^2(2\theta_{13}) \sin^2 \theta_{23}$ )

- overall error on 
$$\delta$$
 larger for  $x=\Delta m^2_{23}$ 

the atmospheric sector strongly affects the measure of  $\theta_{13}$  and  $\delta$ 

 $\chi^2(\theta_{13}, \delta, x) = \chi^2_{app}(\theta_{13}, \delta, x) + \chi^2_{dis}(\theta_{13}, \delta, x)$ 

The atmospheric sector (II)–expected errors from LOI of T2K

"The overall sensitivity of the first phase is expected to be 1% in precision for  $\sin^2 2\theta_{23}$  and better than  $1 \times 10^{-4} \text{ eV}^2$  for  $\Delta m_{23}^2$ "; we estimate the error on  $\bar{\theta}_{23} = 40^\circ$  and  $\bar{\Delta} m_{23}^2 = 2.5 \cdot 10^{-3} eV^2$  $(\bar{\theta}_{13}, \bar{\delta}) = (2/7^\circ, 45^\circ)$ 



significant reduction on the  $\theta_{13}$ -spread is achieved

for larger values of  $\theta_{13}$  the error is enhanced due to the octant ambiguity

also the  $\delta$ -spread is reduced considerably with respect to the results obtained with present uncertainties

 $\chi^2(\theta_{13}, \delta, x) = \chi^2_{app}(\theta_{13}, \delta, x) + \chi^2_{dis}(\theta_{13}, \delta, x)$ 

The atmospheric sector (III)-new analysis of expected performance of T2K

we consider  $\theta_{23} \in [38.2, 42.7]$  and  $\Delta m_{23}^2 \in [2.42, 2.61] \cdot 10^{-3} \ eV^2$  computed for  $\bar{\theta}_{23} = 40^\circ$  and  $\bar{\Delta}m_{23}^2 = 2.5 \cdot 10^{-3} \ eV^2$ 

 $(\bar{\theta}_{13}, \bar{\delta}) = (2/7^{\circ}, 45^{\circ})$ 



If  $\theta_{23}$  and  $\Delta m^2_{atm}$  can be really measured at the T2K-phase I experiment with the expected precision and for any value of  $\bar{\theta}_{23}$ , the results of two-parameters studies can be considered reliable

 $\chi^2(\theta_{13}, \delta, x) = \chi^2_{app}(\theta_{13}, \delta, x) + \chi^2_{dis}(\theta_{13}, \delta, x)$ 

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#### Conclusions on the low-gamma $\beta$ -Beam analysis

with present uncertainties, the measurement of the two unknowns in the PMNS mixing matrix is severely spoiled



errors on  $\theta_{13}$ 

<u>errors on  $\delta$ </u>

#### present uncertainties

 $\delta(\theta_{13})$  as large as  $4^{\circ}$  are found (a little bit smaller for  $\overline{\delta} = 90^{\circ}$ )

expected uncertainties

 $\delta(\theta_{13})$  as large as  $2.5 - 3^{\circ}$  are found (a little bit smaller for  $\overline{\delta} = 90^{\circ}$ ) present uncertainties half of the parameter space is spanned (a little bit smaller for  $\overline{\delta} = 90^{\circ}$ )

#### expected uncertainties

less than half of the parameter space is spanned

A significant reduction in the uncertainties on the atmospheric parameters (especially on  $\theta_{23}$ ) is mandatory if we plan to use such a facility to look for  $\theta_{13}$  (and  $\delta$ )

# **2** The $\nu_{\mu}$ disappearance channel at the SPL SB

useful to measure the atmospheric parameters

$$\begin{split} P(\nu_{\mu} \to \nu_{\mu}) &= 1 - (\sin^2 2\theta_{23} - s_{23}^2 \sin^2 2\theta_{13} \cos 2\theta_{23}) \sin^2 \left(\frac{\Delta_{atm} L}{2}\right) \\ &- \left(\frac{\Delta_{sol} L}{2}\right) [s_{12}^2 \sin^2 2\theta_{23} + \tilde{J} s_{23}^2 \cos \delta] \sin(\Delta_{atm} L) \\ &- \left(\frac{\Delta_{sol} L}{2}\right)^2 [c_{23}^4 \sin^2 2\theta_{12} + s_{12}^2 \sin^2 2\theta_{23} \cos(\Delta_{atm} L)] \end{split}$$

The main uncertainties in the measure of  $\theta_{23}$  and  $\Delta m^2_{23}$  comes from:



intrinsic clone

sign clone



# The $\nu_{\mu}$ disappearance channel at the SPL-SB

full 3-par analysis of disappearance channel:  $\chi^2(\theta_{23},\Delta m^2_{atm},\theta_{13})$ 



#### Combining APP+DIS channels at the SPL SB 3

the previous intervals on  $heta_{23}$  and  $\Delta m^2_{atm}$  are used to study the combination

 $(\bar{\theta}_{13}, \bar{\delta}) = (2/7^{\circ}, 45^{\circ})$ 



-50L

2

8

6

 $\Theta_{13}$ 

10

 $\theta_{13}$  will not be measured better than  $\sim 2^\circ$ 

### **Conclusions on the SPL SB analysis**

atmospheric parameters measurement from disappearance channel is mainly affected by *sign* ambiguity

- errors on  $\Delta m^2_{23}$  roughly doubled but similar to T2K and Enrique's analysis
- errors on  $\theta_{23}$  comparable to the current uncertainties

combining appearance and disappearance channels to measure  $\theta_{13}$  and  $\delta$ 

- degeneracies are still present but reduced in size
- physics reach: it performs slightly better than  $\beta$ -Beam in measuring  $\theta_{13}$
- half of the parameter space for  $\delta$  is spanned,  $\theta_{13}$  measured with error of  $\mathcal{O}(2^{\circ})$
- octant and mixed ambiguities cause a shift towards smaller  $\theta_{13}$  (for any value of  $\overline{\delta}$ )

A significant reduction in the uncertainties on the atmospheric parameters is mandatory if we plan to use such a facility to look for  $\delta$ 

# 4 The $\nu_e$ disappearance channel

$$P_{ee}^{\pm} = 1 - \left(\frac{\Delta_{atm}}{B_{\mp}}\right)^2 \sin^2\left(\frac{B_{\mp}L}{2}\right) \sin^2(2\theta_{13}) - \left(\frac{\Delta_{sol}}{A}\right)^2 \sin^2\left(\frac{AL}{2}\right) \sin^2(2\theta_{12})$$

It well describes the behaviour of the full transition probability in the energy range covered by the considered  $\beta$ -Beam setup

It does not depend on  $\theta_{23}$  and  $\delta$ 

Useful to measure  $\theta_{13}$ 

Two remaining sources of ambiguities that can affect the measure of  $\theta_{13}$ :

- $s_{atm}$  (for large values of  $\theta_{13}$ , i.e. in the "atmospheric" region)
- $(\theta_{13} \theta_{12})$  correlation (for small values of  $\theta_{13}$ , i.e. in the "solar" region)

$$\begin{cases} N_{\nu_{e}}^{\pm}(\bar{\theta}_{13},\bar{\theta}_{12};\bar{s}_{atm}) &= N_{\nu_{e}}^{\pm}(\theta_{13},\theta_{12};s_{atm}=\bar{s}_{atm}) & \text{intrinsic} \\ N_{\nu_{e}}^{\pm}(\bar{\theta}_{13},\bar{\theta}_{12};\bar{s}_{atm}) &= N_{\nu_{e}}^{\pm}(\theta_{13},\theta_{12};s_{atm}=-\bar{s}_{atm}) & \text{sign} \end{cases}$$

The intrinsic ambiguity is not a problem since the clone point is at:  $\theta_{13} = \bar{\theta}_{13}, \theta_{12} \sim \pi/2 - \bar{\theta}_{12}$ 

In principle the sign clone could be dangerous for values of  $\bar{\theta}_{13} < 3^{\circ}$ 

# The $\nu_e$ disappearance channel

the sign ambiguity from the  $\theta_{13} - \theta_{12}$  correlation



two different input points:  $(\bar{\theta}_{13}, \bar{\theta}_{12}) = (0.5/2^{\circ}, 32^{\circ})$ 

-red (blue) lines for neutrinos (antineutrinos)

-also plotted the 1 and 3  $\sigma$  errors on  $\theta_{12}$ 

the sign clone could in principle induce a relatively large spread on the measure of  $\theta_{13}$ 

BUT the statistics is too low to separate the true from the clone region (example for  $\bar{\theta}_{13} = 2^{\circ}$ )

-solid (dotted) line stands for intrinsic (sign) degeneracy

-background and systematics (2 %) included

the effect of the sign clone is screened by low statistics  $\Rightarrow$  the clones from  $\theta_{13} - \theta_{12}$ correlation are a marginal problem



# The $\nu_e$ disappearance channel

sensitivity to  $\theta_{13}$ 



we expect very small improvement combining app+dis channels to measure  $\theta_{13}$ , even with an "optimistic" systematic error

## Comparison with the CP-coverage P. Huber et al., hep-ph/0412199

one builds  $\chi^2(\theta_{13}, \delta, N_{\alpha})$ ,  $N_{\alpha}$  external parameters marginalization over  $N_{\alpha}$  parameters  $\psi$ for each  $\bar{\theta}_{13}, \bar{\delta}$ , a 2-dim  $\chi^2(\theta_{13}, \delta, N_{\alpha}^{min})$   $\psi$ we then minimize the resulting function in  $\theta_{13}$   $\psi$ a one-dimensional function of  $\delta$ :  $\chi^2_{min}(\delta, \bar{\delta})$   $\psi$ imposing  $\chi^2_{min}(\delta, \bar{\delta}) = CL$ , intervals  $\Delta_I(\bar{\delta})$  can be found

$$\xi(\bar{\delta}) = \text{Coverage in } \delta = \frac{1}{2\pi} U_{I=1}^{N_{deg}} \Delta_I(\bar{\delta})$$

 compact way to compare performacies of different experiments

general underestimation of the error on  $ar{\delta}$ 

we performed this procedure starting from a 4-D grid in  $\theta_{13}$ ,  $\delta$  and the atmospheric parameters and compared with the projection of our 3-D fit (vertical gray bands)



 $\beta$ -Beam fits for  $\bar{\delta}=\pm90^\circ,\pm45^\circ,0$  present uncertainties



 $\beta$ -Beam fits for  $\bar{\delta}=\pm90^\circ,\pm45^\circ,0$  expected uncertainties



SPL Super-Beam fits for  $\overline{\delta} = \pm 90^{\circ}, \pm 45^{\circ}, 0$ 



SPL Super-Beam fits for  $\overline{\delta} = \pm 90^{\circ}, \pm 45^{\circ}, 0$ 



## Physics in the first stage of the T2K project

sensitivity of the neutrino oscillation parameters

solid lines: curves for  $\sin^2(2\theta_{23}) = 1$ dashed lines: curves for  $\sin^2(2\theta_{23}) = 0.9$ 



